

# Flavor at low and high $p_T$

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## Outline

Flavor & high  $p_T$  physics interplay in several ways

- Complementary constraints on NP models from low energy precision observables vs. high  $p_T$  searches
- Nontrivial flavor structure affects collider signatures & reach
- Anomalies in B/D/K physics motivate NP searches at high p<sub>T</sub>

#### Flavor bounds on NP vs. LHC reach



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SM as EFT:  $\mathcal{L}_{\text{BSM}} \to \mathcal{L}_{\nu\text{SM}} + \sum_{i,(d>4)} \frac{\mathcal{Q}_i^{(d)}}{\Lambda^{d-4}}$ 



#### Flavor bounds on NP vs. LHC reach



## LHC bad dream scenario: (mini)split SUSY



# LHC bad dream scenario: (mini)split SUSY

30 NP thresholds neutron Kaon W. Altmannshofer et al. EDM m ixin g 1308.3653 beyond direct reach 10 Now  $\tan \beta$ EDM electron ] charm Flavor (& CPV) m ixin g 30 neutron electron EDM EDM powerful probes of 10 PeV sfermions Kaon ~2025  $an\beta$ mixing charm mixing Significant  $10^{2}$  $10^{5}$ 10  $10^{3}$  $10^{4}$  $m_{\tilde{q}} = m_{\tilde{i}} = |\mu|$  (TeV) improvements (+generic FV) **MEG upgrade** Mu3e Mu2e expected in next LHCb 25fb<sup>-1</sup> decade

 $|m_{\tilde{B}}| = |m_{\tilde{W}}| = 3 \text{ TeV}, \ |m_{\tilde{e}}| = 10 \text{ TeV}$ 

## Flavor & high-p<sub>7</sub> as complementary NP probes

Flavor safe NP? Flavor already broken in SM (Higgs).

$$\mathcal{H}_{mat} = \underbrace{\frac{c_{RL}^{IJ}}{\Lambda^n} H^{\dagger} \bar{D}^I Q^J \times X}_{A^n} + \frac{c_{LR}^{IJ}}{\Lambda^n} H \bar{Q}^I D^J \times X}_{A^n} + \underbrace{\frac{c_{RL}^{IJ}}{\Lambda^n} \bar{Q}^I Q^J \times X}_{A^n} + \underbrace{\frac{c_{RR}^{IJ}}{\Lambda^n} \bar{D}^I D^J \times X}_{A^n} + \underbrace{\frac{c_{RR}^{IJ}}{\Lambda^n} + \underbrace{\frac{c_{RR$$

Any (additional) scalar coupling to SM fermions introduces additional breaking (can be aligned with Higgs)

New (massive) vectors coupling to SM fermionic currents can preserve flavor

 $\Rightarrow$  FCNCs loop & GIM suppressed (as in SM)

see e.g. J.F.K. & C. Smith 1111.6402

$$c^{IJ} \to (g/(4\pi))^2 V_{tI}^* V_{tJ} \times c^{33}$$



D'Ambrosio et al. hep-ph/0207036

## Flavor & high- $p_T$ as complementary NP probes

*Example*: simplified DM models with (pseudo)scalar mediators

$$\mathcal{L}_{\rm DM} = i \, g_{\chi} \, A \, \bar{\chi} \gamma^5 \chi \, + \sum_{f=q,\ell,\nu} i \, g_f \, A \, \bar{f} \gamma^5 f$$

- Direct DM detection suppressed
- SM Yukawa-like couplings:  $g_f = \sqrt{2} g_Y m_f / v$
- No missing  $E_T$  signals for  $m_{\chi} > m_A/2$

# Flavor & high- $p_T$ as complementary NP probes

# *Example*: simplified DM models with (pseudo)scalar mediators



## Flavor probes of the Higgs sector

generation of masses in SM through Higgs mechanism

 $\Rightarrow$  Higgs has hierarchical couplings to fermions

$$y_f^{\rm SM} = \sqrt{2}m_f/v$$



How well have we tested this? A. Dery et al., 1302.3229

- proportionality
- factor of proportionality
- diagonality
- CP nature:

Many proposals

1306.5770, 1406.1722, 1503.04830, 1505.03870, 1505.06689, 1507.02916, 1606.09621, 1611.05463 ...

#### Flavor anomalies motivate high- $p_T$ searches

Semileptonic B anomalies: LFU in charged currents



Intriguing exp. situation: O(3σ) combined tension with SM ~20% (upwards) deviation! see also talk by Kitahara

#### Flavor anomalies motivate high- $p_T$ searches

#### Semileptonic B anomalies: neutral currents

see talks by Langenbruch, Watanuki



at O(2.5o), 20% (downwards) deviation!

$$R_{\rm K} = \frac{\int d\Gamma[B^+ \to K^+ \mu^+ \mu^-]/dq^2 \cdot dq^2}{\int d\Gamma[B^+ \to K^+ e^+ e^-]/dq^2 \cdot dq^2}$$

#### Flavor anomalies motivate high- $p_T$ searches

#### Semileptonic B anomalies: neutral currents

see talks by Langenbruch, Watanuki



### Implications for high p<sub>T</sub>: general considerations

B-anomalies in presence of (heavy) NP:

$$\mathcal{L}_{\text{BSM}} \to \mathcal{L}_{\nu\text{SM}} + \sum_{i,(d>4)} \frac{\mathcal{Q}_i^{(d)}}{\Lambda^{d-4}}$$

Deviations in flavor  $\Rightarrow$  indications of NP scale

$$[\text{scale}] = \frac{[\text{mass}]}{[\text{coupling}]}$$

Unitarity/Perturbativity

 $\Rightarrow$  upper bound on coupling

Di Luzio & Nardecchia, 1706.01868

 $\Rightarrow$  upper bound on NP d.o.f. mass

### Implications for high p<sub>T</sub>: general considerations

LFUV in  $R(D^{(*)})$ :  $\Lambda \simeq 2.5 \,\text{TeV}$  e.g.  $\mathcal{Q} = (\bar{c}\gamma_{\mu}P_{L}b)(\bar{\tau}\gamma_{\mu}P_{L}\nu)$ 

 $\Rightarrow$  tree-unitarity  $M_{\rm NP} \lesssim 6.5 {\rm TeV}$ 

#### up to the edge of LHC kinematical reach

see e.g. Altmannshofer et al., 1704.06659 Iguro et al., 1810.05843

LFUV in  $R_{K^{(*)}}$  (& other obs.) :  $\Lambda \sim 40 \text{TeV}$ 

e.g.  $\mathcal{Q} = (\bar{s}\gamma_{\mu}P_{L}b)(\bar{\mu}\gamma^{\mu}P_{L}\mu)$ 

 $\Rightarrow$  NP d.o.f.s accessible at LHC only if their couplings to bs and/or  $\mu\mu$  suppressed!

## Implications of LFUV for NP flavor breaking

NP needs to respect SM gauge symmetry

Si

 $\mathcal{Q}_i[Q, D, U, L, E]$ 

At EW scale: in terms of four-fermion operators

$$R_{K^{(*)}} \underbrace{\left( \begin{array}{c} \epsilon_{ij}^{L} \epsilon_{kl}^{Q}(\bar{L}_{i}L_{j})(\bar{Q}_{k}Q_{l}) \\ \epsilon_{ij}^{E} \epsilon_{kl}^{Q}(\bar{E}_{i}E_{j})(\bar{Q}_{k}Q_{l}) \end{array}}_{\epsilon_{ij}^{E} \epsilon_{kl}^{Q}(\bar{E}_{i}E_{j})(\bar{Q}_{k}Q_{l})} \underbrace{\left( \begin{array}{c} \epsilon_{ij}^{EL} \epsilon_{kl}^{QD}(\bar{E}_{i}H^{\dagger}L_{j})(\bar{Q}_{k}HD_{l}) \\ \epsilon_{ij}^{LE} \epsilon_{kl}^{QU}(\bar{L}_{i}HE_{j})(\bar{Q}_{k}\tilde{H}U_{l}) \end{array} \right)}_{\epsilon_{ij}^{E} \epsilon_{kl}^{E} \epsilon_{kl}^{QU}(\bar{L}_{i}HE_{j})(\bar{Q}_{k}\tilde{H}U_{l})} R(D^{(*)}) \\ R_{L} \\ R_{L}$$

#### Application of crossing - mono-tau production @ LHC

A. Greljo et al., 1811.07920



$$\sigma \sim \mathcal{L} \times \hat{\sigma}$$

★ heavy flavour pdf. suppression  $\mathcal{L}_{ij} \propto \int f_i(x) f_j(\hat{s}/sx) dx$ ✓ compensated by high partonic energy  $\hat{\sigma} \propto (\hat{s}/M^2)^2$ 

#### Application of crossing - mono-tau production @ LHC



Excludes W<sub>R</sub>-v<sub>R</sub> explanation, constrains some LQ scenarios

## Implications of LFUV for NP flavor breaking

# Absence of BSM LFUV, FCNCs in Kaon, Charm, Tau decays requires approximate alignment with the 3rd generation

Fajfer et al., 1206.1872 Bordone et al., 1702.07238



## Immediate implications for LHC

Flavor alignment implies lower NP scale:

$$(\bar{Q}_3 Q_3)(\bar{L}_3 L_3) \to V_{cb}(\bar{c}b)(\bar{\tau}\nu)$$
  
 $\Rightarrow R(D^{(*)}) \text{ anomaly}$   
 $\Lambda \sqrt{|V_{cb}|} \sim 500 \,\text{GeV}$ 

 $(\bar{Q}_3 Q_3)(\bar{L}_2 L_2) \rightarrow V_{tb} V_{ts}(\bar{s}b)(\bar{\mu}\mu)$  $\Rightarrow R_{K^{(*)}}$  anomaly  $\Lambda \sqrt{|V_{ts}|} \sim 8 \,\text{TeV}$ 

Well within LHC reach!

Still only marginally!

Enhanced LFUV in top processes: LHC phenomenology: bb →

 $(\bar{Q}_3Q_3)(\bar{L}_3L_3) \to V_{cb}(\bar{c}b)(\bar{c}c) \to V_{cb}(\bar{c}c)(\bar{c}c) \to V_{cb}(\bar{c}c)(\bar{c}c$ 

Currently tested to O(10%)

 $\mathcal{B}_e = 13.3(4)(4)\%, \ \mathcal{B}_\mu = 13.4(3)(5)\%, \ \mathcal{B}_{\tau_h} = 7.0(3)(5)\%,$ 

ATLAS, 1506.05074



Enhanced LFUV in top processes:

$$(\bar{Q}_3Q_3)(\bar{L}_3L_3) \to V_{cb}(\bar{c}b)(\bar{\tau}\nu) + V_{tb}(\bar{t}b)(\bar{\tau}\nu)$$

Currently tested to O(10%)

 $\mathcal{B}_e = 13.3(4)(4)\%$ ,  $\mathcal{B}_\mu = 13.4(3)(5)\%$ ,  $\mathcal{B}_{\tau_h} = 7.0(3)(5)\%$ ,

ATLAS, 1506.05074



Weak gauge invariance  $\Rightarrow$  neutral currents

 $(\bar{Q}_3Q_3)(\bar{L}_3L_3) \to V_{cb}(\bar{c}b)(\bar{\tau}\nu) + V_{tb}(\bar{t}b)(\bar{\tau}\nu) + (\bar{b}b)(\bar{\tau}\tau)$ 



W'/Z' explanation only allowed if light (M < 400 GeV) or broad ( $\Gamma/M > 20\%$ )

Leptoquark, charged scalar explanations disfavored

Departures from strict  $U(2)_F$  limit can ameliorate the bounds

see e.g. Buttazzo et al., 1706.07808 Importance of EW radiative corrections: B-anomalies without new quark flavor violation

Starting with flavor conserving non-universal operators:

 $(\bar{L}_2 L_2)(\bar{U}_3 U_3)$   $(\bar{E}_2 E_2)(\bar{U}_3 U_3)$ 

EW matching & RGE induce LFUV in rare FCNC B decays

Aebischer et al., 1512.02830 Faroughy et al., 1805.04917

see also Blanger, Delaunay, & Westhoff, 1507.06660 Bauer et al., 1511.01900 Becirevic & Sumensari, 1704.05835

Effective NP scale now loop-suppressed:

 $\frac{\sqrt{|V_{ts}|}}{4\pi} \sim 600 \,\mathrm{GeV}$ 

 $\Rightarrow$  automatically respects 3rd gen. alignment

 $\Rightarrow$  d.o.f.'s mediating R<sub>K</sub> well within LHC kinematical reach

# $R_{K^{(*)}}$ without new flavor violation: a UV completion

JFK, Soreq & Zupan, 1704.06005

VL quark partner of right-handed top (T),

- charged under gauged U(1)'(Z', h')
- T- $t_R$  mix after U(1)' breaking induced U(1)' charge of  $t_R$

(similar mechanism possible to induce muon U(1)' charge)



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Flavor is powerful guide to high- $p_T$  searches at LHC:

 In case new phenomena are discovered at LHC, flavor physics will allow to disentangle different possible interpretations and discriminate between different proposals and scenarios

*Example*: 125GeV Higgs)

 In case no new d.o.f.s are seen at LHC, precision tests of flavor, CP, B & L possibly best probes forward

⇒ their sensitivity in many cases already (far) exceeds energies/scales attainable in present and planned collider & cosmic ray experiments. Flavor is powerful guide to high- $p_T$  searches at LHC:

 In case of significant signals of NP in flavor observables can identify prospective LHC experimental targets

Generally, NP d.o.f.'s accommodating tentative B-anomalies could be beyond LHC reach.

Phenomenological and model-building considerations point towards more optimistic scenarios

- Low energy constraints can point to lighter mediators!
- Example of fruitful interplay between NP searches at energy and intensity frontiers