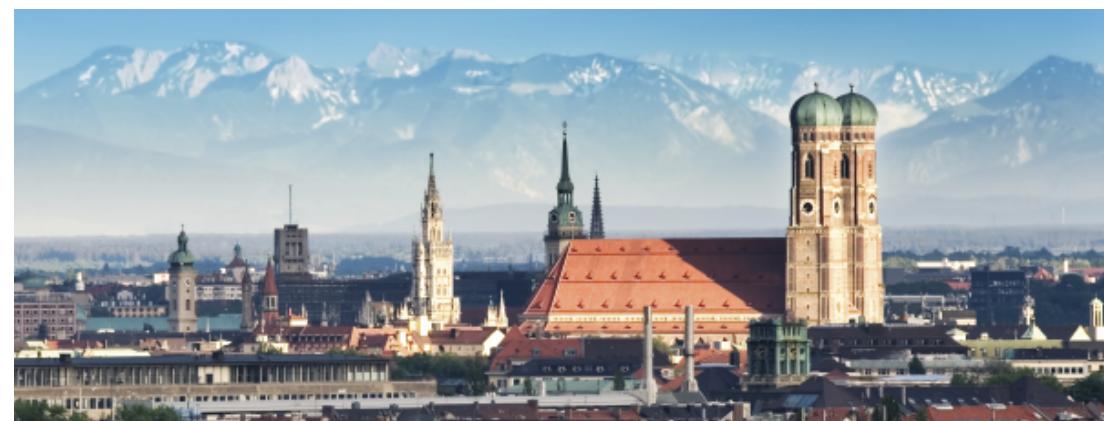
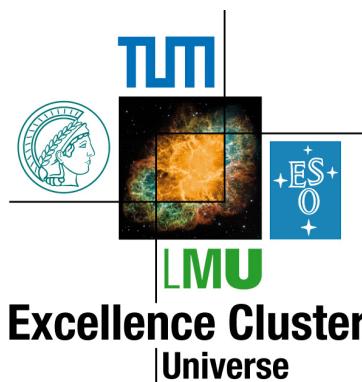


EFT, global analyses, model (in)dependence, and all that

Marco Ciuchini

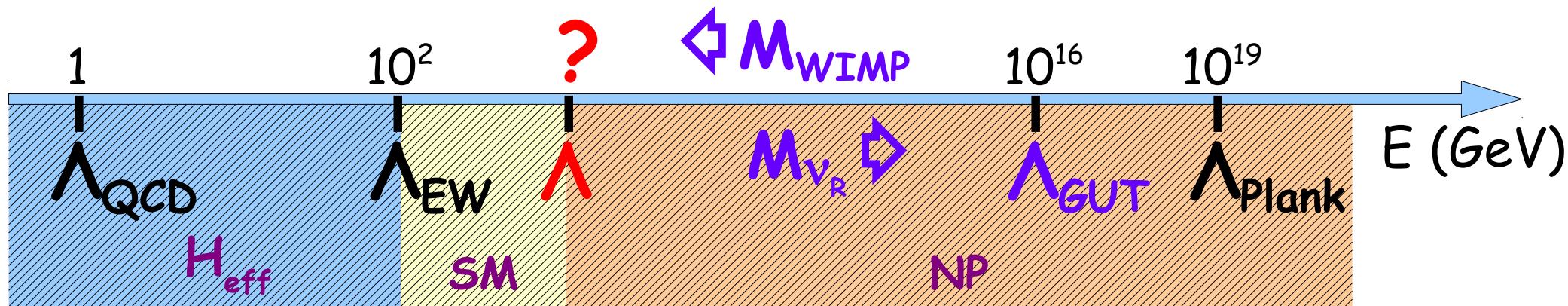
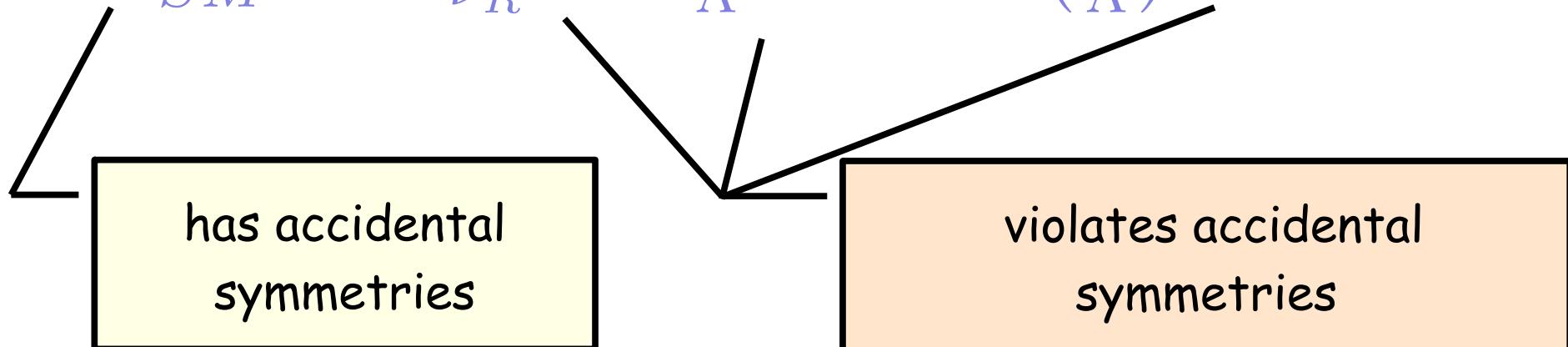


- EFT vs NP models: the UTA example
- From H_{eff} to SMEFT: C_7 - $C_{9/10}$ global analysis
- HEPfit: a tool for combining direct and indirect constraints on HEP models



Standard Model as an EFT

$$\mathcal{L} = \mathcal{L}_{SM}^{D=4} + \mathcal{L}_{\nu_R}^{D=4} + \frac{1}{\Lambda} \mathcal{L}^{D=5} + \left(\frac{1}{\Lambda}\right)^2 \mathcal{L}^{D=6} + \dots$$



$$\rightarrow \mathcal{H}_{eff}^{SM} + \mathcal{H}_{eff}^{NP}$$

$E \ll \Lambda_{EW}$

$SU(3)_c \times U(1)_q$
-invariant, i.e.
low-energy symmetries

When does this construction work?

i) Scale separation

There must be a NP scale much larger than the EW scale (and of the scale of the process)

↳ at least one scale is there: the Plank scale

challenged: e.g. Salvio, Strumia, 1403.4226

ii) Full knowledge of the low-energy physics

The particle content below the NP scale is known, i.e. all low-energy d.o.f.'s are included in the EFT

*challenged: feebly coupled light particles
(dark sectors, portal, etc.)*

Top-down or bottom-up?

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = \frac{1}{\Lambda^2} \sum_i C_i Q_i$$

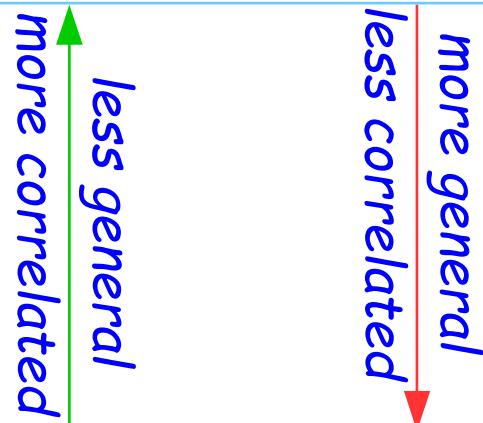
Top-down: a NP correlation pattern is enforced on the WC's and then looked for in the data

Bottom-up: only low-energy correlations are enforced, data fix WCs', hopefully showing additional patterns

Both approaches require the effective theory

Specific NP model

- C_i and Λ are fully calculable in terms of the model parameters



Model-independent analysis

- the free parameters are C_i / Λ^2
- Λ cannot be determined

A well-known example: NP in the UT

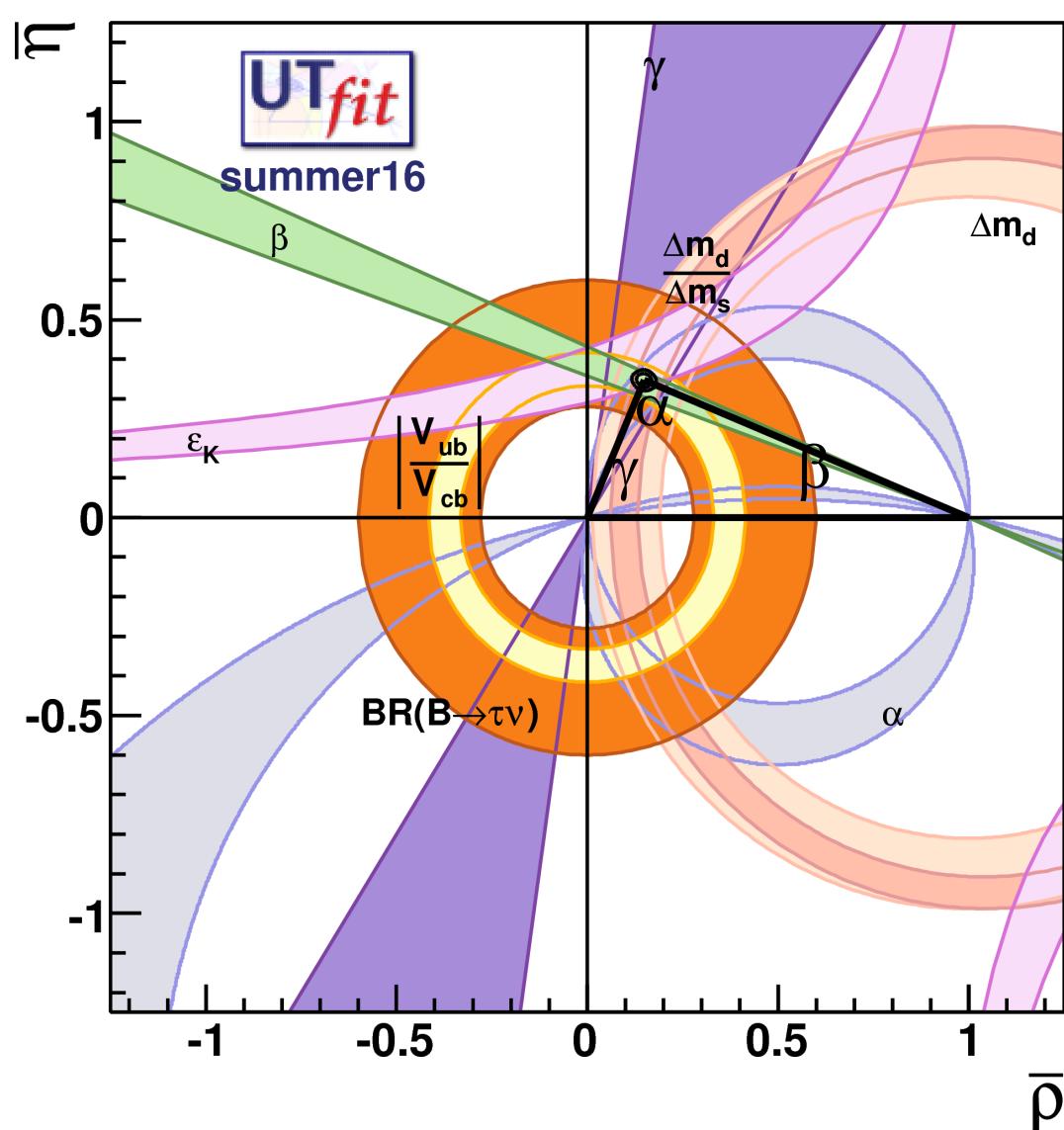
The UT analysis is mainly driven by $\Delta F=2$ transitions and usually done using SM amplitudes

NP can be include with an increasing degree of model-dependence:

Step 0: Phenomenological parametrization

Step 1: EFT analysis

Step 2: Explicit model (MSSM)



0- Phenomenological parametrization(s)

B_d and B_s mixing amplitudes (2+2 real parameters):

$$A_q e^{2i\phi_q} = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

Tree processes

	ρ, η	C_d	φ_d	C_s	φ_s	C_{eK}
γ (DK)	X					
V_{ub}/V_{cb}	X					
Δm_d	X	X				
ACP (J/ Ψ K)	X		X			
ACP (D $\pi(p)$, DK π)	X		X			
A_{SL}		X	X			
α ($\rho\rho, \rho\pi, \pi\pi$)	X		X			
A_{CH}		X	X	X	X	
$\tau(B_s), \Delta\Gamma_s/\Gamma_s$				X	X	
Δm_s				X		
ASL(B_s)				X	X	
ACP (J/ Ψ ϕ)	-X				X	
ε_K	X					X

1 \leftrightarrow 3 family

2 \leftrightarrow 3 family

1 \leftrightarrow 2 family

$q=d, s, \phi_d^{SM}=\beta, \phi_s^{SM}=-\beta_s$

+

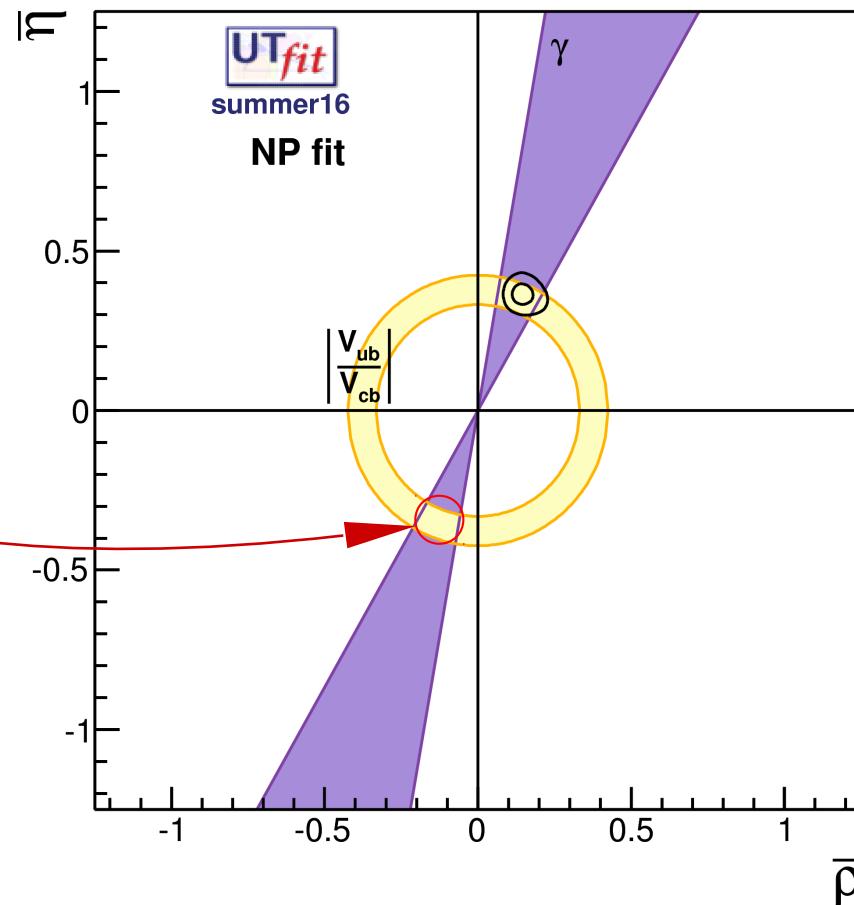
K mixing amplitude (1 real param):

$$\text{Im } A_K = C_\varepsilon \text{Im } A_K^{SM}$$

Correlations at work

assumption: negligible NP in tree decays

excluded by correlations



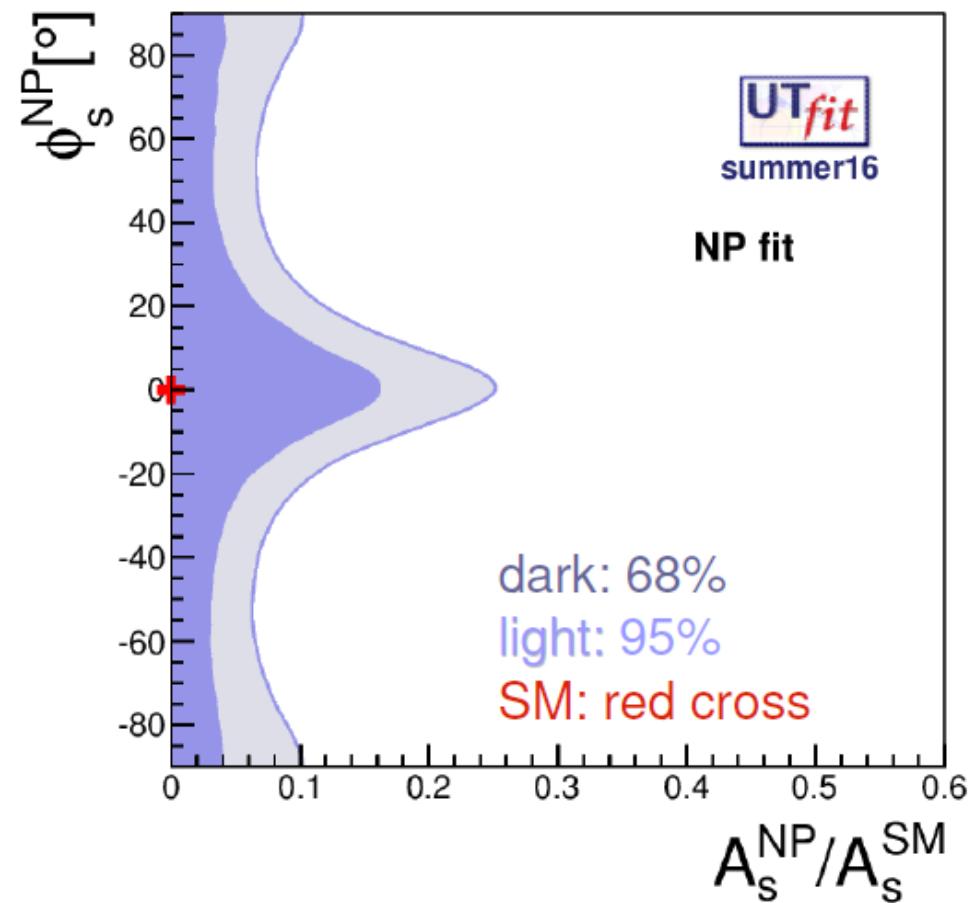
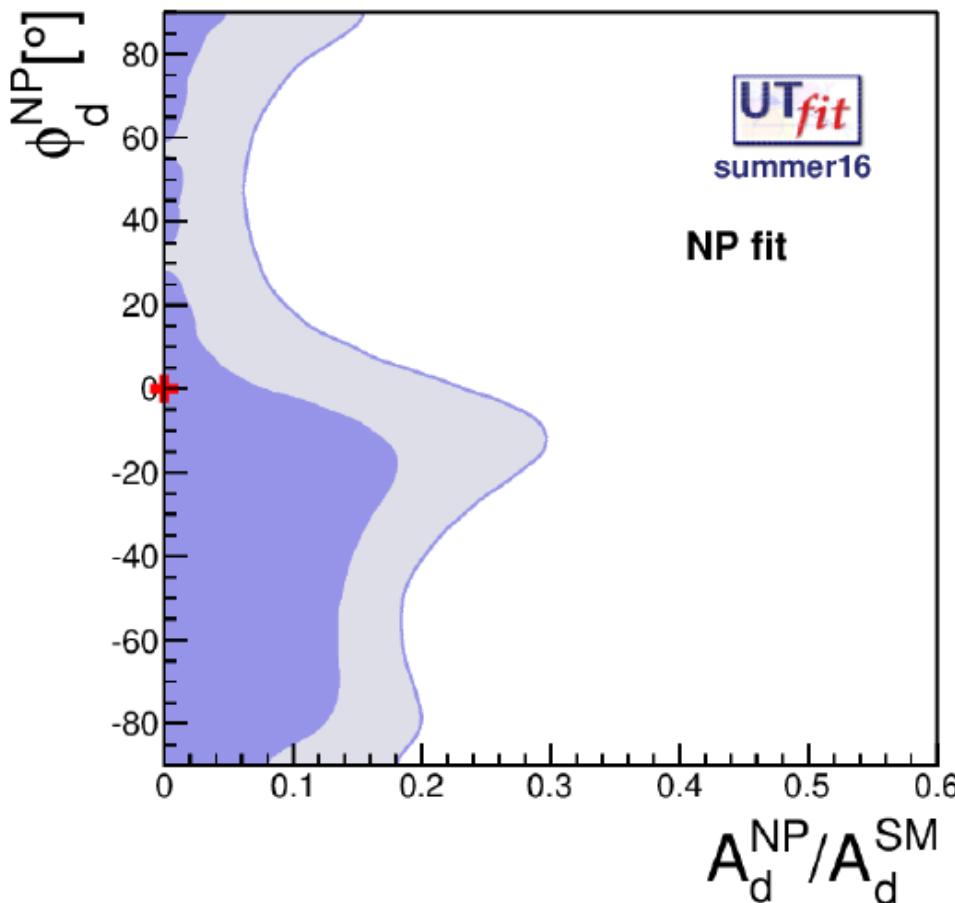
$$\bar{\rho} = 0.147 \pm 0.043$$

$$\bar{\eta} = 0.384 \pm 0.044$$

in the SM fit is:

$$\bar{\rho} = 0.142 \pm 0.018$$
$$\bar{\eta} = 0.357 \pm 0.013$$

Implications for the NP amplitudes



The ratio of NP/SM amplitudes is:

- < ~1-20% @68% prob. (5-30% @95%) in B_d mixing
- < ~3-15% @68% prob. (5-25% @95%) in B_s mixing

1 - EFT analysis of $\Delta F=2$ transitions

The mixing amplitudes $A_q e^{2i\phi_q} = \langle \bar{M}_q | H_{\text{eff}}^{\Delta F=2} | M_q \rangle$

$$H_{\text{eff}}^{\Delta B=2} = \sum_{i=1}^5 C_i(\mu) Q_i(\mu) + \sum_{i=1}^3 \tilde{C}_i(\mu) \tilde{Q}_i(\mu)$$

$$Q_1 = \bar{q}_L^\alpha \gamma_\mu b_L^\alpha \bar{q}_L^\beta \gamma^\mu b_L^\beta \quad (\text{SM/MFV})$$

$$Q_2 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_R^\beta b_L^\beta$$

$$Q_3 = \bar{q}_R^\alpha b_L^\beta \bar{q}_R^\beta b_L^\beta$$

$$Q_4 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_L^\beta b_R^\beta$$

$$Q_5 = \bar{q}_R^\alpha b_L^\beta \bar{q}_L^\beta b_R^\beta$$

$$\tilde{Q}_1 = \bar{q}_R^\alpha \gamma_\mu b_R^\alpha \bar{q}_R^\beta \gamma^\mu b_R^\beta$$

$$\tilde{Q}_2 = \bar{q}_L^\alpha b_R^\alpha \bar{q}_L^\beta b_R^\beta$$

$$\tilde{Q}_3 = \bar{q}_L^\alpha b_R^\beta \bar{q}_L^\beta b_R^\beta$$

$C_i(\Lambda)$ are extracted from data:

- * one by one, "barring accidental cancellations"
- * all together, safer but trickier

Parameter	95% allowed range (GeV^{-2})
$\text{Im}C_K^1$	$[-1.8, 2.5] \cdot 10^{-15}$
$\text{Im}C_K^2$	$[-1.7, 1.2] \cdot 10^{-17}$
$\text{Im}C_K^3$	$[-1.7, 2.5] \cdot 10^{-16}$
$\text{Im}C_K^4$	$[-3.0, 4.3] \cdot 10^{-18}$
$\text{Im}C_K^5$	$[-5.5, 8.1] \cdot 10^{-17}$
$\text{Im}C_D^1$	$[-1.4, 2.0] \cdot 10^{-14}$
$\text{Im}C_D^2$	$[-2.5, 1.7] \cdot 10^{-15}$
$\text{Im}C_D^3$	$[-2.4, 3.5] \cdot 10^{-14}$
$\text{Im}C_D^4$	$[-5.2, 7.7] \cdot 10^{-16}$
$\text{Im}C_D^5$	$[-5.3, 7.9] \cdot 10^{-15}$
$ C_{B_d}^1 $	$< 1.4 \cdot 10^{-12}$
$ C_{B_d}^2 $	$< 3.2 \cdot 10^{-13}$
$ C_{B_d}^3 $	$< 1.2 \cdot 10^{-12}$
$ C_{B_d}^4 $	$< 1.0 \cdot 10^{-13}$
$ C_{B_d}^5 $	$< 2.9 \cdot 10^{-13}$
$ C_{B_s}^1 $	$< 2.3 \cdot 10^{-11}$
$ C_{B_s}^2 $	$< 5.0 \cdot 10^{-12}$
$ C_{B_s}^3 $	$< 1.9 \cdot 10^{-11}$
$ C_{B_s}^4 $	$< 1.7 \cdot 10^{-12}$
$ C_{B_s}^5 $	$< 4.6 \cdot 10^{-12}$

Lower bound on the NP scale Λ from $\Delta F=2$ transitions (TeV @95% prob.)

Already beyond the EFT: it needs assumptions on the UV structure

$$\Lambda = \sqrt{\frac{L \cdot FC}{C_i(\Lambda)}}$$

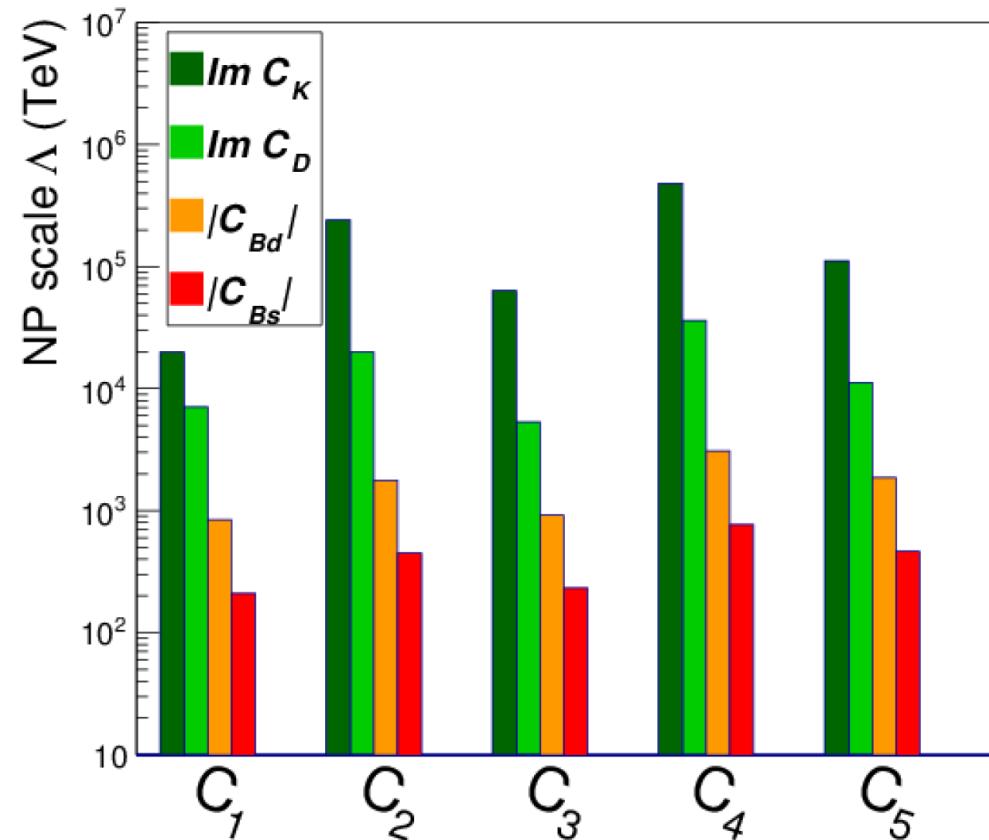
Loop factor L:

tree/strong interact. NP, $L \sim 1$

if perturbative: $L \sim \alpha_s^2, \alpha_W^2$

Flavor couplings FC:

$|FC| \sim 1$, arbitrary phases; ...

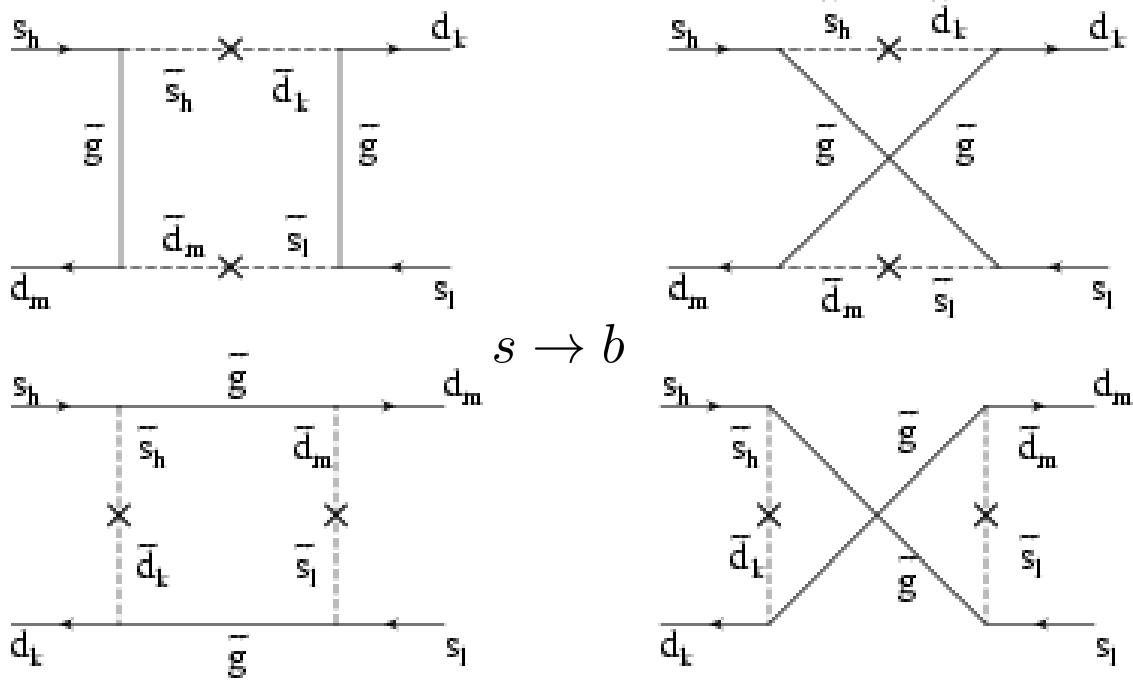


Λ (TeV)	K CPV	D CPV	B_d CPC	B_s CPC
-----------------	-------	-------	-----------	-----------

lower bound	4.8×10^5	3.6×10^4	3.1×10^3	760
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for the case of loop and flavour NP couplings equal to 1

2- MSSM: gluino-squark contributions

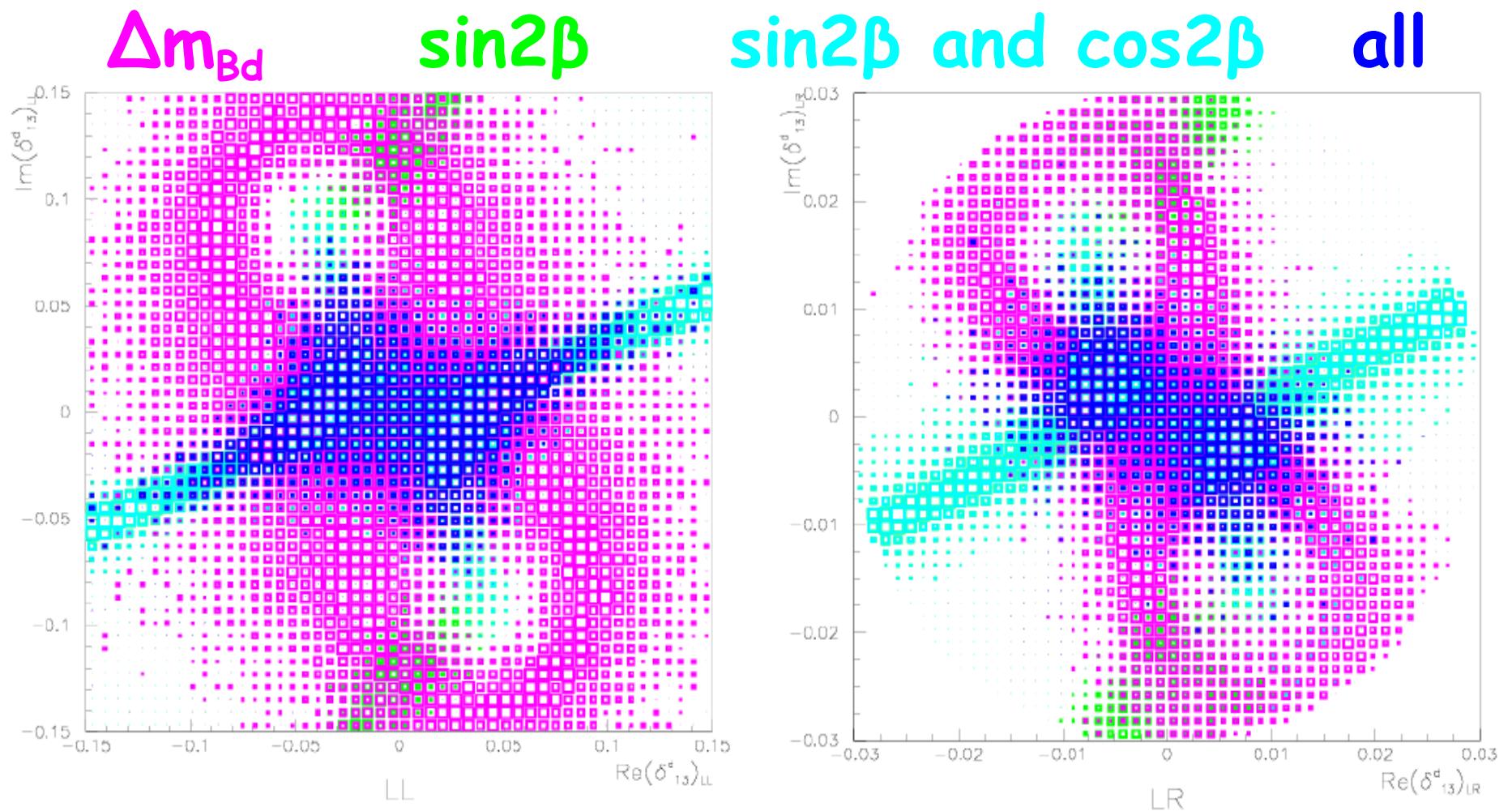


Gabbiani et al.,
hep-ph/9604387

$$\begin{aligned}
 C_1 &= \alpha_s^2 \frac{(\delta_{13}^d)_{LL}^2}{\tilde{m}^2} f_1(x) & C_2 &= \alpha_s^2 \frac{(\delta_{13}^d)_{RL}^2}{\tilde{m}^2} f_2(x) & C_3 &= \alpha_s^2 \frac{(\delta_{13}^d)_{RL}^2}{\tilde{m}^2} f_3(x) \\
 \tilde{C}_1 &= \alpha_s^2 \frac{(\delta_{13}^d)_{RR}^2}{\tilde{m}^2} f_1(x) & \tilde{C}_2 &= \alpha_s^2 \frac{(\delta_{13}^d)_{LR}^2}{\tilde{m}^2} f_2(x) & \tilde{C}_3 &= \alpha_s^2 \frac{(\delta_{13}^d)_{LR}^2}{\tilde{m}^2} f_3(x) \\
 C_4 &= \alpha_s^2 \left[\frac{(\delta_{13}^d)_{LL} (\delta_{13}^d)_{RR}}{\tilde{m}^2} f_4(x) + \frac{(\delta_{13}^d)_{LR} (\delta_{13}^d)_{RL}}{\tilde{m}^2} \tilde{f}_4(x) \right] \\
 C_5 &= \alpha_s^2 \left[\frac{(\delta_{13}^d)_{LL} (\delta_{13}^d)_{RR}}{\tilde{m}^2} f_5(x) + \frac{(\delta_{13}^d)_{LR} (\delta_{13}^d)_{RL}}{\tilde{m}^2} \tilde{f}_5(x) \right]
 \end{aligned}$$

$$x = \frac{m_{\tilde{g}}^2}{\tilde{m}}$$

NP model:
 → less parameters
 → more correlations



$\text{Re}(\delta_{13}^d)_{\text{LL,RR}}$ vs $\text{Im}(\delta_{13}^d)_{\text{LL,RR}}$

$\text{Re}(\delta_{13}^d)_{\text{LR,RL}}$ vs $\text{Im}(\delta_{13}^d)_{\text{LR,RL}}$

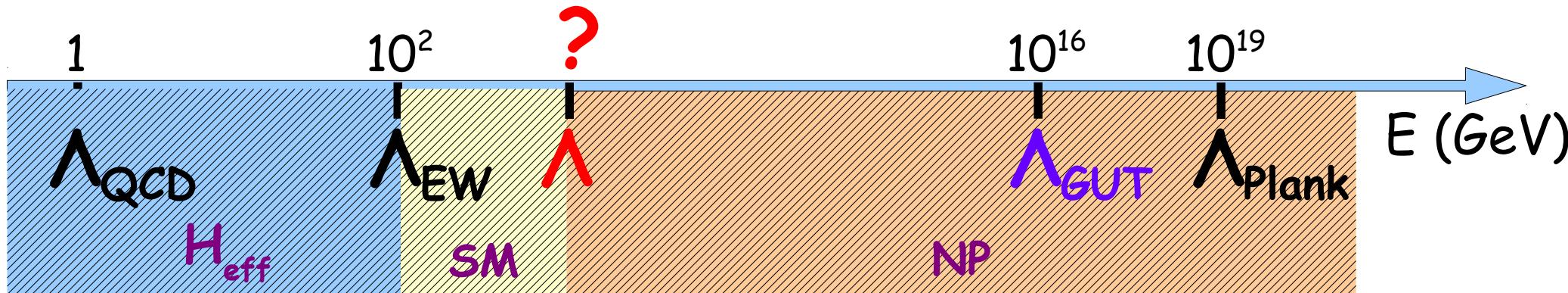
MC et al, hep-ph/9808328 + updates

$m_{sq}=m_g=-\mu=350 \text{ GeV} (!!), \text{ scale as } m_{sq}/350 \text{ GeV}$

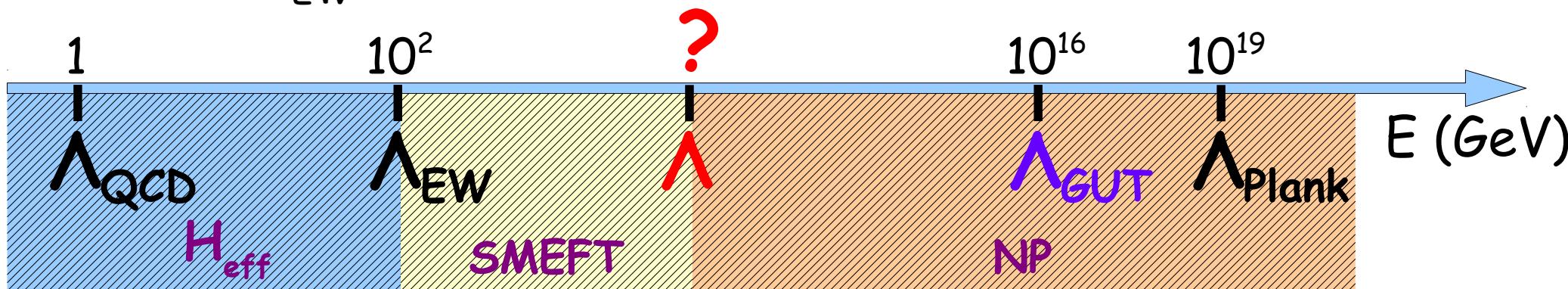
$ \delta_{13}^d _{\text{LL,RR}}$	$ \delta_{13}^d _{\text{LL=RR}}$	$ \delta_{13}^d _{\text{LR}}$	$ \delta_{13}^d _{\text{RL}}$
$7 \cdot 10^{-2}$	$5 \cdot 10^{-3}$	$1 \cdot 10^{-2}$	$1 \cdot 10^{-2}$

A step upward: SMEFT

$$\mathcal{L} = \mathcal{L}_{SM}^{D=4} + \mathcal{L}_{\nu_R}^{D=4} + \frac{1}{\Lambda} \mathcal{L}^{D=5} + \left(\frac{1}{\Lambda}\right)^2 \mathcal{L}^{D=6} + \dots$$



If $\Lambda \gg \Lambda_{EW}$ (as LHC data might suggest):



$\rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y$ EFT $\rightarrow SU(3)_c \times U(1)_q$ EFT
 $E \ll \Lambda$ $E \ll \Lambda_{EW}$

Facts about the SMEFT

see talks by Bobeth, Gonzalez-Alonso, and Straub

- ▶ 2499 operators (59 operators considering only flavour-diagonal and family-universal) built out of $SU(2)\times U(1)$ multiplets (including the Higgs doublet)
- ▶ It allows for studying SM processes at large (e.g. EWPO, Higgs couplings, flavour, ...) Buchmuller & Wyler;
Grzadkowski et al;
- ▶ Once matched onto the usual weak H_{eff} , it induces additional constraints and correlations on the WC's

CAVEAT:

- ▶ LHC may still prove that Λ is not much larger than Λ_{EW}
- ▶ No additional constraint on the H_{eff} WC's if the EWSB is not linearly realized Catà, Jung

EWSB and low-energy processes

It is another floor added to the tower of EFT's built going from low to high energies (or vice versa).

As energy much larger than the EW scale, the EW symmetry is restored.

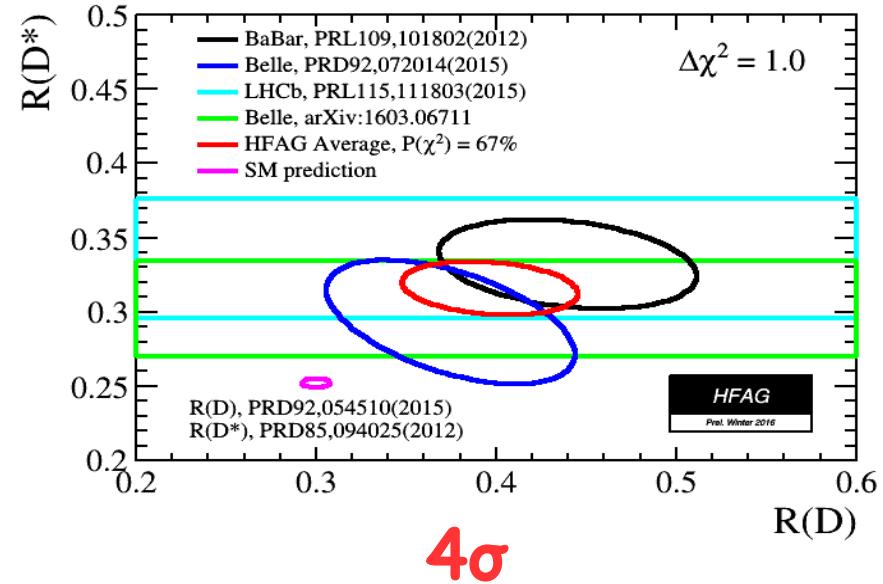
It is found that low-energy processes are sensitive to the details of EWSB (!!)

Catà, Jung,
arXiv:1505.05804

For instance, if studying $b \rightarrow s \ell \ell$ transitions, it is found that the Wilson coefficient of the operator $(\bar{s} \sigma_{\mu\nu} b) \bar{l} \sigma^{\mu\nu} l$ is non-vanishing, then not only the presence of NP is established, but also the non-standard character of the EWSB is ascertained

SMEFT & flavour anomalies

V.Vaanoni, ICHEP



$$R(X) = \frac{\Gamma(B \rightarrow X\tau\nu)}{\Gamma(B \rightarrow X\ell\nu)}$$

$$R(D) = -1.9\sigma$$

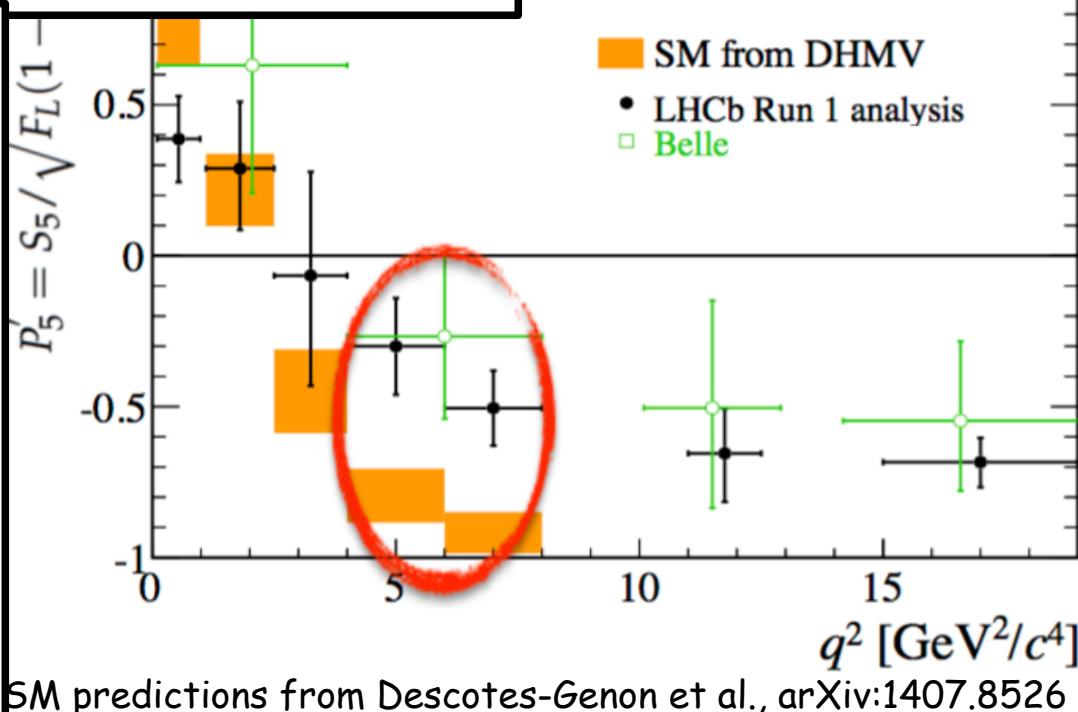
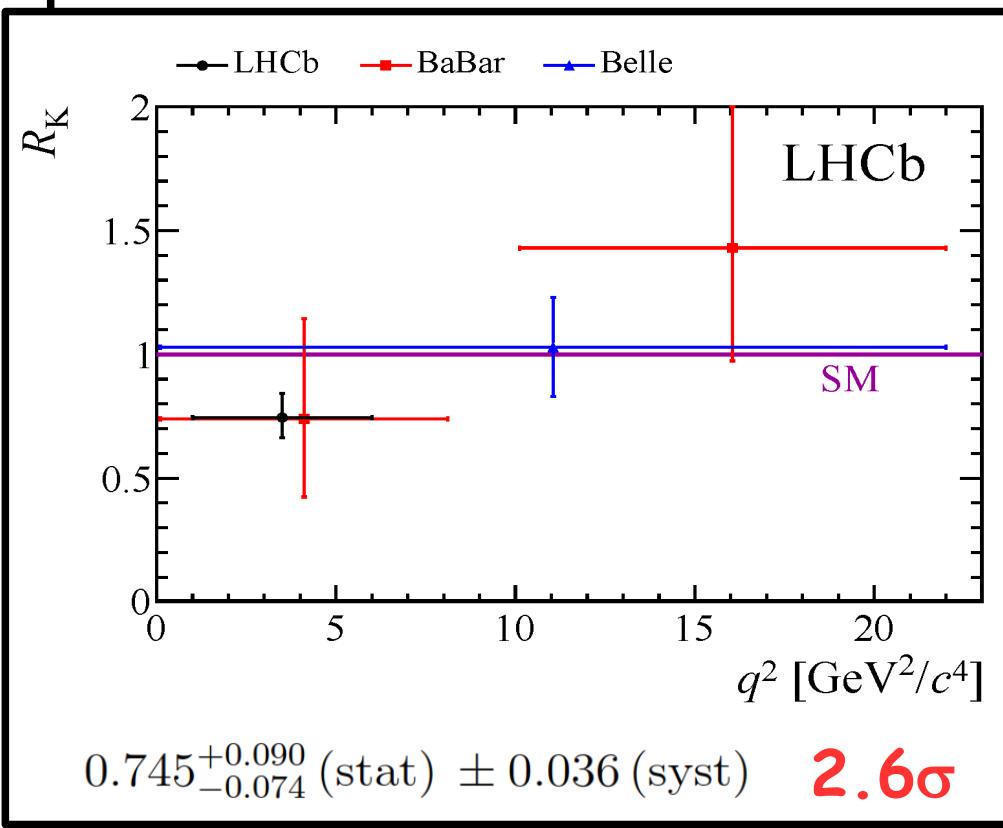
$$0.397 \pm 0.040 \pm 0.028$$

$$R(D)_{\text{SM}} = 0.300 \pm 0.008$$

$$R(D^*) = -3.3\sigma$$

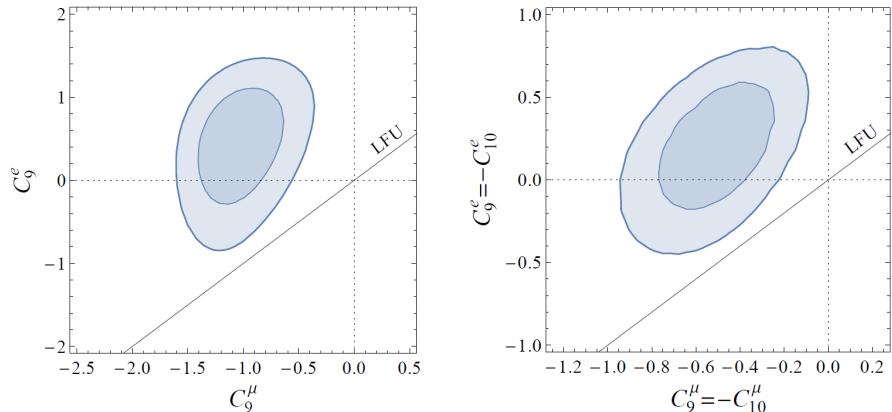
$$0.316 \pm 0.016 \pm 0.010$$

$$R(D^*)_{\text{SM}} = 0.252 \pm 0.003$$

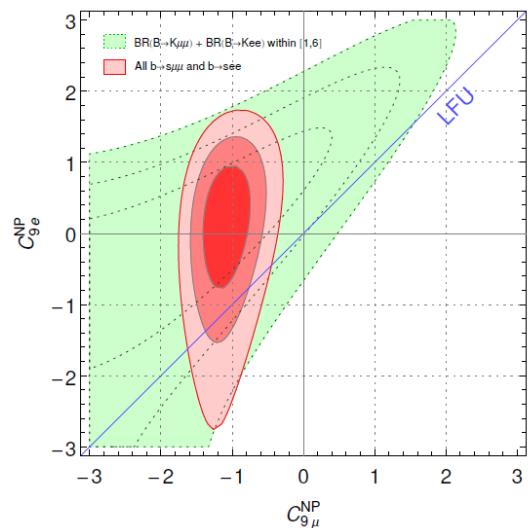


EFT global analysis

Altmannshofer, Straub., arXiv:1411.3161

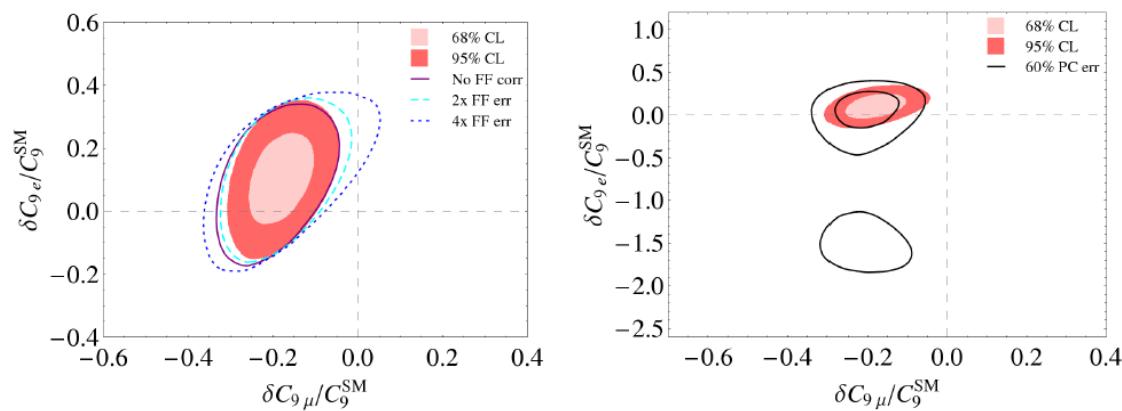


Descotes-Genon et al., arXiv:1605.06059



- * $B \rightarrow K^{(*)} \mu\mu$
- * $B \rightarrow X_s \gamma$
- * $B_s \rightarrow \phi \mu\mu$
- * R_K
- * $B \rightarrow K^* \gamma$

Hurth et al., arXiv:1603.00865



point to an $O(1)$ correction to
the WC of $Q_9^\mu = \bar{s}_L \gamma_\alpha b_L \bar{\mu} \gamma^\alpha \mu$

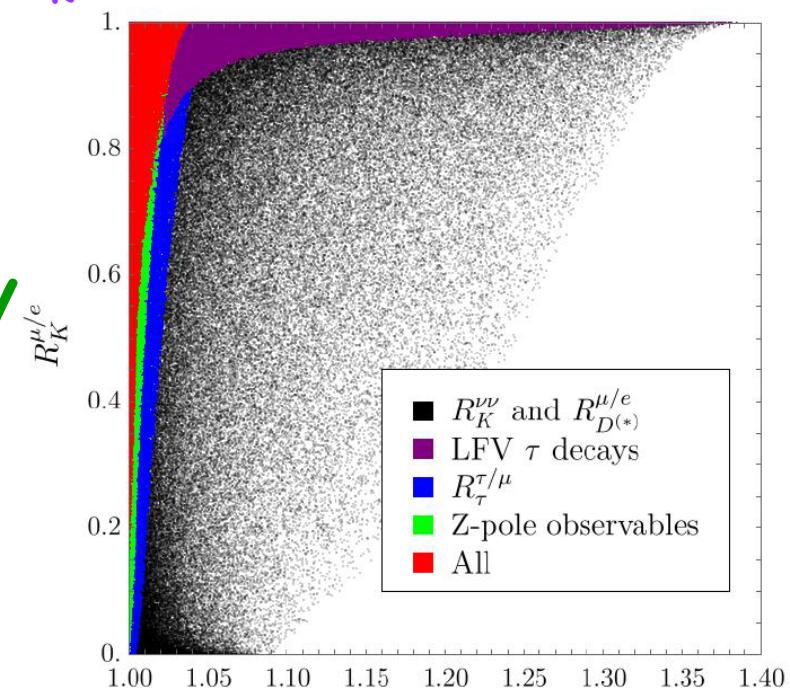
If LFU is induced by NP coupled to the 3rd generation at a scale much larger than the weak scale, in the SMEFT there are 2 four-fermion operator candidates:

$$Q'_{L,3}\gamma_\mu Q'_{L,3}L'_{L,3}\gamma^\mu L'_{L,3}, \quad Q'_{L,3}\gamma_\mu\sigma^i Q'_{L,3}L'_{L,3}\gamma^\mu\sigma^i L'_{L,3}$$

i) give typically rise to large LFV Glashow et al, arXiv:1411.0565

ii) can account for the anomalies in R_K , $R(D)$ & $R(D^*)$ Bhattacharya et al, arXiv:1412.7164

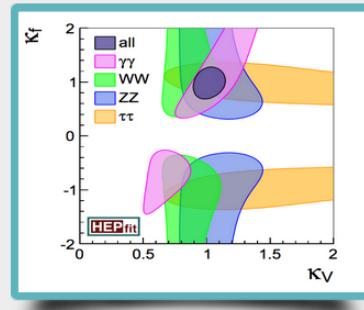
iii) RGE running in SMEFT produces LFV corrections to the $V_{l\ell}$ vertices and induces large purely leptonic FV transitions, which disfavour a common explanation of the anomalies Feruglio et al, arXiv:1606.00524



what about RGE effects in models?

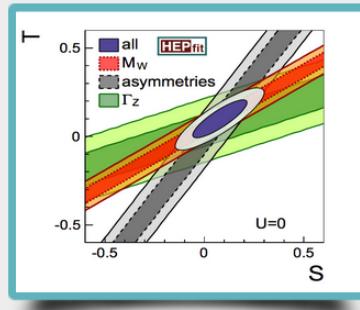
$$R_{D^{(*)}}^{\tau/\ell}$$

HEPfit: a Code for the Combination of Indirect and Direct Constraints on High Energy Physics Models.



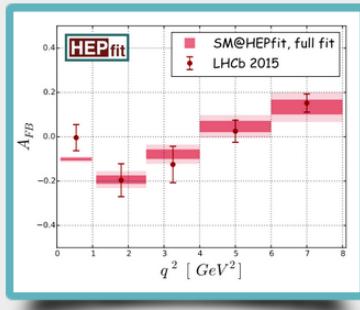
Higgs Physics

HEPfit can be used to study Higgs couplings and analyze data on signal strengths.



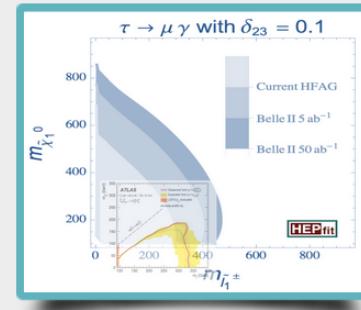
Precision Electroweak

Electroweak precision observables are included in **HEPfit**.



Flavour Physics

The Flavour Physics menu in **HEPfit** includes both quark and lepton flavour dynamics.



BSM Physics

Dynamics beyond the Standard Model can be studied by adding models in **HEPfit**.

Support

Support email: [hepfit-support\[at\]roma1.infn.it](mailto:hepfit-support[at]roma1.infn.it).

You can also connect to us through our social network pages linked below.

The HEPfit Collaboration

Downloads

Current Version: **HEPfit** v1.0-RC1

Developer Version: **HEPfit** @ GitHub

Previous Versions:

Dependencies: GSL, ROOT, BOOST, BAT

The HEPfit team



Rome I&III

Jorge de Blas

Debtosh Chowdhury

Marco Ciuchini

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Marco Fedele

Enrico Franco

Ayan Paul

Luca Silvestrini

SISSA Trieste

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KEK

Satoshi Mishima

CERN

Maurizio Pierini

Florida State University

Laura Reina

Tohoku University

Norimi Yokozaki

Lanzhou University

Fu-Sheng Yu

+ you!

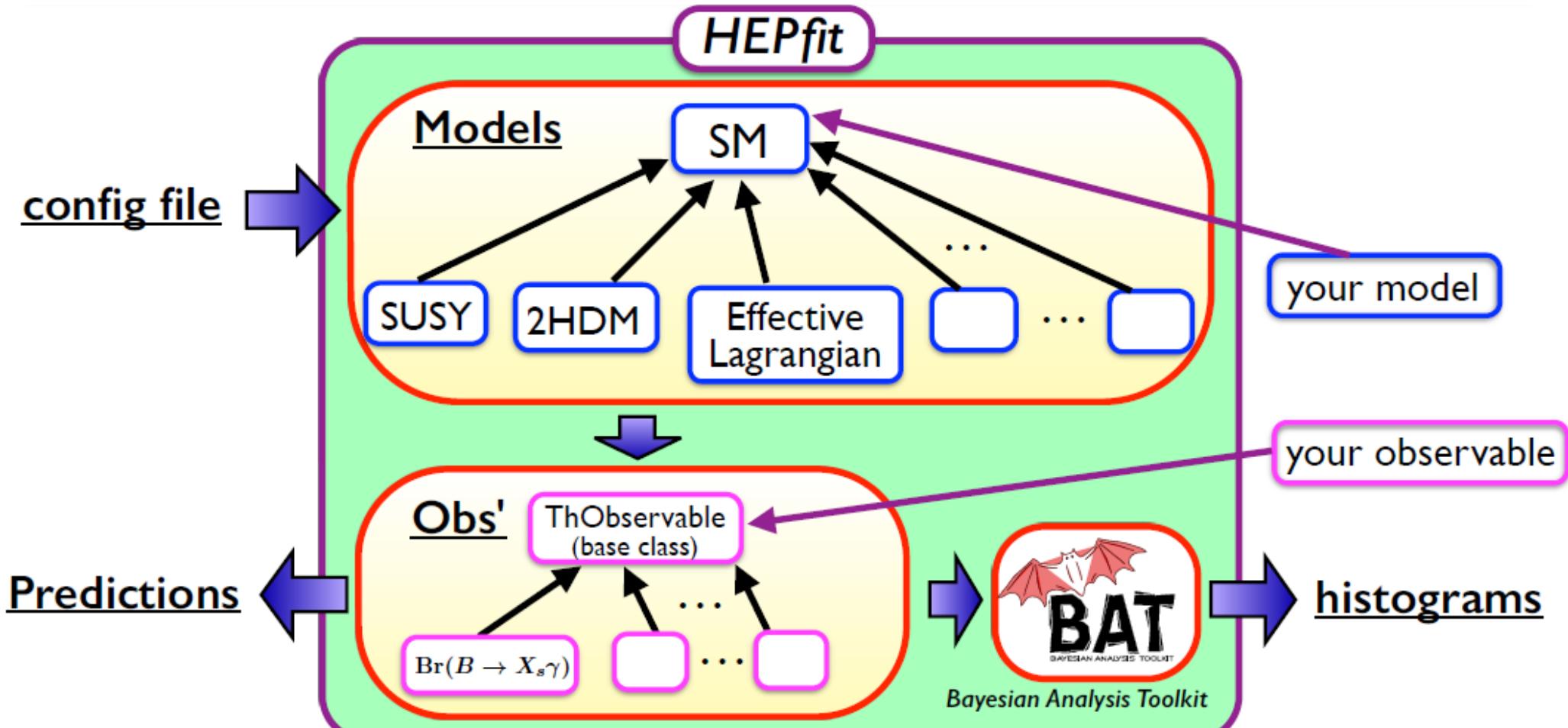


The HEPfit philosophy

- * **Open source**
 - C++ source released under GPL available on github
- * **Easy to use**
 - CMake installation, optionally include external libraries
 - configurable with few human-readable config files
 - fully doxygen-ed
- * **Highly flexible and customizable**
 - includes an expanding set of models and observables
 - provides interfaces to add new models and observables
 - can be used for quick estimate as well as global analyses
 - interfaces to the Bayesian Analysis Tools based on MCMC
- * **Fast**
 - optional support for MPI parallelization

Caldwell et al, arXiv:0808.2552
<http://mpp.mpg.de/bat>

The HEPfit structure



courtesy of S. Mishima

Installation and usage

Download HEPfit-x.x.tar.gz from the HEPfit website -
<http://hepfit.roma1.infn.it>, then

```
tar zxvf HEPfit-x.x.tar.gz  
cd HEPfit-x.x  
cmake . -DLOCAL_INSTALL_ALL=ON -DMPIBAT=ON  
make  
make install
```

this will also download, compile and install the required and optional libraries (ROOT, BAT, gsl, BOOST, MPI), if needed. Then, it is ready to run:

```
./analysis StandardModel.conf MonteCarlo.conf
```

getting in input two config files (or just one if you do not wish to do the MCMC statistical analysis)

The config file

StandardModel.conf

```
1 StandardModel
2 # Model parameters:
3 ModelParameter mtop      173.2      0.9      0.
4 ModelParameter mH1       125.6      0.3      0.
5 ...
6 CorrelatedGaussianParameters V1_lattice 2
7 ModelParameter a_0V     0.496     0.067     0.
8 ModelParameter a_1V     -2.03     0.92      0.
9 1.00     0.86
10 0.86    1.00
11
12 <All the model parameters have to be listed here>
13
14 # Observables:
15 Observable Mw          Mw          M_{W}        80.3290 80.4064 MCMC weight 80.385 0.015 0.
16 Observable GammaW      GammaW      #Gamma_{W} 2.08569 2.09249 MCMC weight 2.085 0.042 0.
17 #
18 # Correlated observables:
19 CorrelatedGaussianObservables Zpole2 7
20 Observable Alepton     Alepton     A_{1}        0.143568 0.151850 MCMC weight 0.1513 0.0021 0.
21 Observable Rbottom    Rbottom     R_{b}        0.215602 0.215958 MCMC weight 0.21629 0.00066 0.
22 Observable Rcharm    Rcharm     R_{c}        0.172143 0.172334 MCMC weight 0.1721 0.0030 0.
23 Observable AFBbottom AFBbottom  A_{FB}^{b}   0.100604 0.106484 MCMC weight 0.0992 0.0016 0.
24 Observable AFBcharm  AFBcharm  A_{FB}^{c}   0.071750 0.076305 MCMC weight 0.0707 0.0035 0.
25 Observable Abottom   Abottom   A_{b}        0.934320 0.935007 MCMC weight 0.923 0.020 0.
26 Observable Acharm   Acharm   A_{c}        0.666374 0.670015 MCMC weight 0.670 0.027 0.
27 1.00    0.00    0.00    0.00    0.09    0.05
28 0.00    1.00    -0.18   -0.10   0.07    -0.08   0.04
29 0.00    -0.18   1.00    0.04    -0.06   0.04    -0.06
30 0.00    -0.10   0.04    1.00    0.15    0.06    0.01
31 0.00    0.07    -0.06   0.15    0.15    -0.02   0.04
32 0.09    -0.08   0.04    0.06    -0.02   1.00    0.11
33 0.05    0.04    -0.06   0.01    0.04    0.11    1.00
34 #
35 # Output correlations:
36 Observable2D MwvsGammaW Mw M_{W} 80.3290 80.4064 noMCMC noweight GammaW #Gamma_{W} 2.08569 2.09249
37 ...
38 Observable2D Bd_Bsbar_mumu noMCMC noweight
39 Observable BR_Bdmumu    BR(B_{d}\rightarrow\mu\mu)    1. -1. 1.05e-10 0. 0.
40 Observable BRbar_Bsmumu BR(B_{s}\rightarrow\mu\mu)    1. -1. 3.65e-9 0. 0.
41 ...
42 Observable2D S5_P5 noMCMC noweight
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The config file

StandardModel.conf

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2 # Model parameters:
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5 ...
6 CorrelatedGaussianParameters V1_lattice 2
7 ModelParameter a_0V     0.496     0.067     0.
8 ModelParameter a_1V     -2.03     0.92      0.
9 1.00     0.86
10 0.86    1.00
11
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25 Observable Abottom   Abottom    A_{b}      0.934320 0.935007 MCMC weight 0.923 0.020 0.
26 Observable Acharm   Acharm    A_{c}      0.666374 0.670015 MCMC weight 0.670 0.027 0.
27 1.00     0.00    0.00    0.00    0.09    0.05
28 0.00     1.00    -0.18   -0.10   0.07    -0.08   0.04
29 0.00     -0.18   1.00    0.04    -0.06   0.04    -0.06
30 0.00     -0.10   0.04    1.00    0.15    0.06    0.01
31 0.00     0.07    -0.06   0.15    0.15    -0.02   0.04
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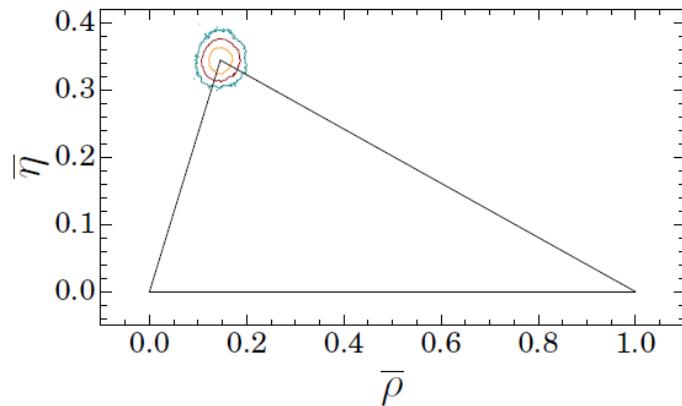
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1 StandardModel
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HEPfit model-observable matrix

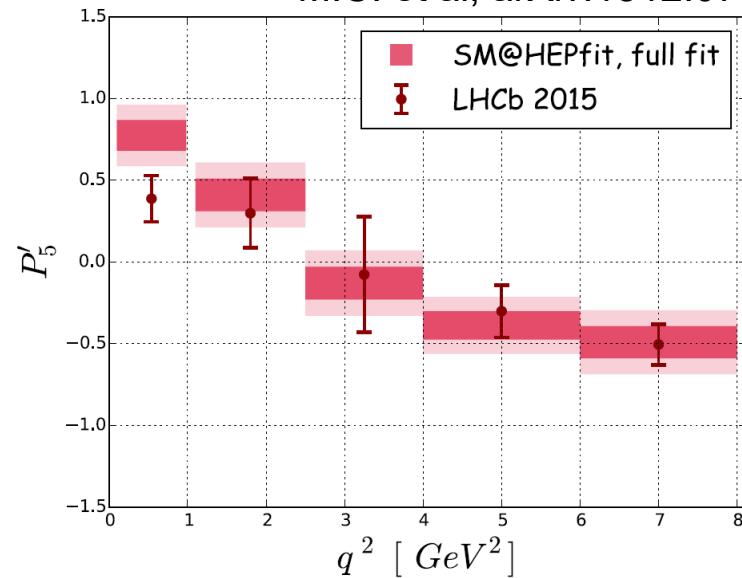
	SM	THDM	MSSM	H_{eff}		SM	THDM	MSSM	L_{EWSB}	SMEFT
$\Delta F=2$ amp's	X	O	O	O	EWPO + LEP2	X	X		X	X
$B \rightarrow \tau\nu$	X	X	O	O	$H \rightarrow VV,$ $H \rightarrow ff \mu's$	X	X	O		X
$B_{s/d} \rightarrow \mu\mu$	O	O	O	O	direct searches		X	O		
rare K decays	O			O	LFV $\ell_i \rightarrow \ell_j \gamma, 3\ell_j$				X	
$B \rightarrow X_s \gamma$	X	X	O	O	X: done					
$B \rightarrow V \gamma$ $B \rightarrow P/V \ell\ell$	X			O	O: in progress					
$B \rightarrow X_s \ell\ell$	O			O	more to come					
$B \rightarrow PP, PV$	O			O	(L-R model, aligned THDM,...)					

Here's what you get

HEPfit UTA, soon adopted by UTfit

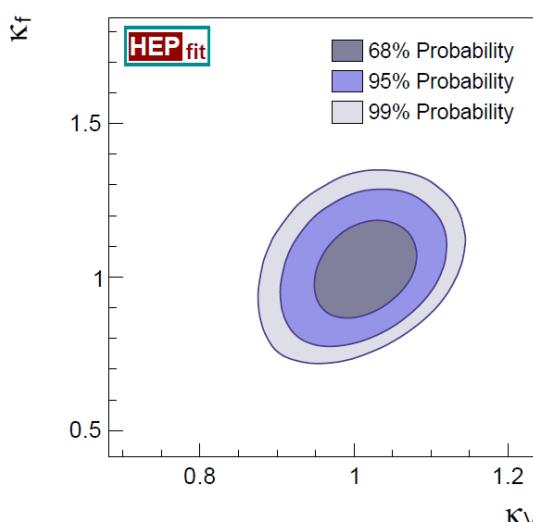


M.C. et al, arXiv:1512.07157

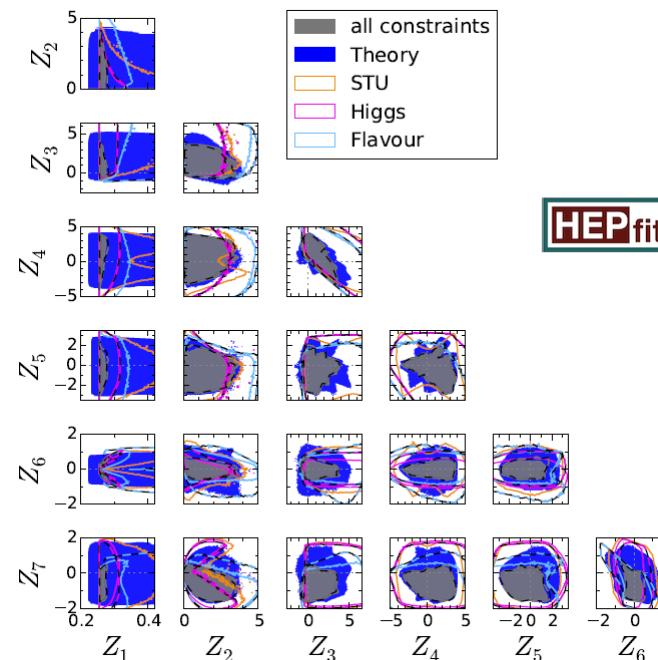


SM flavour

Fits to EWPO



De Blas et al.,
arXiv:1608.01509



2HDM with Z_2
soft symmetry
breaking

Cacchio et al.,
arXiv:1609.01290

More studies with HEPfit on the way

- * SMEFT analysis of EWPO and Higgs μ 's
- * global EFT analysis of radiative B decays
- * analysis of MSSM with generic flavour structure
- * phenomenological study of the aligned THDM
- * full-fledged NP UTA and $\Delta F=2$ EFT analysis
- * ...

What about your next phenomenological analysis?
And your next background study?

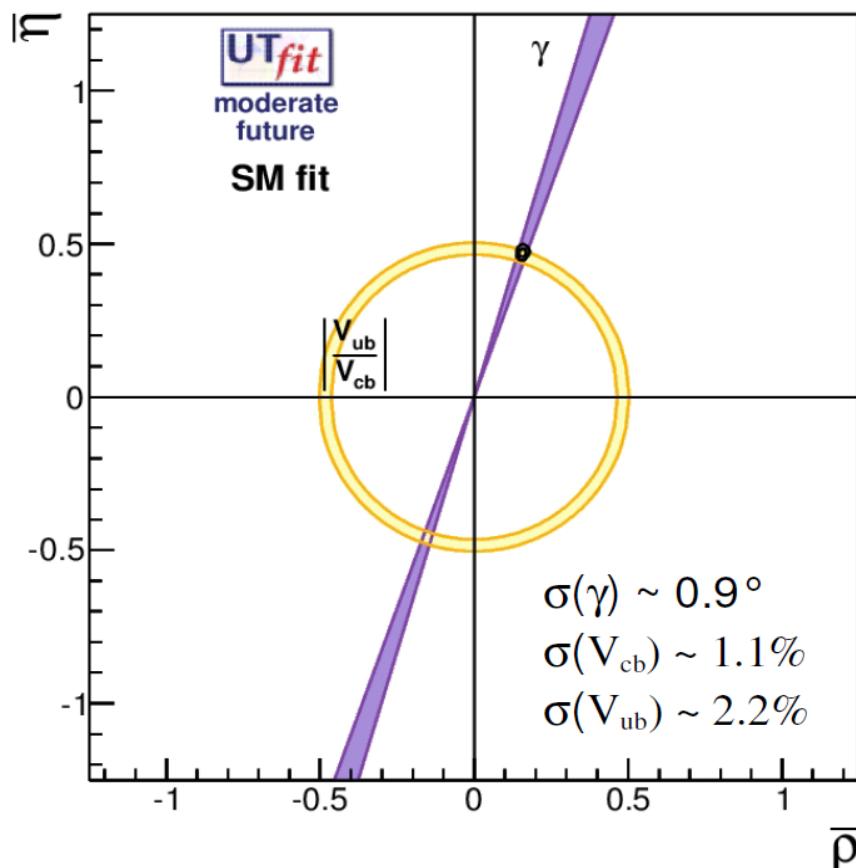
Give  a try!

<http://hepfit.roma1.infn.it>

Backup

Outlook

Look at the future



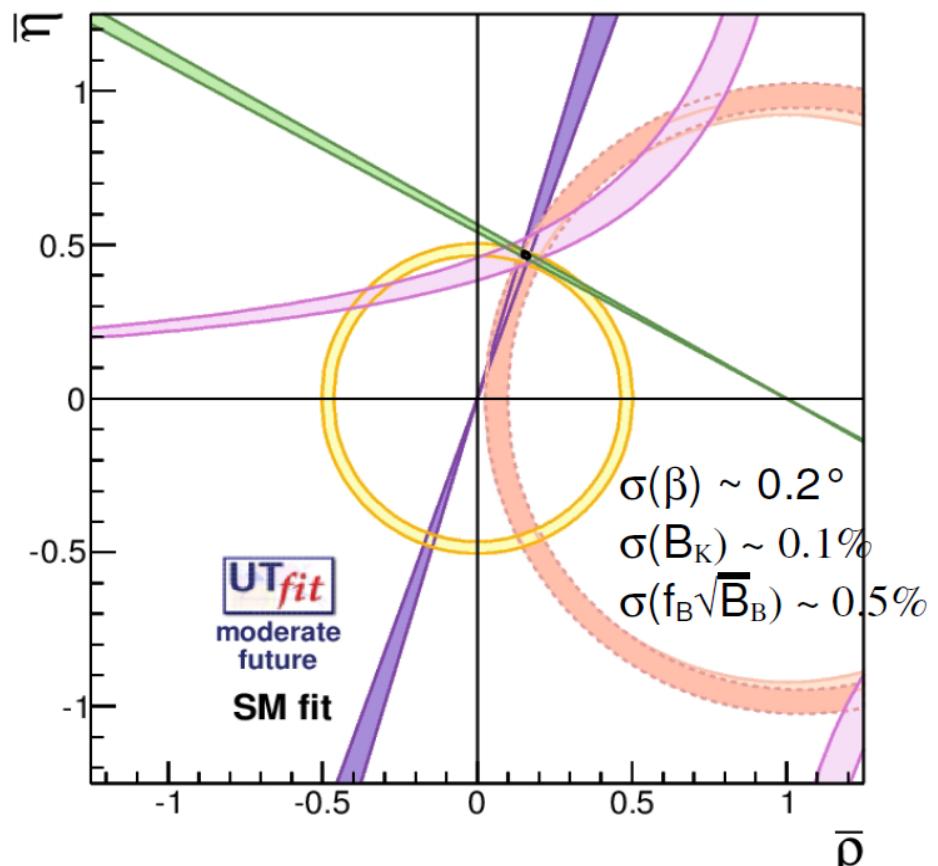
errors from tree-only fit on ρ and η :

$$\sigma(\rho) = 0.008 \text{ [currently } 0.050\text{]}$$

$$\sigma(\eta) = 0.010 \text{ [currently } 0.035\text{]}$$

errors predicted from
Belle II + LHCb upgrade

M. Bona, ICHEP 2016

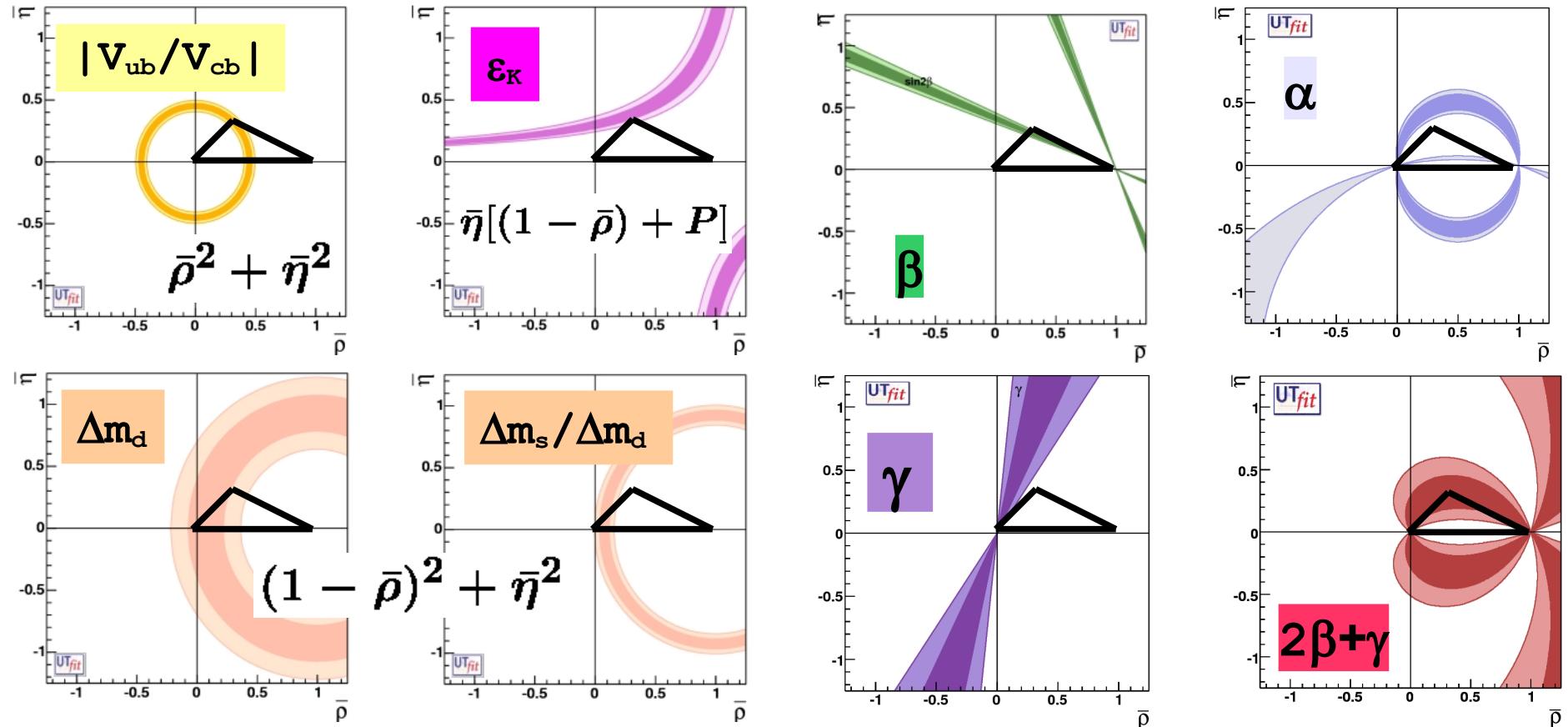


errors from 5-constraint fit on ρ and η :

$$\sigma(\rho) = 0.005 \text{ [currently } 0.015\text{]}$$

$$\sigma(\eta) = 0.004 \text{ [currently } 0.013\text{]}$$

Unitarity Triangle analysis: $V_{ub}^*V_{ud} + V_{cb}^*V_{cd} + V_{tb}^*V_{td} = 0$



Original goal:

- determine the UT apex and the CKM matrix parameters

Overconstrained fit:

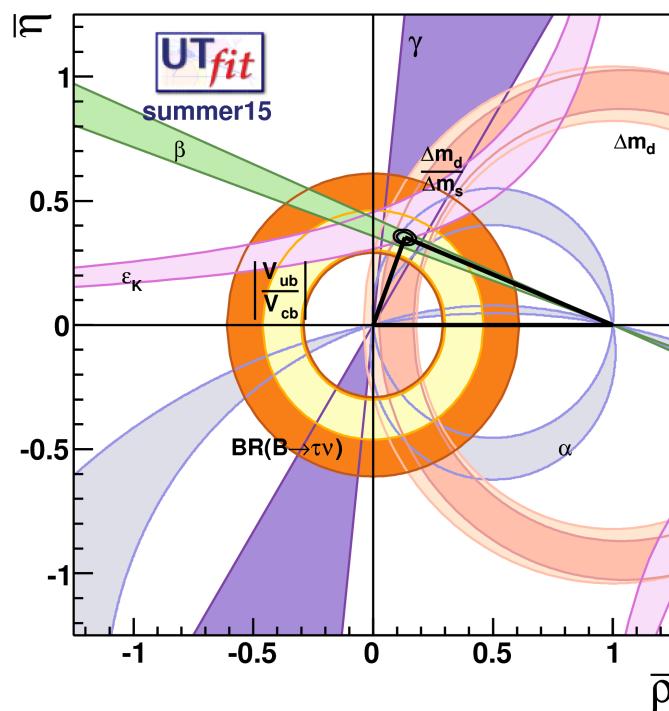
- predict observables, hadronic parameters and constrain NP

ε_K & B_K

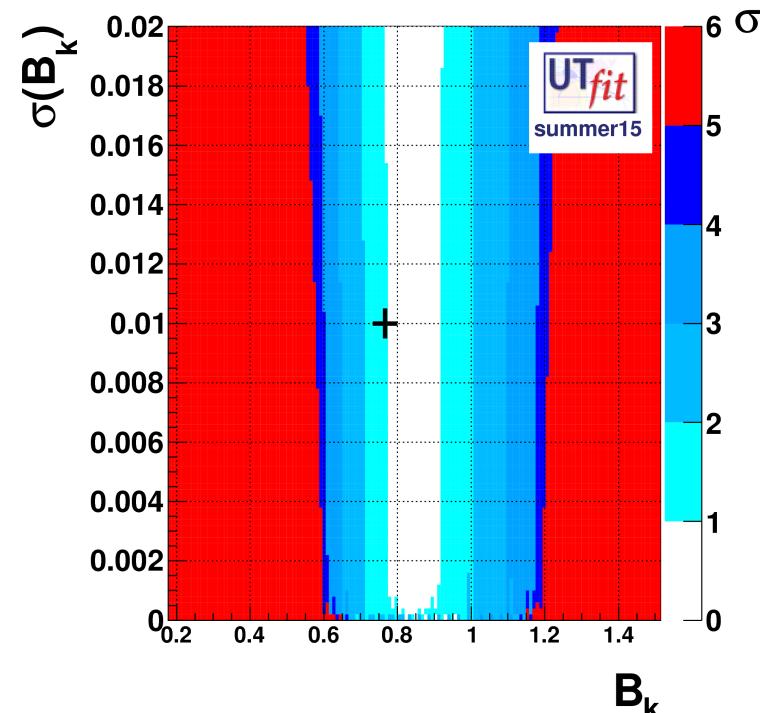
$$B_K^{\text{lattice}} = 0.766 \pm 0.010$$

$$B_K^{\text{prediction}} = 0.845 \pm 0.074$$

$\sim 1\sigma$

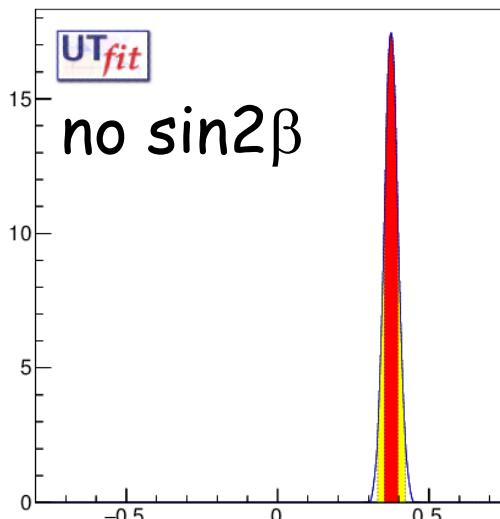


given B_K^{lattice} , ε_K calls
for large A or η

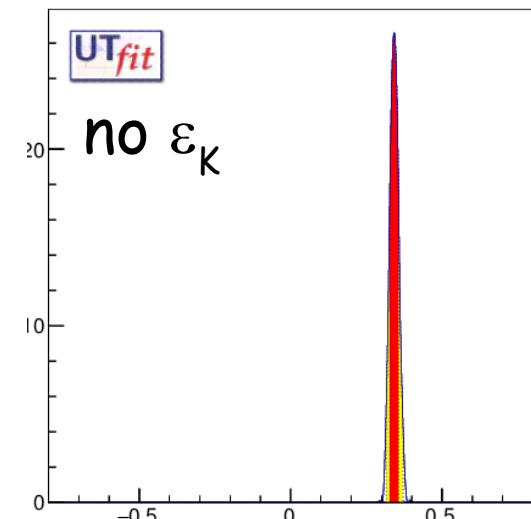


$$\bar{\eta} = 0.375 \pm 0.023$$

Probability density



$$\bar{\eta} = 0.343 \pm 0.015$$



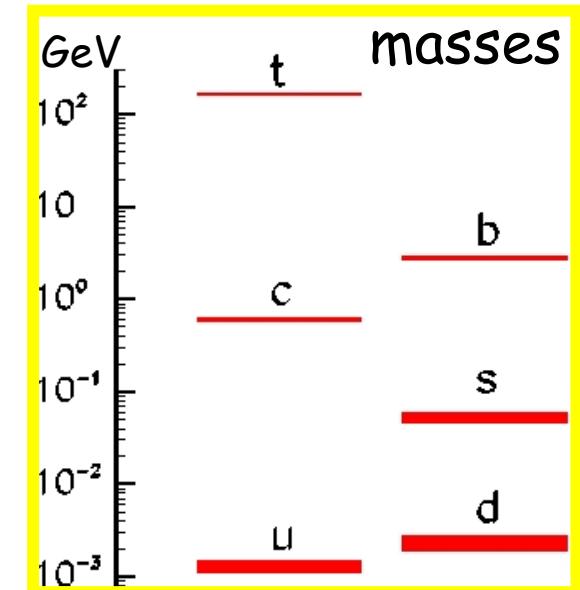
Flavour physics in the SM: rich phenomenology (FCNC suppression, mixing, CP violation, ...) but little understanding of the "why" and the "how"

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{EWSB}} + \mathcal{L}_{\text{kin}} + \mathcal{L}_{\text{gauge}} + \mathcal{L}_Y$$

The Yukawa Lagrangian describes quark flavour physics in terms of 10 physical parameters:

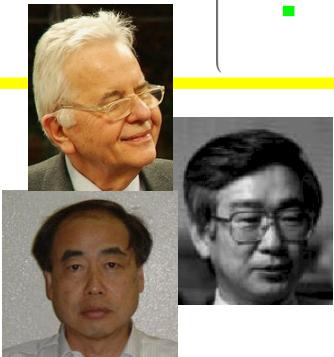
the Cabibbo-Kobayashi-Maskawa matrix

6 masses, 3 mixing angles + 1 CPV phase



mixing

$$V_{\text{CKM}} = \begin{pmatrix} \text{red} & \text{blue} & \cdot \\ \text{blue} & \text{red} & \cdot \\ \cdot & \cdot & \text{red} \end{pmatrix}$$



Beyond the SM: a powerful indirect probe of the New Physics scale Λ

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{\mathcal{L}_5}{\Lambda} + \frac{\mathcal{L}_6}{\Lambda^2} + \dots$$

has accidental
(approximate) symmetries

may violate
accidental symmetries

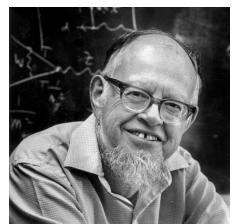
The CKM matrix in the SM

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Standard parametrization (PDG): $s_{12}, s_{13}, s_{23}, \delta$

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

Wolfenstein parametrization: λ, A, ρ, η



$$\begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

The CKM matrix in the SM

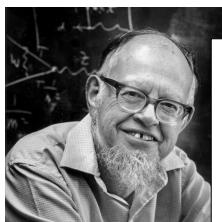
$$\begin{matrix} & d & s & b \\ u & 0.9743(2) & 0.2251(7) & 3.7(1) \cdot 10^{-3} e^{-i66(2)^\circ} \\ c & -0.2250(7) e^{i0.035(1)^\circ} & 0.9734(2) e^{-i0.0019(1)^\circ} & 4.26(7) \cdot 10^{-2} \\ t & 8.7(1) \cdot 10^{-3} e^{-i22(1)^\circ} & -4.16(6) \cdot 10^{-2} e^{i1.04(4)^\circ} & 0.99910(2) \end{matrix}$$

Standard parametrization (PDG): $s_{12}, s_{13}, s_{23}, \delta$

$$s_{12} = 0.2250 \pm 0.0007 \quad s_{23} = (4.229 \pm 0.057) \times 10^{-2}$$

$$s_{13} = (3.68 \pm 0.10) \times 10^{-3} \quad \delta = (65.9 \pm 2.0)^\circ$$

Wolfenstein parametrization: λ, A, ρ, η



$$\lambda = 0.2250 \pm 0.0007$$

$$\rho = 0.157 \pm 0.014$$

$$A = 0.833 \pm 0.012$$

$$\eta = 0.352 \pm 0.011$$

SM results

Summer 2016

SM determination of the Unitarity Triangle

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

$$R_u e^{i\gamma} + R_+ e^{-i\beta} = 1$$

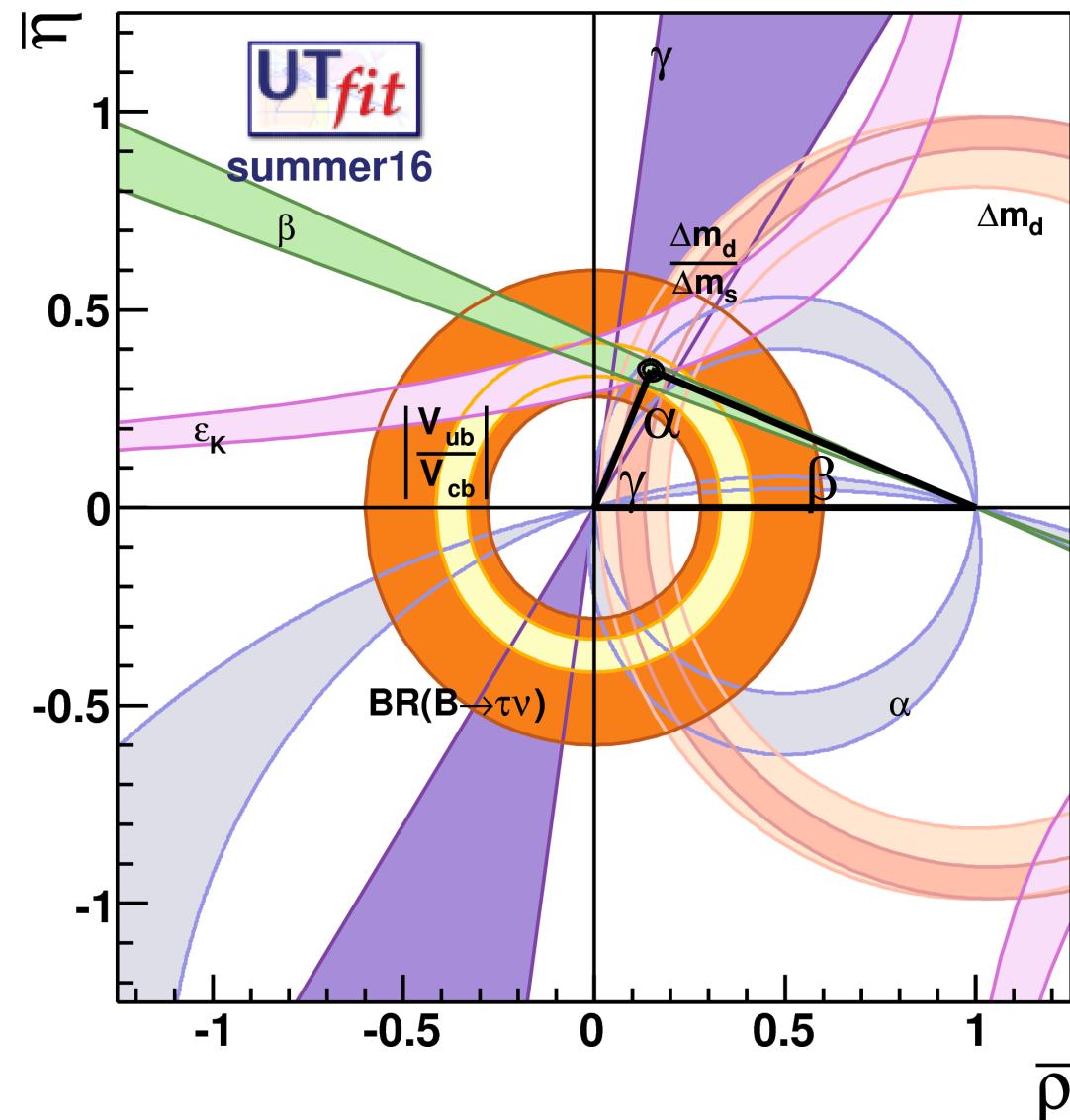
$$R_u = 0.372 \pm 0.013$$

$$R_+ = 0.917 \pm 0.022$$

$$\gamma = (65.8 \pm 1.9)^\circ$$

$$\beta = (22.11 \pm 0.76)^\circ$$

$$\alpha = (92.0 \pm 2.0)^\circ$$



apex coordinates

$$\bar{\rho} = 0.153 \pm 0.013$$

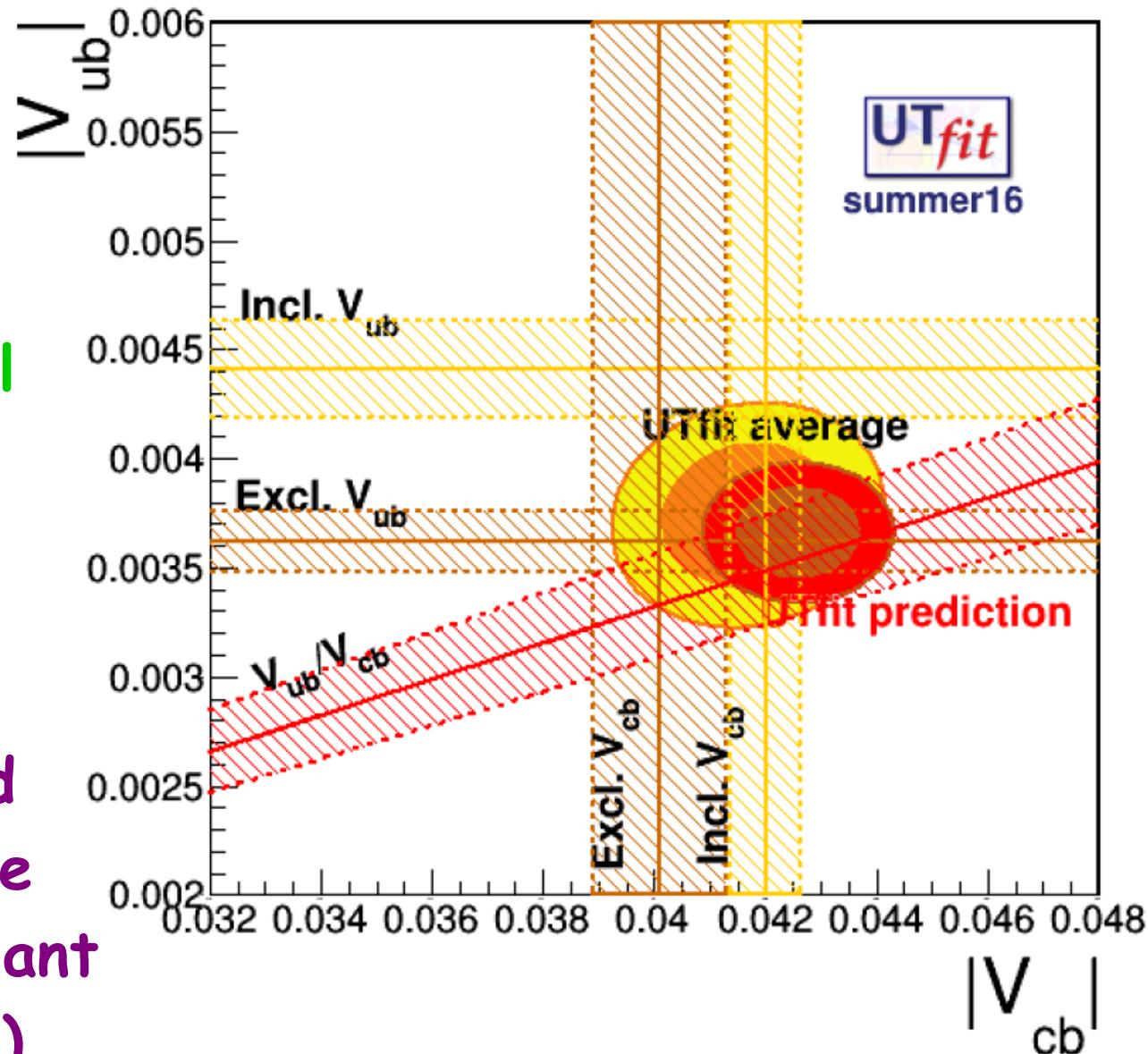
$$\bar{\eta} = 0.343 \pm 0.011$$

SM predictions: B_d & K

	Measurement	%	Prediction	Pull(σ)
$\sin 2\beta$	0.680 ± 0.023	3.4	0.724 ± 0.028	+1.2
$\gamma [^\circ]$	70.5 ± 5.7	8	65.3 ± 2.0	< 1
$\alpha [^\circ]$	94.2 ± 5.0	5	91.0 ± 2.5	< 1
$ V_{cb} \cdot 10^3$	41.7 ± 1.0	2	42.6 ± 0.7	< 1
$ V_{cb} _{\text{incl}} \cdot 10^3$	42.0 ± 0.6	1	" "	< 1
$ V_{cb} _{\text{excl}} \cdot 10^3$	40.0 ± 1.1	3	" "	+1.7
$ V_{ub} \cdot 10^3$	3.73 ± 0.21	6	3.66 ± 0.12	< 1
$ V_{ub} _{\text{incl}} \cdot 10^3$	4.40 ± 0.22	5	" "	-3.0
$ V_{ub} _{\text{excl}} \cdot 10^3$	3.61 ± 0.13	4	" "	< 1
$\varepsilon_K \cdot 10^3$	2.228 ± 0.011	0.5	2.03 ± 0.18	-1.1
$\text{BR}(B \rightarrow \tau\nu) \cdot 10^{-4}$	1.06 ± 0.20	20	0.81 ± 0.06	-1.3

V_{cb} & V_{ub}

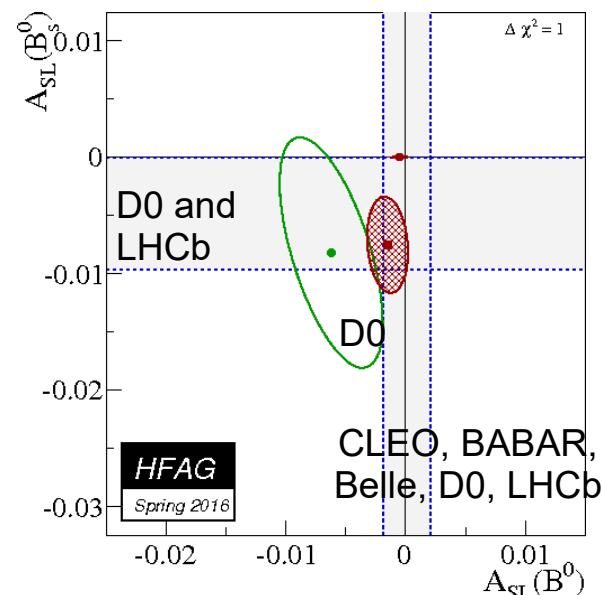
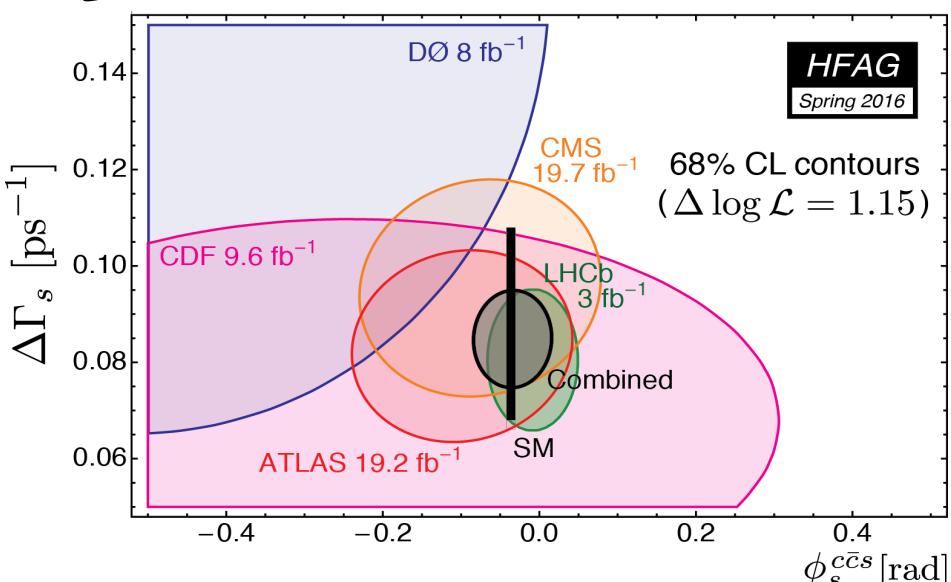
- * data favour inclusive V_{cb} and exclusive V_{ub}
- * no clear evidence to inculpate a theoretical method (a suspect: V_{ub} inclusive)
- * difficult to explain with NP (right-handed currents are not viable in a $SU(2) \times U(1)$ invariant effective field theory)



Crivellin, Pokorski, arXiv:1407.1320

SM predictions: B_s

Measurement	%	Prediction	Pull (σ)
$\Delta m_s [\text{ps}^{-1}]$	17.757 ± 0.021	0.1	17.7 ± 0.9
$\beta_s [^\circ]$	0.94 ± 0.94	100	1.04 ± 0.03
$\Delta \Gamma_s / \Gamma_s$	0.124 ± 0.009	7.2	0.154 ± 0.012
$A_{SL}^s \cdot 10^4$	-75 ± 41	57	0.13 ± 0.01



Deviations from the SM to keep an eye on

- ▶ ε'/ε
- ▶ $\text{BR}(B_s \rightarrow \mu\mu), \text{BR}(B \rightarrow \mu\mu)$
- ▶ $R(D), R(D^*)$
- ▶ R_K
- ▶ q^2 spectrum of $B \rightarrow K^* \mu\mu$

Deviations from the SM to keep an eye on

► ϵ'/ϵ

► $\text{BR}(B_s \rightarrow \mu\mu)$, $\text{BR}(B \rightarrow \mu\mu)$

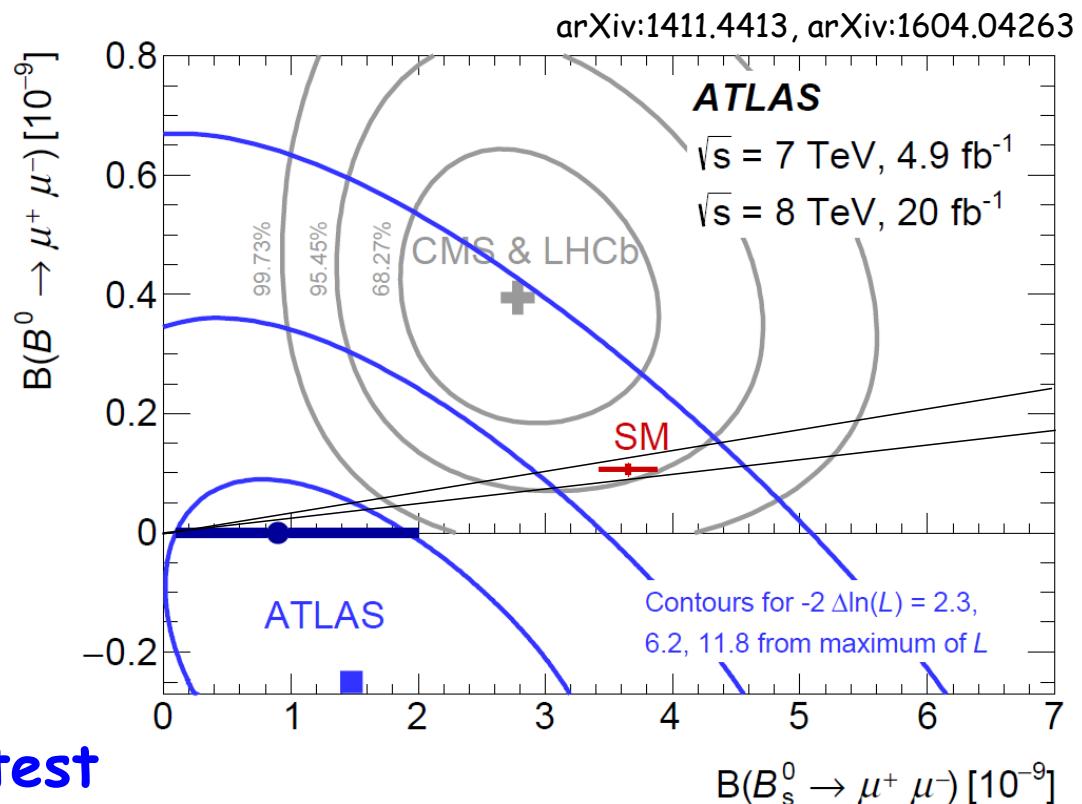
$$\text{BR}(B_s \rightarrow \mu^+ \mu^-) = (2.8^{+0.7}_{-0.6}) \times 10^{-9}$$

$$\text{BR}(B_s \rightarrow \mu^+ \mu^-)_{SM} = (3.65 \pm 0.23) \times 10^{-9} + 1.2\sigma$$

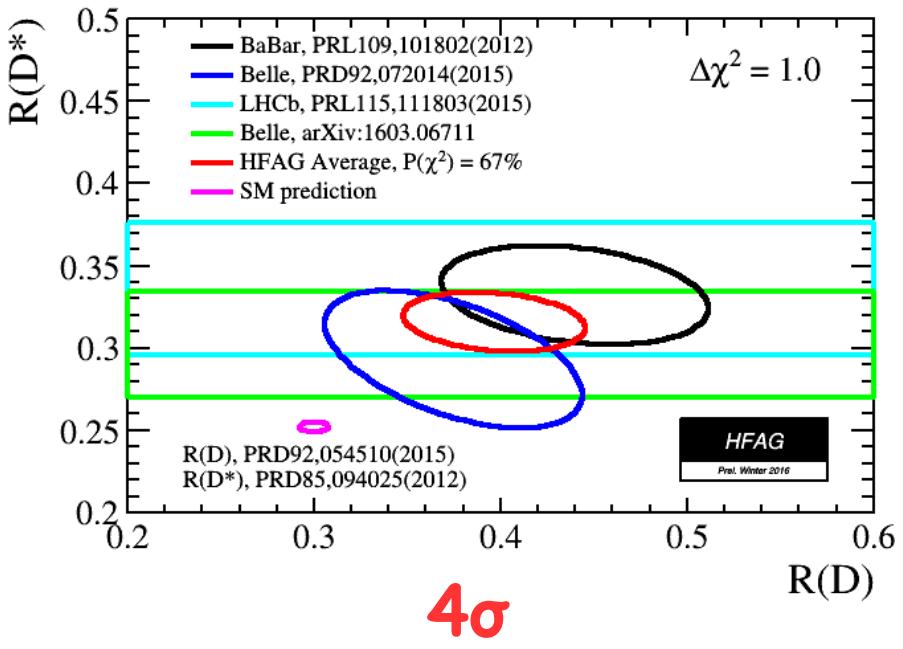
$$\text{BR}(B \rightarrow \mu^+ \mu^-) = (3.9^{+1.6}_{-1.4}) \times 10^{-10}$$

$$\text{BR}(B \rightarrow \mu^+ \mu^-)_{SM} = (1.06 \pm 0.09) \times 10^{-10} - 2.2\sigma$$

SM predictions from
Bobeth et al., arXiv:1311.0903



Minimal Flavour Violation test



$$R(X) = \frac{\Gamma(B \rightarrow X\tau\nu)}{\Gamma(B \rightarrow X\ell\nu)}$$

$$R(D) = -1.9\sigma$$

$$0.397 \pm 0.040 \pm 0.028$$

$$R(D)_{\text{SM}} = 0.300 \pm 0.008$$

Na et al., arXiv:1505.03925

$$R(D^*) = -3.3\sigma$$

$$0.316 \pm 0.016 \pm 0.010$$

$$R(D^*)_{\text{SM}} = 0.252 \pm 0.003$$

Fajfer et al., arXiv:1203.2654

an eye on

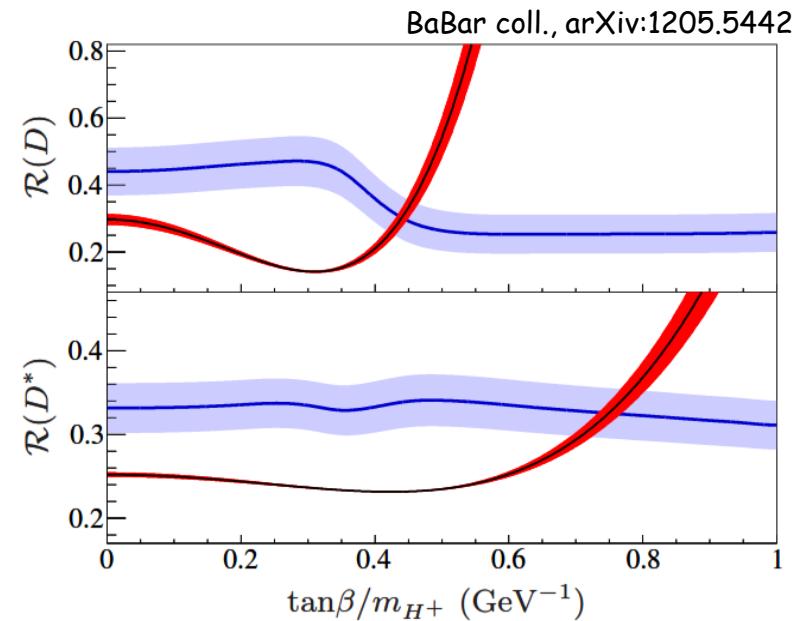
► $R(D)$, $R(D^*)$

► $\Gamma(B)$

simplest realizations of
2HDM cannot explain the
excess in the two
channels simultaneously

► q^2

more exotic NP required,
e.g. 2HDM-type III,
leptoquarks,
compositeness, ...



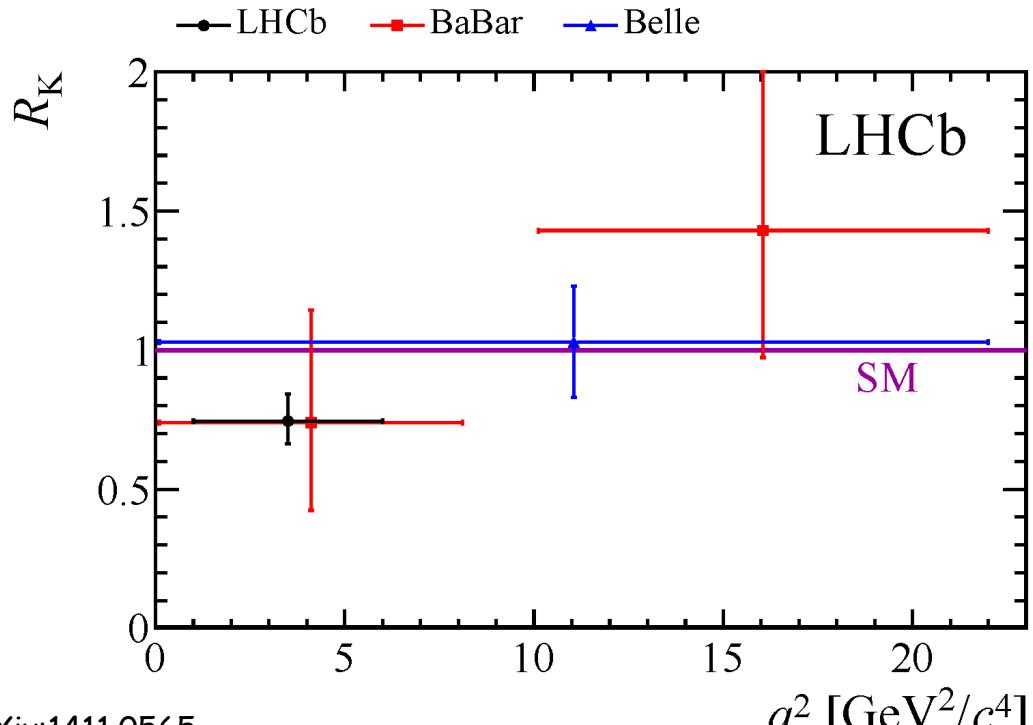
Deviations

Also beyond the SM,
such a large violation
of lepton universality
is not easily obtained
(e.g. leptoquarks)

it may be correlated
to large LFV in B
decays

$$b \rightarrow s \ell_i^+ \ell_j^-$$

Glashow et al., arXiv:1411.0565



► **R_K**

$$R_K = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ e^+ e^-]}{dq^2} dq^2} = 0.745^{+0.090}_{-0.074} (\text{stat}) \pm 0.036 (\text{syst})$$

LHCb Collaboration, arXiv:1406.6482

$R_K^{\text{SM}} = 1.0003 \pm 0.0001$

Bobeth et al., arXiv:0709.4174

+2.6 σ

If LFU is induced by NP coupled to the 3rd generation at a scale much larger than the weak scale, in the effective theory, one gets

$$Q'_{L,3}\gamma_\mu Q'_{L,3}L'_L\gamma^\mu L'_L, \quad Q'_{L,3}\gamma_\mu\sigma^i Q'_{L,3}L'_L\gamma^\mu\sigma^i L'_L,$$

i) can account for the anomalies in R_K , $R(D)/R(D^*)$

ii) give typically rise to large LFV

Glashow et al., arXiv:1411.0565

Bhattacharya et al., arXiv:1412.7164

iii) running effects produce large corrections to the $V_{l\ell}$ vertices and induce purely leptonic LFV transitions

Feruglio et al., arXiv:1606.00524

Explicit models face more severe constraints, still there are viable proposals:

* non-minimal scalar leptoquarks

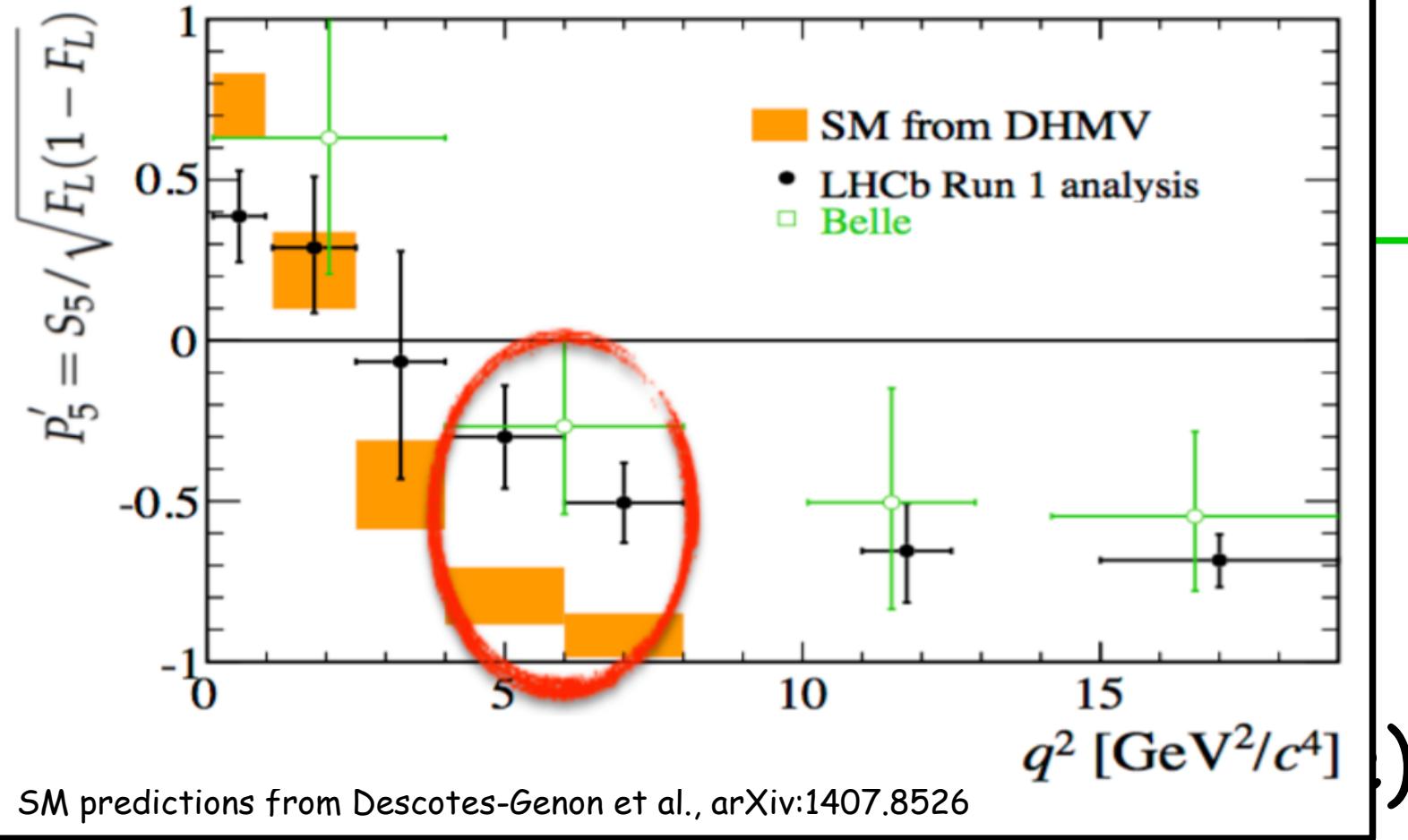
Becirevic et al., arXiv:1608.08501

* triplet heavy gauge bosons

Greljo et al., arXiv:1506.01705

*

...



► q^2 spectrum of $B \rightarrow K^* \mu\mu$

$$\frac{d^{(4)}\Gamma}{dq^2 d(\cos\theta_l)d(\cos\theta_k)d\phi} = \frac{9}{32\pi} \left(I_1^s \sin^2\theta_k + I_1^c \cos^2\theta_k + (I_2^s \sin^2\theta_k + I_2^c \cos^2\theta_k) \cos 2\theta_l \right. \\ \left. + I_3 \sin^2\theta_k \sin^2\theta_l \cos 2\phi + I_4 \sin 2\theta_k \sin 2\theta_l \cos\phi \right. \\ \left. + I_5 \sin 2\theta_k \sin\theta_l \cos\phi + (I_6^s \sin^2\theta_k + I_6^c \cos^2\theta_K) \cos\theta_l \right. \\ \left. + I_7 \sin 2\theta_k \sin\theta_l \sin\phi + I_8 \sin 2\theta_k \sin 2\theta_l \sin\phi \right. \\ \left. + I_9 \sin^2\theta_k \sin^2\theta_l \sin 2\phi \right)$$

angular
analysis

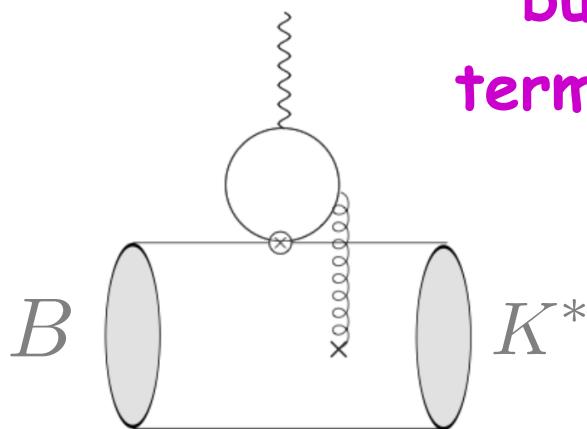
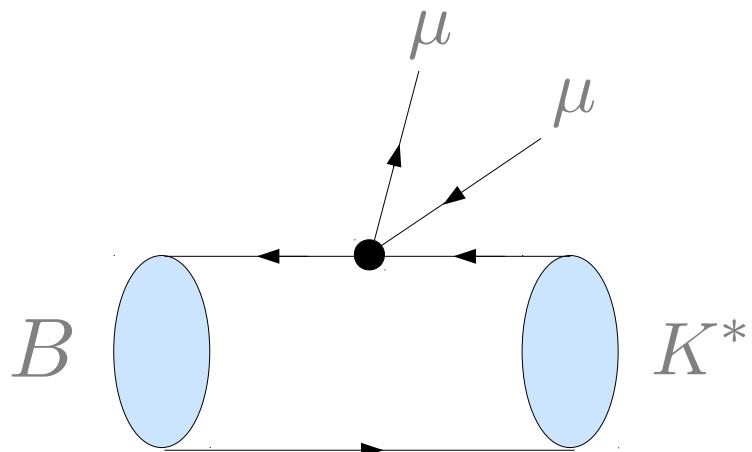
$$S_i = \left(I_i^{(s,c)} + \bar{I}_i^{(s,c)} \right) / \Gamma' \\ (2\Gamma' \equiv d\Gamma/dq^2 + d\bar{\Gamma}/dq^2)$$

8 CP-AVERAGED OBSERVABLES

$$F_L, A_{FB}, S_{3,4,5,7,8,9}$$

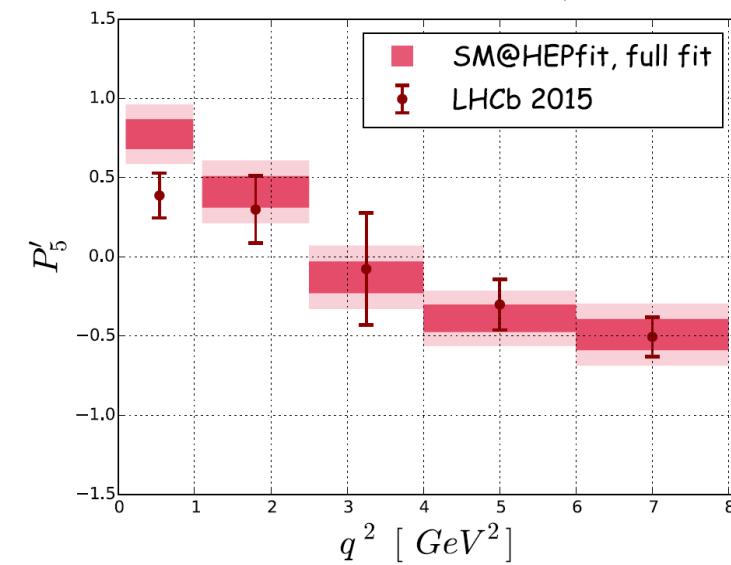
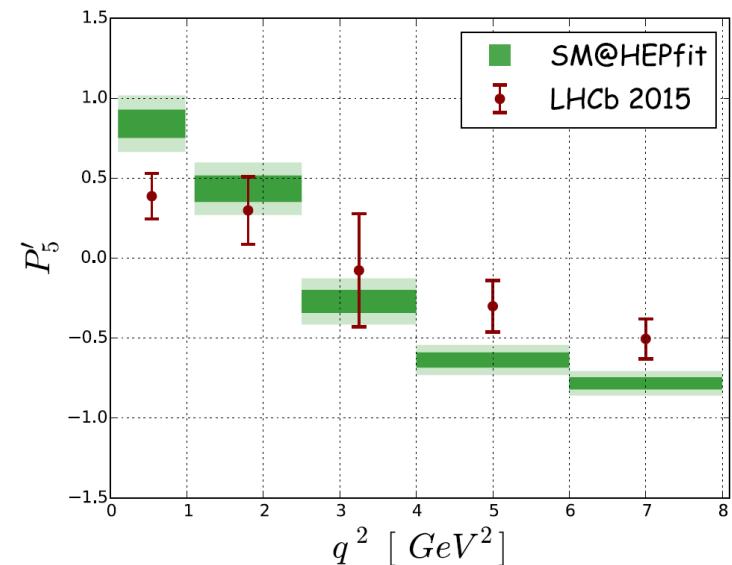
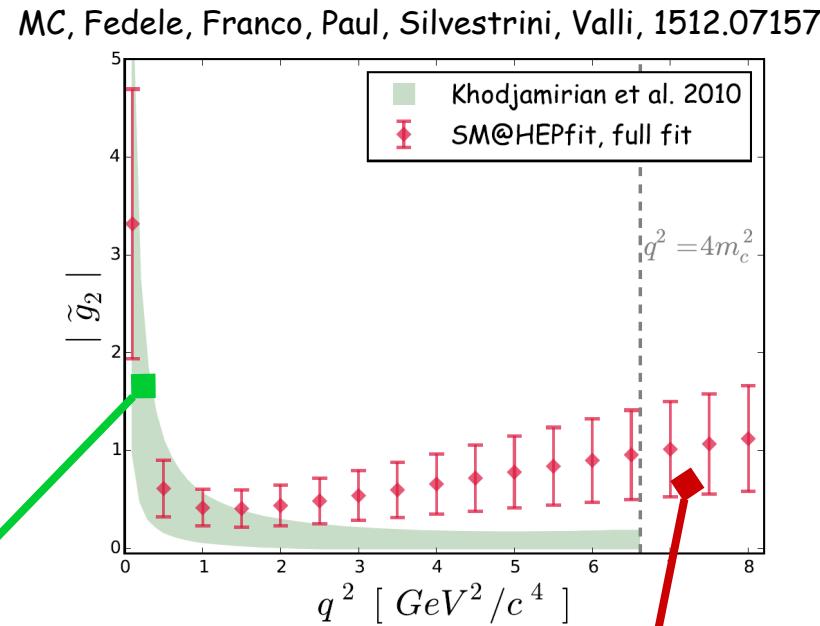
$B \rightarrow K^* \ell^+ \ell^-$

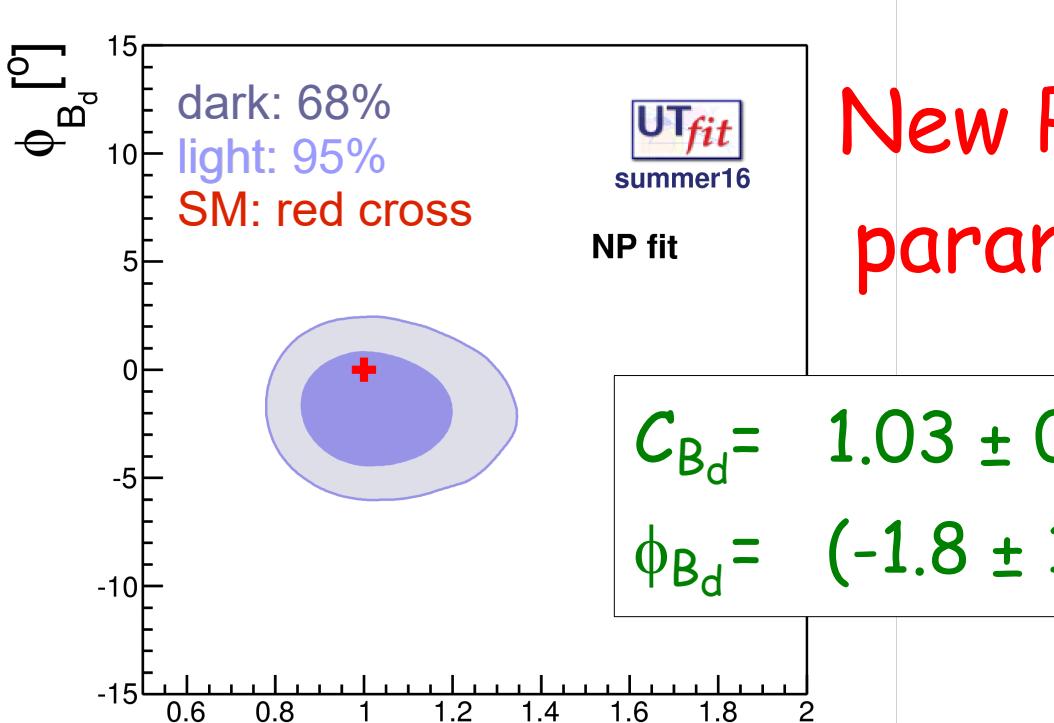
The leading amplitude is factorizable



BSM sensitivity could be hindered by hadronic uncertainties

but non-factorizable terms may be important:





New Physics parameters

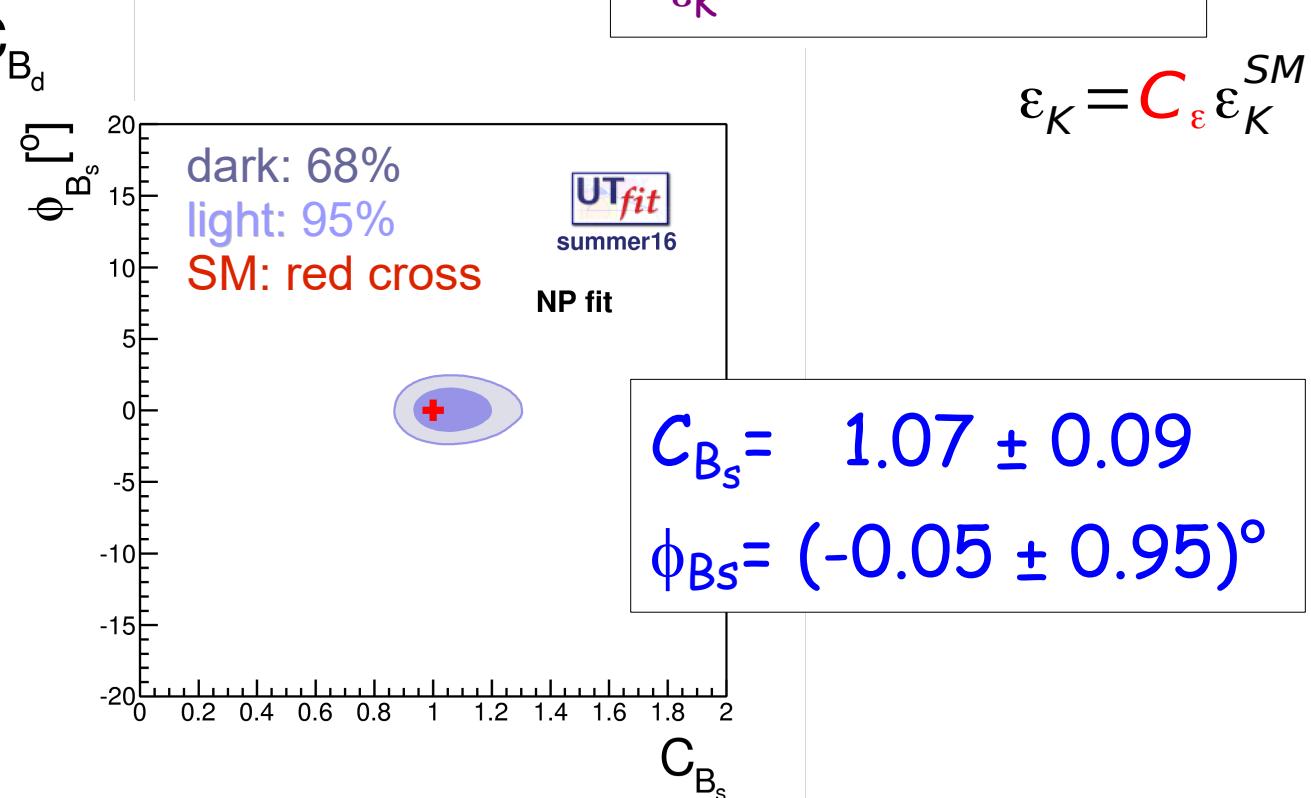
$$\Delta m_q = C_{B_q} (\Delta m_q)^{SM}$$

$$a_{CP}^{B_d \rightarrow J/\psi K_s} \rightarrow \sin 2(\beta + \phi_{B_d})$$

$$a_{CP}^{B_s \rightarrow J/\psi \phi} \rightarrow -\beta_s + \phi_{B_s}$$

$$A_{SL}^q = \text{Im} \left(\Gamma_{12}^q / A_q \right)$$

$$\Delta \Gamma^q / \Delta m_q = \text{Re} \left(\Gamma_{12}^q / A_q \right)$$



sample user codes for MCMC

```
#include <iostream>
#include <HEPfit.h> ← the header

#ifndef _MPI
#include <mpi.h>
#endif

int main(int argc, char** argv)
{

#ifndef _MPI
    MPI::Init();
    int rank = MPI::COMM_WORLD.Get_rank();
    MPI::Status status;
#else
    int rank = 0;
#endif

    try {
        if(argc != 3){
            if (rank == 0) std::cout << "usage: " << argv[0] << " ModelConf.conf MonteCarlo.conf\n" << std::endl;
            return EXIT_SUCCESS;
        }
        std::string ModelConf = argv[1];
        std::string MCMCConf = argv[2];
        std::string FileOut = "";
        std::string JobTag = "";

        ThObsFactory ThObsF;
        ModelFactory ModelF;

        MonteCarlo MC(ModelF, ThObsF, ModelConf, MCMCConf, FileOut, JobTag);
        MC.Run(rank);
    } ← the user code
    }

#ifndef _MPI
    MPI::Finalize();
#endif

    return EXIT_SUCCESS;
} catch (const std::runtime_error& e) {
    std::cerr << e.what() << std::endl;
    return EXIT_FAILURE;
}
}
```

to implement your own statistical analysis

```
#include <iostream>
#include <ComputeObservables.h>

int main(int argc, char** argv)
{
    try {
        std::string ModelConf = argv[1];
        std::map<std::string, double> DPars;

        ThObsFactory ThObsF;
        ModelFactory ModelF;

        ComputeObservables CO(ModelF, ThObsF, ModelConf);
        { CO.AddObservable("Mw");
        CO.AddObservable("GammaZ");
        CO.AddObservable("AFBbottom"); } ← list observables
        std::map<std::string, double> D0bs = CO.getObservables();

        for (int i = 0; i < 2; i++) {
            DPars["Mz"] = 91.1875 + 0.0001 * i;
            DPars["AlsMz"] = 0.1184 + 0.000001 * i;

            D0bs = CO.compute(DPars);

            std::cout << "\nParameters[" << i + 1 << "]:\n" << std::endl;
            for (std::map<std::string, double>::iterator it = DPars.begin(); it != DPars.end(); it++) {
                std::cout << it->first << " = " << it->second << std::endl;
            }
            std::cout << "\nObservables[" << i + 1 << "]:\n" << std::endl;
            for (std::map<std::string, double>::iterator it = D0bs.begin(); it != D0bs.end(); it++) {
                std::cout << it->first << " = " << it->second << std::endl;
            }
        }

        return EXIT_SUCCESS;
    } catch (const std::runtime_error& e) {
        std::cerr << e.what() << std::endl;
        return EXIT_FAILURE;
    }
}
```

list observables

your own statistical analysis

```
## Number of chains
NChains                      96
#
## Max iterations in prerun
PrerunMaxIter                 100000
#
## Analysis iterations
Iterations                     1000000
#
## Write Markov Chain
WriteChain                     false
#
## Use a particular seed
#Seed                          0
#
## Find mode with Minuit
FindModeWithMinuit            false
#
## Calculate the evidence (total normalization)
CalculateNormalization        false
#
## Print all marginalized plots
PrintAllMarginalized           true
#
## Print correlation matrix
PrintCorrelationMatrix         false
#
## Print knowledge update plots
PrintKnowledgeUpdatePlots      false
#
## Print parameter plots
PrintParameterPlot             false
#
## Use ordering of parameters in the MonteCarlo run
OrderParameters                false
```