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# Full 4D angular analysis of $B \rightarrow D^*Iv$ at BaBar

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On behalf of the BaBar collaboration

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### $B \rightarrow D^* I_V$ with $D^* \rightarrow D\pi$ decays

- B→D<sup>\*</sup>ℓν (D<sup>\*</sup> → Dπ) decay has a 4body topology
  - $q^2 = (p_B p_{D^*})^2$
  - 3 angles:  $\Omega = \{\Theta_{\ell}, \Theta_{\vee}, \chi\}$
  - The spin of D\* retains information on the spin of the recoiling W\*
    - Rich phenomenology
    - Sensitive to New Physics

 $B \rightarrow D^*W^*, W^* \rightarrow \ell v$ 



• The differential decay rate is

$$rac{d\Gamma}{dq^2 d\Omega} \propto \sum_{i=1}^{14} f_i(\Omega) \Gamma_i(q^2)$$

- Functions  $f_i(\Omega)$  depend only on angles  $\Omega = \{\theta_{i}, \theta_{i}, \chi\}$
- $\Gamma_i(q^2)$  are functions of the helicity amplitudes  $H_{+,-,0}$ 
  - Can be expressed in terms of axial and vector form factors {A<sub>1</sub>, A<sub>2</sub>, V}

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### Data sample: the hadronic tagging

- Analysis based on 426 fb<sup>-1</sup> at Y(4S)
- Hadronic tagging
  - Suppress continuum e<sup>+</sup>e<sup>-</sup> → qq and combinatorial background
  - Improve the resolution on the kinematics of the signal decay
    - Boost kinematics in the B<sub>sig</sub> rest frame
    - Increase the signal/background separation
- Improved B<sub>tag</sub> algorithm used also in
  - B→D(\*)ππℓν PRL 116 (2016) 041801
  - R(D)-R(D\*) PRL 109 (2012) 101802
- 2968 modes, different seeds considered (D<sup>0</sup>, D<sup>+</sup>, D<sub>s</sub>, J/ψ) and looser cuts on intermediate states



 But no requirements on purity of the tag side: sample is very clean

### Data sample: the hadronic tagging

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 Most precise previous measurements to the data are from BaBar and Belle



Phys.Rev.D77:032002(2008)



- Untagged samples: high efficiency but higher background and unconstrained kinematics
- Fit the projections to the 4-dimensions  $q^2$ , $\theta_{\chi}$ , $\theta_{\chi}$ , $\chi$ 
  - Strong statistical correlation between the bins in the various projections need to be considered
  - Reduced sensitivity to form factor shapes

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### Selection

- Full exclusive event topology reconstructed: no additional tracks
  - $B_{tag}^{0}$  &  $B^{0} \rightarrow D^{*-} \ell^{+} v$ , with  $D^{*-} \rightarrow D^{0} \pi^{-}$ ,  $\ell = e, \mu$
  - $B_{tag}^{-}$  & B<sup>+</sup> $\rightarrow$ D<sup>\*0</sup> $\ell^+v$ , with D<sup>\*0</sup> $\rightarrow$  D<sup>0</sup> $\pi^0$ ,  $\ell$ =e,µ
    - D<sup>0</sup> reconstructed in the cleanest mode:  $K^-\pi^+$ ,  $K^-\pi^+\pi^0$ ,  $K^-\pi^+\pi^-\pi^+$
  - Positive Particle identification for all particles
- Minimal selection:
  - $|\mathbf{p}_{\pi,lab}| < 400 \text{ MeV}, |\mathbf{p}_{e,lab}| > 200 \text{ MeV} + \text{brem. recovery, } |\mathbf{p}_{e,lab}| > 300 \text{ MeV}$
  - $\Delta m = m(D\pi)-m(D)$  consistent with PDG at  $4\sigma$

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  - $\Delta m = m(D\pi)-m(D)$  consistent with PDG at  $4\sigma$
- Kinematic fit of the full event topology:  $e^+e^- \rightarrow Y(4S) \rightarrow B_{tag} \& B \rightarrow D^*\ell v$ 
  - Mass constraint: B<sub>tag</sub>, B<sub>sig</sub>, D, D\*,v
  - Vertex constraint: beam spot, secondary vertexes
  - Probability of the  $\chi^2$  of the kinematic fit used as discriminating variable
  - Event further cleaned requiring  $E_{extra} = \Sigma E_{\gamma} < 0.4-0.6$  GeV (depending on the mode)

### The discriminating variable U

Signal variables computed in B signal variables computed in B

$$U = E_{\rm miss} - |\vec{p}_{\rm miss}| = E_{\nu} - |\vec{p}_{\nu}|$$

From kinematic fit without v-mass constraint

- Small background ~ 2%
- Good data-MC agreement in all variables



### Fit of the Form Factors

- Signal is very clean: background modeled from simulation
- <u>Unbinned ML fit</u> using the full 4D differential rate

 $\frac{d\Gamma}{dq^2d\Omega}$ 

- All selected events enter
  - Background contribution described in the ML
- Signal acceptance in the full 4-D phase space obtained from a large sample of signal events generated flat in dq<sup>2</sup>dΩ

- External parameters are used to obtain |V<sub>cb</sub>| from the FF's shape given by the angular variables, added as gaussian constraints in the likelihood
  - Rely on BF(B→D\*ℓv) and lifetimes from HFLAV
  - h<sub>A1</sub>(w=1)=0.906 ± 0.013 FNAL/MILC

$$\Gamma_{tot} \equiv \int \frac{d\Gamma}{dq^2 d\Omega} dq^2 d\Omega = \frac{BF(B \to D^* \ell \nu)}{\tau_B}$$

### CLN fit results: |V<sub>cb</sub>|

• Using CLN parameterization



 $1.122 \pm 0.024$  $1.270 \pm 0.026$  $0.852 \pm 0.018$  $38.76 \pm 0.042 \pm 0.055$ HFLAV Spring 2019 (BaBar'19 not-included)

- Result includes systematics
- Many cross checks performed
  - Analysis separated in B<sup>0</sup> and B<sup>+</sup>:  $\pi^0$  and  $\pi^+$  slow
  - Separation in lepton mode and D decay modes
  - Effect of resolution on kinematic variables ( $\sigma(q^2)=0.072 \text{ GeV}^2$ ) negligible
  - Dominant contribution to systematics is from the residual background

### BGL fit results: |V<sub>cb</sub>|

- In the BGL fit, the z-expansion is truncated at N=1
  - With the present sensitivity, N=2 terms statistical insignificant and mostly violates unitarity
  - $|V_{cb}|$  + 5 coefficients:  $\mathbf{a}_0^f$ ,  $\mathbf{a}_1^f$ ,  $\mathbf{a}_1^{F1}$ ,  $\mathbf{a}_0^g$ ,  $\mathbf{a}_1^g$
- Two relations used to connect the coefficients

$$\begin{cases} f(q_{max}^2) = 2\sqrt{m_B m_D^*} \cdot h_{A_1}(q_{max}^2) & \rightarrow a_0^{\mathsf{f}} \text{ is constrained by } \mathsf{h}_{\mathsf{A}_1} \\ F_1(q_{max}^2) = (m_B - m_{D^*})f(q_{max}^2) & \rightarrow a_0^{\mathsf{F}_1} \text{ is not independent} \end{cases}$$

$a_0^f \times 10^2$	$a_1^f \times 10^2$	$a_1^{F_1} \times 10^2$	$a_0^g \times 10^2$	$a_1^g \times 10^2$	$ V_{cb}  \times 10^3$
1.29	1.63	0.03	2.74	8.33	38.36
$\pm 0.03$	$\pm 1.00$	$\pm 0.11$	$\pm 0.11$	$\pm 6.67$	$\pm 0.90$

•  $|V_{cb}|$  consistent with BaBar-CLN fit and CLN-WA

 $|V_{cb}|$ 

- Both CLN and BGL values for  $|V_{cb}|$  are consistent with the World Average



• Still in tension with inclusive  $|V_{cb}|$ 

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#### BGL results: form factors shapes

LCSR'08: Faller et al EPJC60 (2009) 603



W

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#### M. Rotondo

W

### Impact on R(D\*) prediction

- The BGL form factors lead to an updated prediction for R(D\*)
  - R(D\*)=0.253 ± 0.005
  - The uncertainty on the additional form factors needed for massive leptons follows the assumptions in Gambino'17 JHEP11 (2017) 061 (dominant source of uncertainty)

	R(D)	R(D*)	RD-RD* #σ from SM	RD* only #σ from SM
Bernlochner et al. PRD95(2017)115008	0.299±0.003	0.257 ± 0.003		
Bigi et al. JHEP1711(2017)061		0.260 ± 0.008		
Jaiswal el al. JHEP1712(2017)060	0.299±0.004	0.257 ± 0.005		
HFLAV	<u>0.299±0.004</u>	<u>0.258 ± 0.005</u>	3.08	<u>2.5</u>
BaBar PRL123(2019),091801		0.253 ± 0.005	3.43	2.8
Gambino et al. PLB795(2019)386		0.254 ± 0.007	3.16	2.6
Bordone et al. ArXiv:1908.09398 (no exp.)	0.298±0.003	0.247 ± 0.006	3.77	3.2
Bordone el al. ArXiv:1908.09398	0.297±0.003	0.250 ± 0.003	3.87	3.2

Prediction using BaBar only result is compatible with the most recent predictions BaBar result can be included in the most recent calculations

The predictions that use the unpublished Belle result are systematically higher than the most recents

### Summary

- First tagged full 4-D angular analysis of  $B \rightarrow D^* \ell v$ 
  - Used both B<sup>0</sup> and B<sup>+</sup>
  - $|V_{cb}|$  consistent with world average for both CLN and BGL
    - persistent difference with inclusive determination
  - Updated prediction for R(D\*) consistent with most recent predictions
- Result published Phys.Rev.Lett. 123 (2019) no.9, 091801
  - Long PRD under internal review: more information + combined fit with  $B{\rightarrow}D\ell_V$
- Waiting for Lattice calculations at w>1
  - Many in the pipeline
- Results need confirmation from a Belle full 4-D angular analysis

## BACKUP



### Data/MC comparison

 1D-projections: data and simulated data re-weighted with fit results obtained with BGL parameterization



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