

$B_c \rightarrow J/\psi$ Form Factors and $R(J/\psi)$ using Lattice QCD

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$B_c \rightarrow J/\psi \ell \bar{\nu}_\ell$ Semileptonic decays

Semileptonic differential decay rate for $B_c \rightarrow J/\psi \ell \bar{\nu}_\ell$:

$$\frac{d\Gamma(B_c \rightarrow J/\psi \ell \bar{\nu}_\ell)}{dq^2} = \frac{G^2}{(2\pi)^3} |V_{cb}|^2 \frac{(q^2 - M_\ell^2)^2 p'}{12M_{B_c}^2 q^2} \left[(H_-^2 + H_0^2 + H_+^2) + \frac{M_\ell^2}{2q^2} (H_-^2 + H_0^2 + H_+^2 + 3H_t^2) \right]$$

Easy part: Kinematics, Electroweak perturbation theory

Hard part: QCD matrix elements between B_c and J/ψ states

$$H_{\pm}(q^2) = (M_{B_c} + M_{J/\psi}) A_1(q^2) \mp \frac{2M_{B_c} p'}{M_{B_c} + M_{J/\psi}} V(q^2),$$

$$\begin{aligned} H_0(q^2) = & \frac{1}{2M_{J/\psi}\sqrt{q^2}} \left((M_{B_c}^2 - M_{J/\psi}^2 - q^2) A_1(q^2) \right. \\ & \left. - 4 \frac{M_{B_c}^2 p'^2}{M_{B_c} + M_{J/\psi}} A_2(q^2) \right), \end{aligned}$$

$$H_t(q^2) = \frac{2M_{B_c} p'}{\sqrt{q^2}} A_0(q^2).$$

$$\langle J/\psi(p', \epsilon) | \bar{c} \gamma^\mu b | B_c^-(p) \rangle = \frac{2iV(q^2)}{M_{B_c} + M_{J/\psi}} \epsilon^{\mu\nu\rho\sigma} \epsilon_\nu^* p'_\rho p_\sigma$$

$$\langle J/\psi(p', \epsilon) | \bar{c} \gamma^\mu \gamma^5 b | B_c^- \rangle = 2M_{J/\psi} A_0(q^2) \frac{\epsilon^* \cdot q}{q^2} q^\mu$$

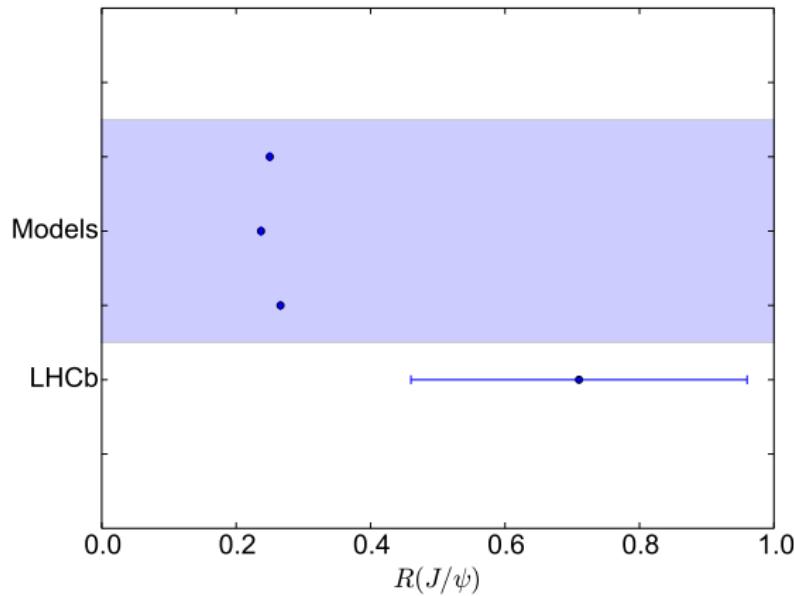
$$+ (M_{B_c} + M_{J/\psi}) A_1(q^2) \left[\epsilon^{*\mu} - \frac{\epsilon^* \cdot q}{q^2} q^\mu \right]$$

$$- A_2(q^2) \frac{\epsilon^* \cdot q}{M_{B_c} + M_{J/\psi}} \left[p^\mu + p'^\mu - \frac{M_{B_c}^2 - M_{J/\psi}^2}{q^2} q^\mu \right].$$

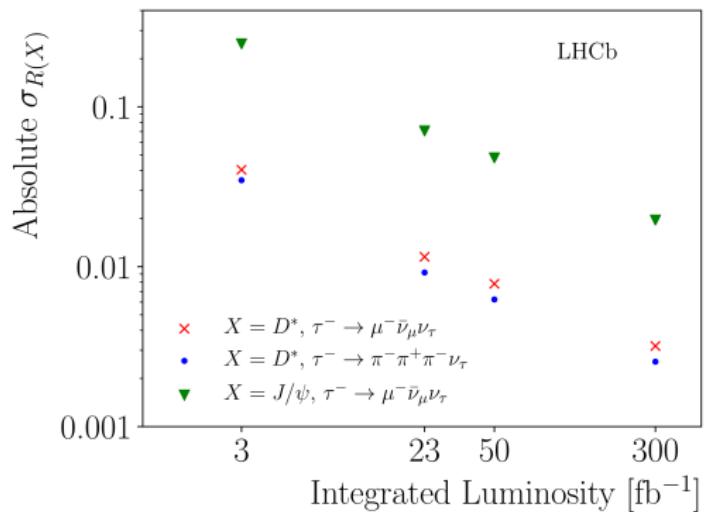
Experimental Status - $R(J/\psi)$

Useful to define dimensionless ratio of total decay rates to μ/e and τ final state.

$$R(J/\psi) = \frac{\Gamma(B_c \rightarrow J/\psi \tau^- \bar{\nu}_\tau)}{\Gamma(B_c \rightarrow J/\psi \mu^- \bar{\nu}_\mu)}$$



Experimental Status



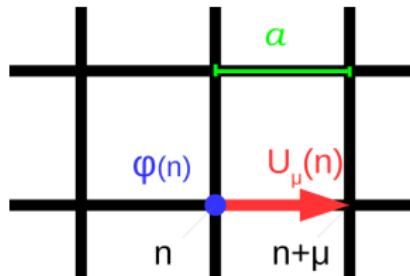
Projected uncertainties in $R(D^*)$ and $R(J/\psi)$ reproduced from
arXiv:1808.08865v4

Lattice QCD

We want to extract matrix elements, amplitudes and energies from Euclidean correlation functions computed in the path integral formalism,

$$\int \mathcal{D}[\psi, \bar{\psi}, A] \mathcal{O}_1(t) \mathcal{O}_2(0) e^{-S^E[\psi, \bar{\psi}, A]} = \sum_n \langle 0 | \hat{\mathcal{O}}_1 | n \rangle \langle n | \hat{\mathcal{O}}_2 | 0 \rangle e^{-E_n t},$$

- ▶ discretise QCD onto a lattice
- ▶ Fermion integrals exact \rightarrow need to invert dirac operator
- ▶ Monte-carlo integral over gauge fields U



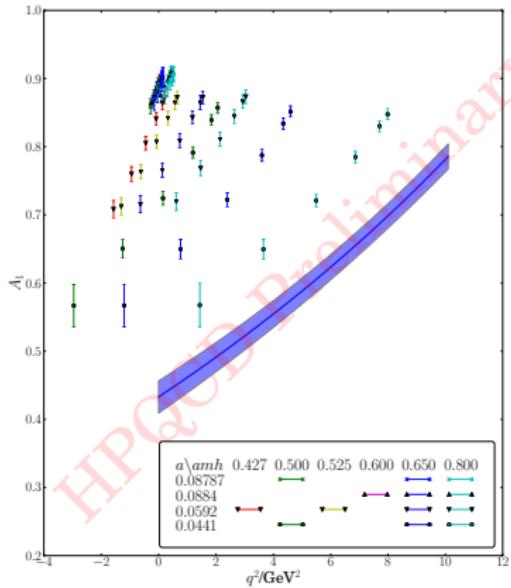
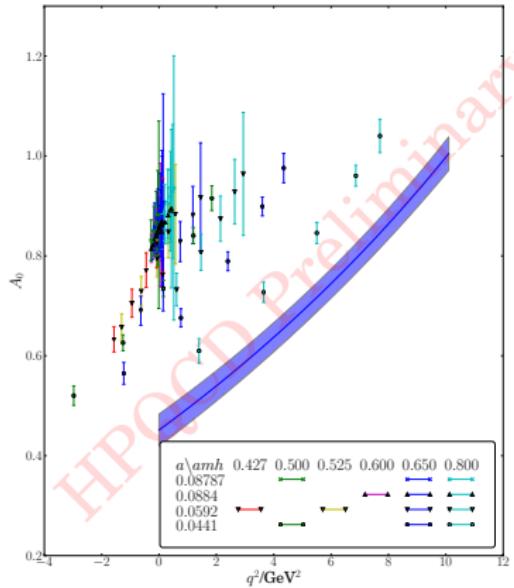
Heavy-HISQ

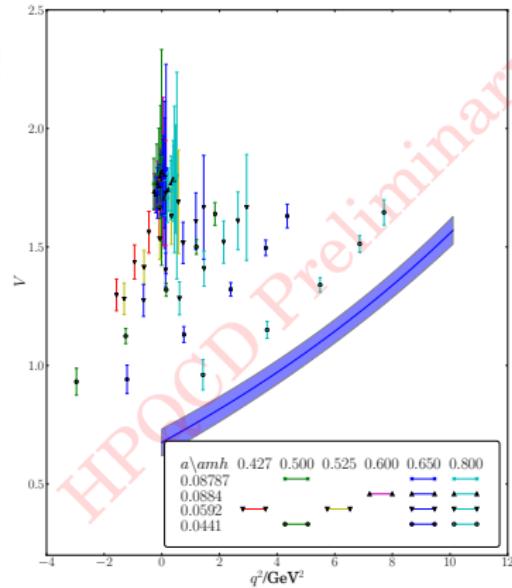
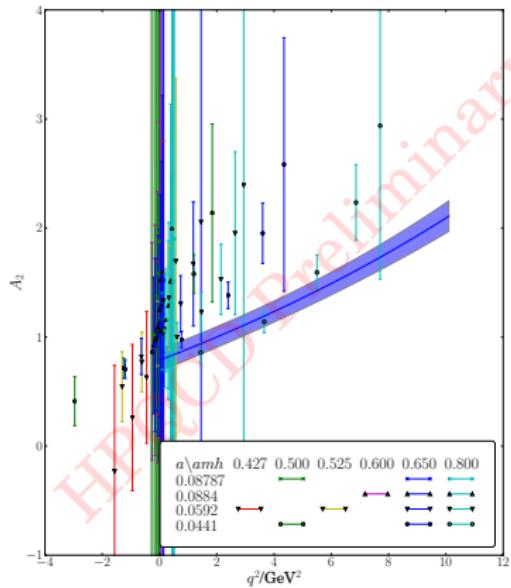
Discretisation effects enter as powers of am^n . Finest lattice has $am_b \approx 0.8 \rightarrow$ use unphysically light heavy quarks with $am_h < 1$ to constrain discretisation effects and heavy mass dependence in order to extract continuum physics.

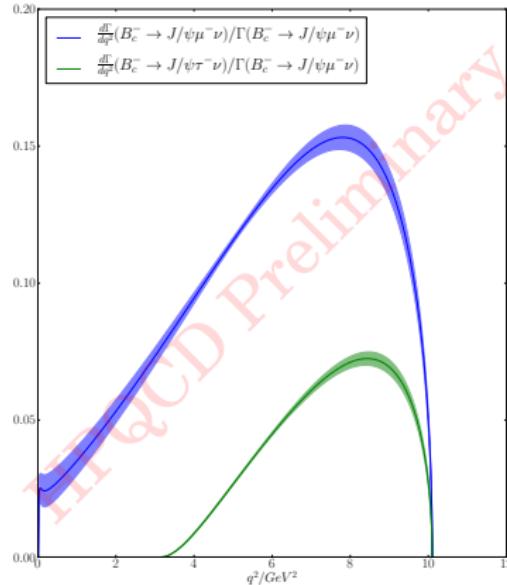
- ▶ use several am_h on each lattice.
- ▶ fit data to polynomial in $z(q^2)$ with coefficients dependent upon Λ_{QCD}/M_{η_h} , am_h and am_c .
- ▶ Set $M_{\eta_h} = M_{\eta_b}$, $am_h = am_c = 0$ in fit function to extract physical q^2 dependence.

We use Highly Improved Staggered Quarks (HISQ) \rightarrow significantly reduced discretisation errors, crucial for calculations involving heavy quarks. Using all HISQ set up also allows for non-perturbative renormalisation of lattice currents.

Results







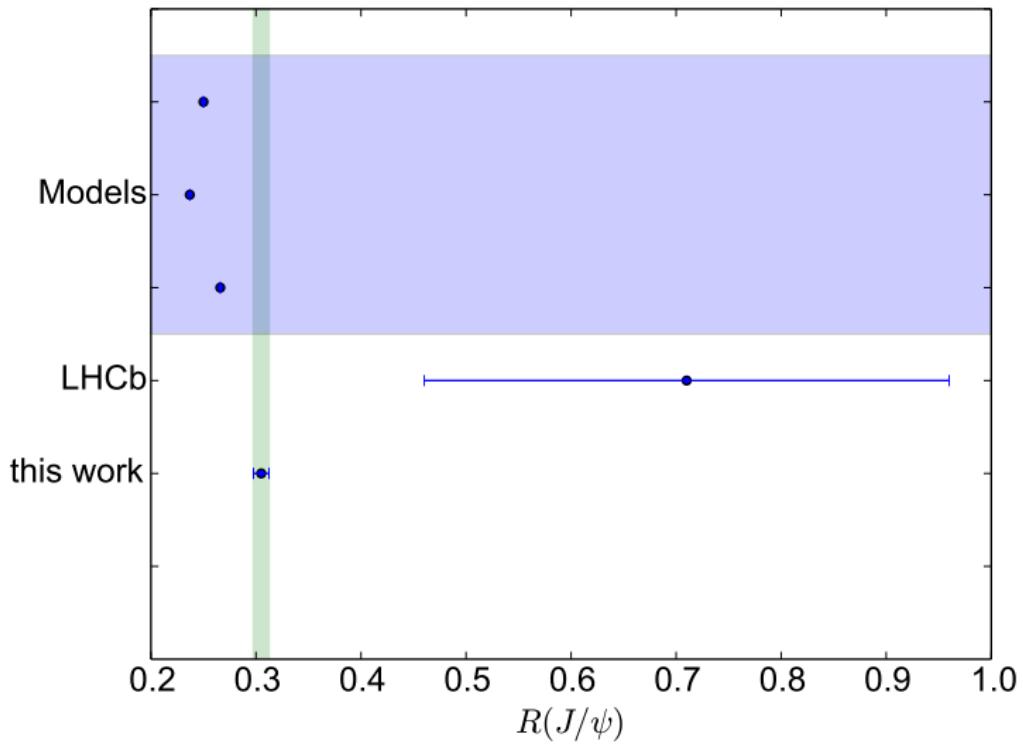
We find

$$\Gamma(B_c^- \rightarrow J/\psi \mu^- \bar{\nu}_\mu) = 2.14(14) \times 10^{10} s^{-1} \text{(Preliminary)},$$

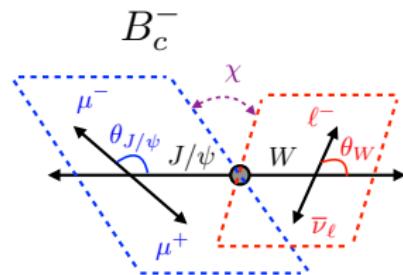
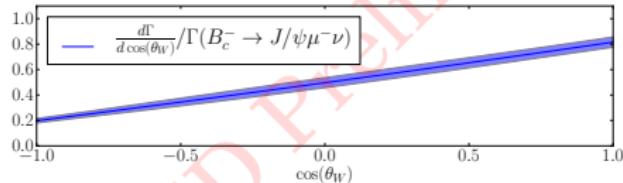
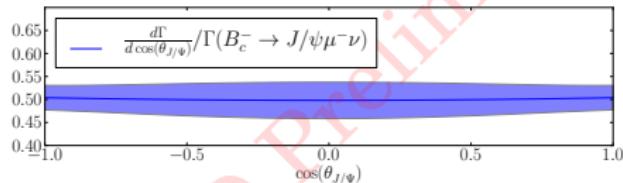
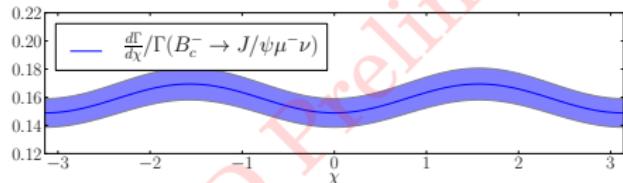
$$\Gamma(B_c^- \rightarrow J/\psi \tau^- \bar{\nu}_\tau) = 6.53(34) \times 10^9 s^{-1} \text{(Preliminary)},$$

and the ratio

$$R(J/\psi) = 0.3050(74) \text{(Preliminary)}.$$



Can also construct angular differential rates, with $J/\psi \rightarrow \mu^+ \mu^-$



Outlook

- ▶ We have computed $R(J/\psi)$ in the full SM for the first time.
- ▶ The precision of our result is expected to remain competitive with experiment for the foreseeable future.
- ▶ The similar calculation for $B_s \rightarrow D_s^*$ is underway, as is $B \rightarrow D^*$
 - Calculation of $R(D^*)$
 - Model independent determination of V_{cb}

Thanks