Introduction	Experiment	Lattice QCD approach	Interpolating operators	GEVP	Preliminary results	Conclusion

Z_b tetraquark channel with lattice QCD

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- fundamental quantum field theory of quarks (q) and gluons (g)
- the Lagrangian:
 - $\mathscr{L} = \sum_{q} \overline{\psi}_{q,s} \left(i \gamma^{\mu} \partial_{\mu} \delta_{ab} g_{s} \gamma^{\mu} t^{C}_{ab} A^{C}_{\mu} m_{q} \delta_{ab} \right) \psi_{q,b} \frac{1}{4} F^{A}_{\mu\nu} F^{A\mu\nu}$

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- ψ quark field
- A_{μ} gluon gauge field
- ab-initio predictive methods for QCD:
 - quantum chromodynamics on the lattice (lattice QCD)
 - perturbative expansions in the coupling

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- fundamental quantum field theory of quarks (q) and gluons (g)
- the Lagrangian:

$$\mathscr{L} = \sum_{q} \overline{\psi}_{q,a} \left(i \gamma^{\mu} \partial_{\mu} \delta_{ab} - g_{s} \gamma^{\mu} t^{C}_{ab} A^{C}_{\mu} - m_{q} \delta_{ab} \right) \psi_{q,b} - \frac{1}{4} F^{A}_{\mu\nu} F^{A\mu\nu}$$

$$\psi$$
 – quark field

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$$A_{\mu}$$
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- ab-initio predictive methods for QCD:
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- The theory is discretised onto a four-dimensional spacetime lattice
 - quark fields lattice sites
 - gauge fields links between sites
- lattice spacing a
- quantities extracted from Euclidean (imaginary-time) correlation functions
- Monte-Carlo methods



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Experimental discovery of the Z_b tetraquarks

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• the Belle collaboration first observed two charged bottomonium-like resonances in $\Upsilon(5S) \rightarrow \Upsilon(nS)\pi^+\pi^-, h_b(mP)\pi^+\pi^-,$ (n = 1, 2, 3; m = 1, 2)

Introduction

name	valence quark content	m [MeV]	$I^G(J^{PC})$	discovery year	reference
$Z_b(10610)$		10607.2	$1^+(1^{+-})$	2011	 [A. Bondar et al. (Belk), Phys. Rev. Lett. 108, 122001 (2012).] and
$Z_b(10650)$	ЬБ	10652.2	$1^+(1^{+-})$	2011	[A. Garmash et al. (Belle), Phys. Rev. D91, 072003 (2015).] for both

- $Z_b(10610)$ and $Z_b(10650)$ decay also to $B\overline{B}^*$ and $B^*\overline{B}^*$ respectively
 - these are the dominant decay channels



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Experimental discovery of the Z_b tetraquarks

- properties incompatible with a $q\overline{q}$ structure
- masses few MeV above the thresholds for the open beauty channel $B\overline{B}^*$ and $B^*\overline{B}^*$
- this suggests a "molecular" nature of these states, which might explain most of their observed properties

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Lattice QCD - scattering analysis

- Lüscher scattering formalism:
 - eigen-energies of a correlation function
 ⇒ two hadron scattering amplitudes
 ⇒ mass and decay width
- rigorously treating Z_b tetraquark with lattice QCD:
 - scattering matrix for at least 7 coupled channels
 - a severe challenge



Figure: non-interacting energies $E(L) = \sum_{i=1,2} \sqrt{m_i^2 + p_i^2} + \Delta E,$ $\vec{p}_i = \frac{2\pi}{L} \vec{n}_i$

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- b and \overline{b} infinitely heavy $(m_b
 ightarrow \infty)$ and static
 - on a distance r
 - their spins conserved quantities
- compute the potential V(r) of the static quarks in the presence of the light quarks, \Rightarrow potential of B and \overline{B}

Born-Oppenheimer approximation:

 solve the Schrödinger equation with the computed potential V(r)



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Symmetries of the static limit

- j_{z,light} = z-component of the angular momentum of the light quarks
- $C \circ P =$ combined parity and charge conjugation
- ε = P_x = x-parity = reflection along an axis orthogonal to the separation axis e. g. x-axis
- my advisor has considered the case $S_{\text{heavy}} = 1$, $j_{z,\text{light}} = 0$, CP = -1, $\varepsilon = -1$ [S. Prelovšek, H. Bahtiyar and J. Petković (2019), [arXiv:1912.02656].]







shematic presentation of operators for a specific set of quantum numbers $(S_{\text{heavy}} = 0, j_{z,\text{light}} = 0, CP = +1, \varepsilon = +1)$: $O_{7} = O^{h_{b}a_{0}(0)} \propto \left[\overline{b}P_{-}\gamma_{5}b\right] \left[\overline{q}\mathbb{I}q\right]_{\vec{n}=\vec{0}}$ $O_8 = O^{h_b a_0(1)} \propto \left[\overline{b}P_-\gamma_5 b\right] \left(\left[\overline{q}\mathbb{I}q\right]_{\vec{p}=\hat{e}_z} + \left[\overline{q}\mathbb{I}q\right]_{\vec{p}=-\hat{e}_z} \right)$

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• two-point correlation functions

$$C_{ij}(t) = \langle 0 | O_i(t) O_j^{\dagger}(0) | 0
angle = \sum_n \langle 0 | O_i | n
angle \langle n | O_j^{\dagger} | 0
angle e^{-E_n t}$$

• overlaps:

$$Z_i^{(n)} = \langle 0 | \hat{O}_i | n \rangle, \qquad Z_j^{(n)*} = \langle n | \hat{O}_j^{\dagger} | 0 \rangle$$

eigen-energies E_n extracted using the GEVP approach

$$\sum_{j=1}^{N} C_{ij}(t) v_{j}^{(n)}(t,t_{0}) = \sum_{j=1}^{N} \lambda^{(n)}(t,t_{0}) C_{ij}(t_{0}) v_{j}^{(n)}(t,t_{0}),$$

where $n = 1, \ldots, N$ and $t > t_0$

$$\lambda^{(n)}(t,t_0) = e^{-E_n(t-t_0)} = Ae^{-E_n t}$$

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• effective energies:

$$E_{\mathrm{eff}}^{(n)}(t,t_0)=-rac{\partial\log\left(\lambda^{(n)}(t,t_0)
ight)}{\partial t},\qquad E_n=\lim_{t o\infty}E_{\mathrm{eff}}^{(n)}(t,t_0)$$

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example for
$$r = 1$$
, $t_0 = 2$ and these 5 interpolating operators
a $O_1 = O^{B\bar{B}^*} \propto (P_-\gamma_3)_{AB}(P_-\gamma_5)_{CD} \ (\bar{b}_C q_B) \ (\bar{q}_A b_D)$
a $O_2 = O^{B\bar{B}^{*'}} \propto (P_-\gamma_3)_{AB}(P_-\gamma_5)_{CD} \ (\bar{b}_C \nabla^2 q_B) \ (\nabla^2 \bar{q}_A b_D)$
a $O_3 = O^{\eta_b \rho(0)} \propto [\bar{b}P_-\gamma_5 b] \ [\bar{q}\gamma_3 q]_{\vec{p}=\vec{0}}$
a $O_4 = O^{\eta_b \rho(1)} \propto [\bar{b}P_-\gamma_5 b] \ ([\bar{q}\gamma_3 q]_{\vec{p}=\hat{e}_z} + [\bar{q}\gamma_3 q]_{\vec{p}=-\hat{e}_z})$
b $O_5 = O^{\eta_b \rho(2)} \propto [\bar{b}P_-\gamma_5 b] \ ([\bar{q}\gamma_3 q]_{\vec{p}=2\hat{e}_z} + [\bar{q}\gamma_3 q]_{\vec{p}=-2\hat{e}_z})$



• effective energies:

$$E_{\rm eff}^{(n)}(t,t_0) = -\frac{\partial \log(\lambda^{(n)}(t,t_0))}{\partial t}, \qquad E_n = \lim_{t \to \infty} E_{\rm eff}^{(n)}(t,t_0)$$

example for $S_{\rm heavy}=$ 0, $j_{z,{\rm light}}=$ 0, CP=+ 1, $\epsilon=+$ 1, r= 1, $t_0=$ 2:





Preliminary results



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Conclu	sion					

- there are many phenomenological studies on spectroscopy of exotic hadrons with heavy quarks
- but only one group (except of my supervisor) made studies based on the fundamental theory lattice QCD considering the Z_b tetraquark
- this tetraquark is due to experimental discovery of great interest



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• looks like no attraction in my channel

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