Pseudo-Dirac heavy neutrinos in low scale seesaw models

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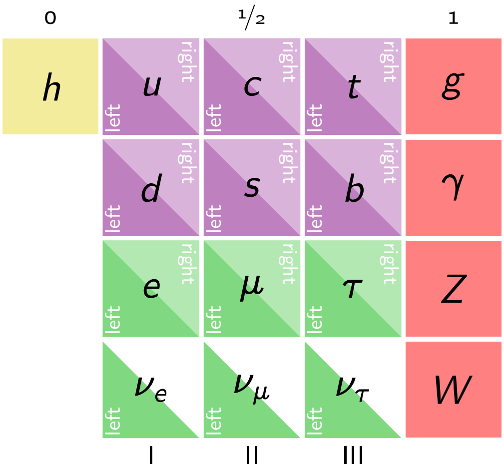
Standard Model neutrinos

Standard Model particle content

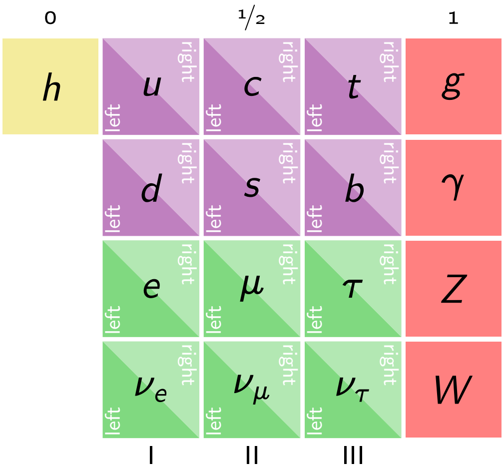


Neutrinos ν_α stand out

purely left-chiral and massless

Standard Model neutrinos

Standard Model particle content



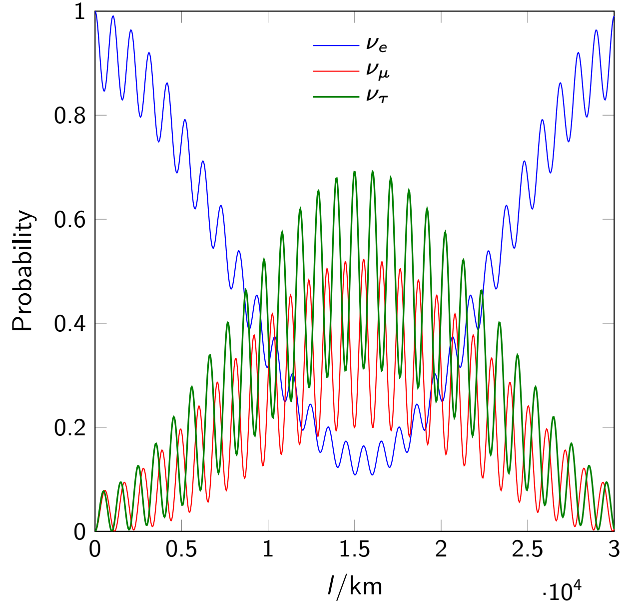
Neutrinos ν_α stand out

purely left-chiral and massless

Right-chiral or sterile Neutrinos

neutral under SM symmetries

Observed neutrino flavour oscillations



Flavour oscillations are explained by

right-chiral neutrinos allowing mass terms

Seesaw models generating neutrino masses

Dirac mass

$$\mathcal{L}_D = -m_{D\alpha}\bar{\nu}_\alpha N + \text{h.c.}\,, \qquad \boldsymbol{m}_D = v\boldsymbol{y}$$

Majorana mass

$$\mathcal{L}_M = -\frac{1}{2} m_M \overline{N} N^c + \text{h.c.}$$

Coupling strength is determined by

$$\boldsymbol{\theta} = \boldsymbol{m}_D / m_M$$

Majorana mass introduces

lepton number violation (LNV)

Majorana mass vanishes if

lepton-number L is conserved

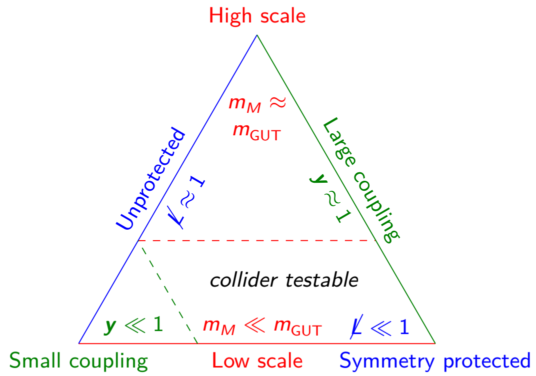
Neutrino oscillation pattern requires

at least two massive neutrinos

Neutrino mass matrix from two sterile neutrinos

$$M_\nu = \frac{\boldsymbol{m}_D^{(1)} \otimes \boldsymbol{m}_D^{(1)}}{m_M^{(1)}} + \frac{\boldsymbol{m}_D^{(2)} \otimes \boldsymbol{m}_D^{(2)}}{m_M^{(2)}}$$

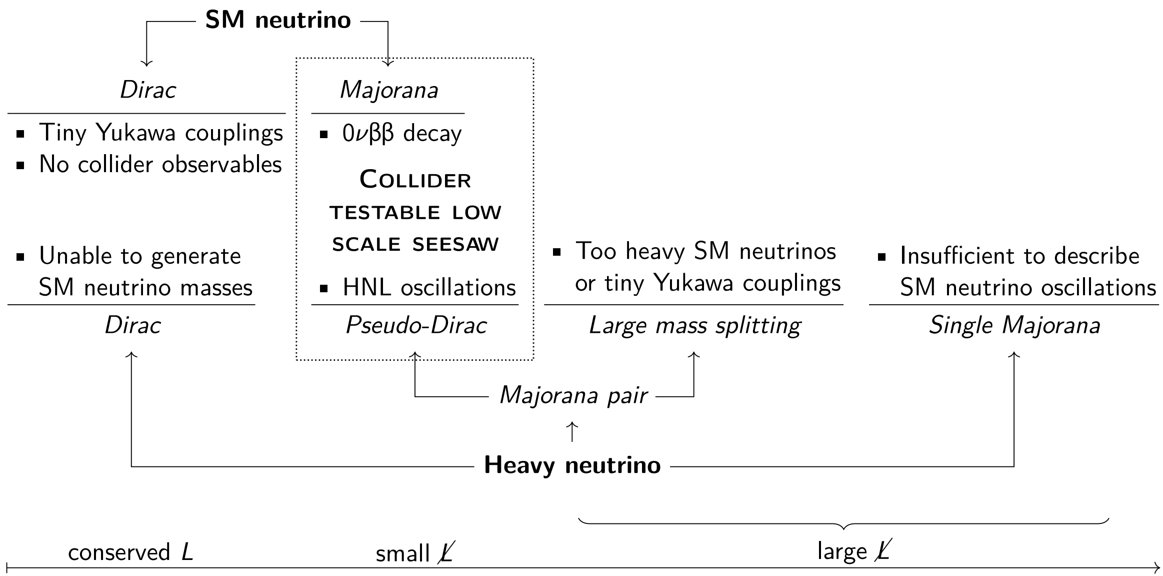
Viable seesaw models



Neutrino masses are small for

- small $\boldsymbol{y}$
- large m_M
- symmetry protected cancellation

Are HNLs Majoran or Dirac Fermions? [2210.10738]



Single Majorana and Dirac HNLs are

- not predicted by low-scale seesaw models

Unique phenomenology of pseudo-Dirac HNLs

- Heavy neutrino-antineutrino oscillations
- $0 < R_{ll} = \frac{N_{\text{LNV}}}{N_{\text{LNC}}} < 1$
- Governed by mass splitting Δm

Symmetry protected seesaw scenario (SPSS)

[2210.10738]

Exact limit

$$\mathcal{L}^{L}_{\text{SPSS}} = -m_M \overline{N}_1 N_2^c - y_1 \widetilde{H}^\dagger \bar{\ell} N_1^c + \text{h.c.}$$

Small breaking terms $v y_2 \approx \mu_M \approx \mu'_M \ll m_M$

$$\mathcal{L}^{\not L}_{\text{SPSS}} = -y_2 \widetilde{H}^\dagger \bar{\ell} N_2^c - \mu'_M \overline{N}_1 N_1^c - \mu_M \overline{N}_2 N_2^c + \text{h.c.}$$

Lepton number-like symmetry

generalises accidental SM lepton number L

One simple choice of charges

	ℓ	N_1	N_2
L	$+1$	-1	$+1$

Other new fields

further terms in Lagrangian

Neutrino mass matrix M_n

contains seesaw information

Basis

$n = (\nu, n_4, n_5)$

Dirac masses

$\boldsymbol{m}_D = \boldsymbol{y}_1 v, \quad \boldsymbol{\mu}_D = \boldsymbol{y}_2 v$

Symmetric limit

$$M_n^L = \begin{pmatrix} 0 & \boldsymbol{m}_D & 0 \\ \boldsymbol{m}_D^\mathsf{T} & 0 & m_M \\ 0 & m_M & 0 \end{pmatrix}$$

- Massless neutrinos $M_\nu = 0$
- Dirac HNL

Mild symmetry breaking

$$M_n^{\not L \ll 1} = \begin{pmatrix} 0 & \boldsymbol{m}_D & \boldsymbol{\mu}_D \\ \boldsymbol{m}_D^\mathsf{T} & \mu'_M & m_M \\ \boldsymbol{\mu}_D^\mathsf{T} & m_M & \mu_M \end{pmatrix}$$

- Pseudo-Dirac HNL (small Δm Majorana pair)
- Phenomenology governed by small parameters μ

Large symmetry breaking

$$M_n^{\not L \gg 0} = \begin{pmatrix} 0 & \boldsymbol{m}_D & \widehat{\boldsymbol{m}}_D \\ \boldsymbol{m}_D^\mathsf{T} & \widehat{m}'_M & m_M \\ \widehat{\boldsymbol{m}}_D^\mathsf{T} & m_M & \widehat{m}_M \end{pmatrix}$$

- Large Δm Majorana pair
- Requires large m_M or tiny $\boldsymbol{\theta}$

Special cases captured by the symmetry protected seesaw

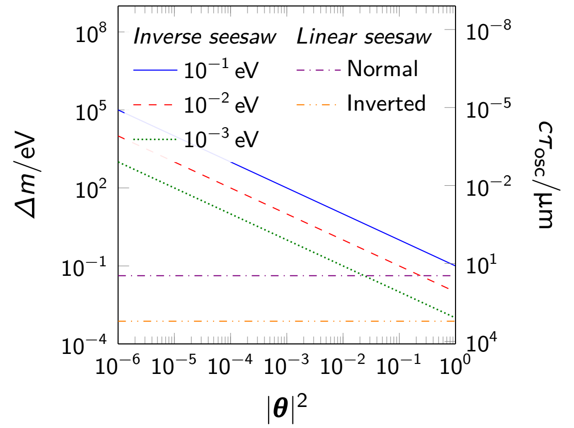
	Linear seesaw $\boldsymbol{\mu}_D$	Inverse seesaw μ_M	Seesaw independent μ'_M
$M_n =$	$\begin{pmatrix} 0 & \boldsymbol{m}_D & \boldsymbol{\mu}_D \\ \boldsymbol{m}_D^\mathsf{T} & 0 & m_M \\ \boldsymbol{\mu}_D^\mathsf{T} & m_M & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & \boldsymbol{m}_D & 0 \\ \boldsymbol{m}_D^\mathsf{T} & 0 & m_M \\ 0 & m_M & \mu_M \end{pmatrix}$	$\begin{pmatrix} 0 & \boldsymbol{m}_D & 0 \\ \boldsymbol{m}_D^\mathsf{T} & \mu'_M & m_M \\ 0 & m_M & 0 \end{pmatrix}$
$M_\nu =$	$\boldsymbol{\mu}_D \otimes \boldsymbol{\theta}$	$\mu_M \boldsymbol{\theta} \otimes \boldsymbol{\theta}$	0 (at tree level)
$\Delta m =$	Δm_ν	$m_\nu \lvert\boldsymbol{\theta}\rvert^{-2}$	$\lvert\mu'_M\rvert$

Benchmark models

Seesaw	Hierarchy	BM
Linear	Normal	$\Delta m_\nu = 42.3\,\text{meV}$
	Inverted	$\Delta m_\nu = 748\,\mu\text{eV}$
Inverse		$m_\nu = 1\,\text{meV}$
		$m_\nu = 10\,\text{meV}$
		$m_\nu = 100\,\text{meV}$

Generic seesaw

All small parameter μ are nonzero



Heavy neutrino-antineutrino oscillations

Heavy neutrino-antineutrino oscillations

Oscillations

between LNC and LNV processes

Damping due to decoherence

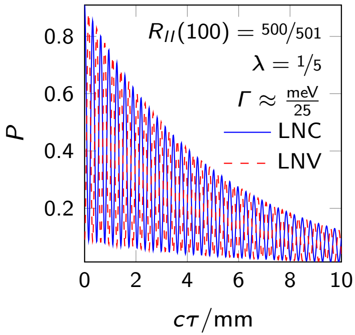
governed by λ

Oscillation length

governed by mass splitting Δm

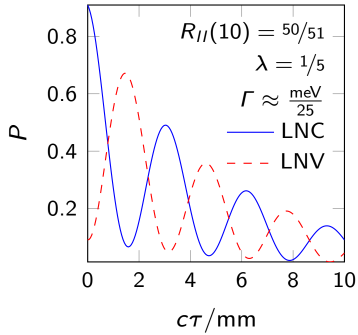
$$P_{\text{osc}}^{\text{LNC/LNV}}(\tau) = \frac{1 \pm \cos(\Delta m \tau)\exp(-\lambda)}{2}$$

Short oscillation length



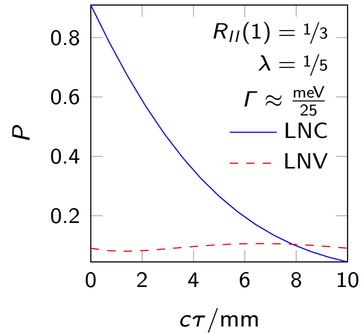
- Oscillations not resolvable
- Large R_{ll}
- 'Majorana' limit

Intermediate oscillation length



- Oscillations potentially measurable
- Pseudo-Dirac character crucial

Long oscillation length



- LNV strongly suppressed
- Small R_{ll}
- 'Dirac' limit

Integrated effect

Oscillation probability

$$P_{\text{osc}}^{\text{LNC/LNV}}(\tau) = \frac{1 \pm \cos(\Delta m \tau) \exp(-\lambda)}{2}$$

Decay probability density

$$P_{\text{decay}}(\tau) = -\frac{d}{d\tau} \exp(-\Gamma \tau) = \Gamma \exp(-\Gamma \tau)$$

Probability to decay with oscillation between τ_{min} and τ_{max}

$$P_{ll}^{\text{LNC/LNV}}(\tau_{\text{min}}, \tau_{\text{max}}) = \int_{\tau_{\text{min}}}^{\tau_{\text{max}}} P_{\text{osc}}^{\text{LNC/LNV}}(\tau) P_{\text{decay}}(\tau)\, d\tau$$

Integrated

$$P_{ll}^{\text{LNC/LNV}}(\tau_{\text{min}}, \tau_{\text{max}}) = \Gamma \frac{P^{\text{LNC/LNV}}(\tau_{\text{max}}) - P^{\text{LNC/LNV}}(\tau_{\text{min}})}{2}$$

where

$$P^{\text{LNC/LNV}}(\tau) = P(\tau, \Gamma, 0) \pm \frac{P(\tau, \Gamma_-, \lambda) + P(\tau, \Gamma_+, \lambda)}{2}$$

In the limit

$$\tau_{\text{min}} \to 0\,, \quad \tau_{\text{max}} \to \infty\,, \quad \lambda \to 0$$

Probability simplifies

$$P_{ll}^{\text{LNC/LNV}} = \frac{1}{2} \begin{cases} \frac{\Gamma^2}{\Delta m^2 + \Gamma^2} + 1 & \text{LNC} \\ \frac{\Delta m^2}{\Delta m^2 + \Gamma^2} & \text{LNV} \end{cases}$$

Ratio is easily measurable

$$R_{ll} = \frac{P_{ll}^{\text{LNV}}}{P_{ll}^{\text{LNC}}} = \frac{\Delta m^2}{\Delta m^2 + 2\Gamma^2}$$

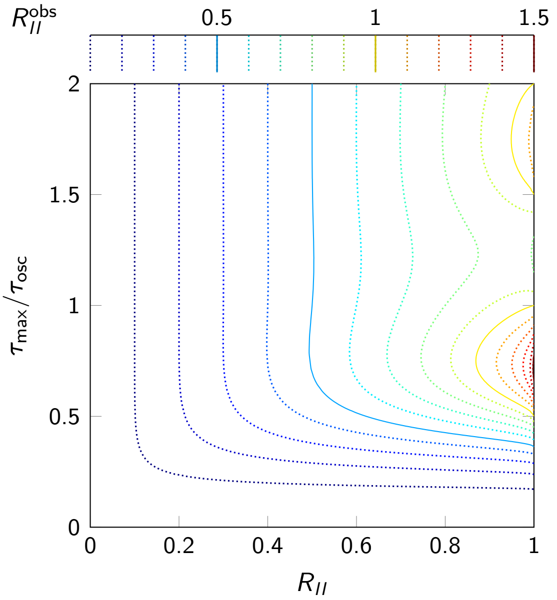
with

$$P(\tau, \Gamma, \lambda) = -\frac{e^{-\lambda - \Gamma \tau}}{\Gamma}$$

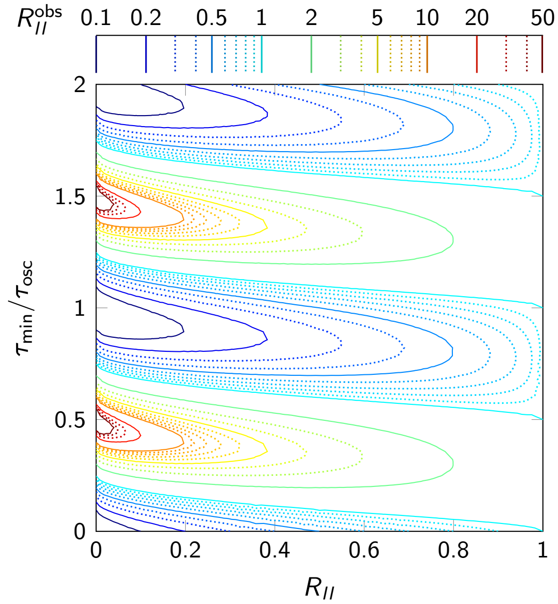
$$\Gamma_\pm = \Gamma \pm i \Delta m$$

$R_{//}$ with finite detector size: $R_{//}^{\mathrm{obs}}$

$R_{//}^{\mathrm{obs}}(R_{//}, \tau_{\mathrm{max}}/\tau_{\mathrm{osc}})$ while $\tau_{\mathrm{min}} \to 0$



$R_{//}^{\mathrm{obs}}(R_{//}, \tau_{\mathrm{min}}/\tau_{\mathrm{osc}})$ while $\tau_{\mathrm{max}} \to \infty$



Software implementation of the phenomenological SPSS

[pSPSS; 2210.10738]

Mass splitting

$$m_{4/5} = m_M\left(1 + |\boldsymbol{\theta}|^2/2\right) \mp \Delta m/2$$

Phenomenological SPSS (pSPSS) adds

Δm Heavy neutrino-antineutrino oscillations

λ Decoherence damping

FeynRules model file

Pseudo-Dirac HNLs in the pSPSS

Available online

feynrules.irmp.ucl.ac.be/wiki/pSPSS

Parameter

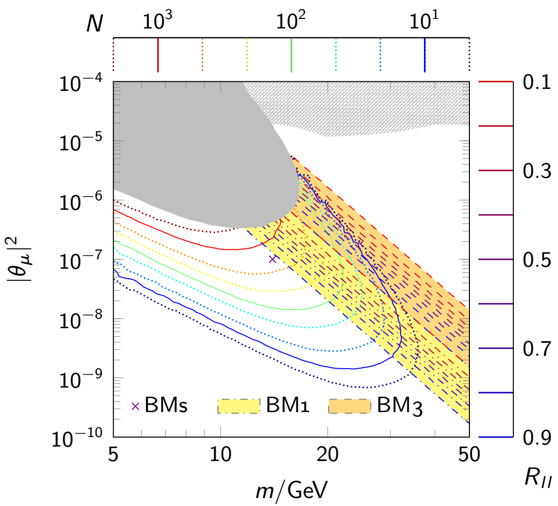
BLOCK PSPSS #		
1	1.000000e+02	# mmaj
2	1.000000e-12	# deltam
3	0.000000e+00	# theta1
4	1.000000e-03	# theta2
5	0.000000e+00	# theta3
6	0.000000e+00	# damping

Oscillations implemented in MadGraph

```
mass_splitting = param_card.get_value('PSPSS', 2)
damping = param_card.get_value('PSPSS', 6)
for event in lhe:
    leptonnumber = 0
    write_event = True
    for particle in event:
        if particle.status == 1:
            if particle.pid in [11, 13, 15]:
                leptonnumber += 1
            elif particle.pid in [-11, -13, -15]:
                leptonnumber -= 1
    for particle in event:
        id = particle.pid
        width = param_card['decay'].get((abs(id),)).value
        if width:
            if id in [8000011, 8000012]:
                tau0 = random.expovariate(width / cst)
                if 0.5 * (1 + math.exp(-damping)*math.cos(
                      mass_splitting * tau0 / cst)) >= random.
                      random():
                    write_event = (leptonnumber == 0)
                else:
                    write_event = (leptonnumber != 0)
                vtim = tau0 * c
            else:
                vtim = c * random.expovariate(width / cst)
            if vtim > threshold:
                particle.vtim = vtim
    # write this modify event
    if write_event:
        output.write(str(event))
output.write('</LesHouchesEvents>\n')
output.close()
```

Monte Carlo Simulation

HL-LHC event number with $\mathscr{L} = 3\,\mathrm{ab}^{-1}$

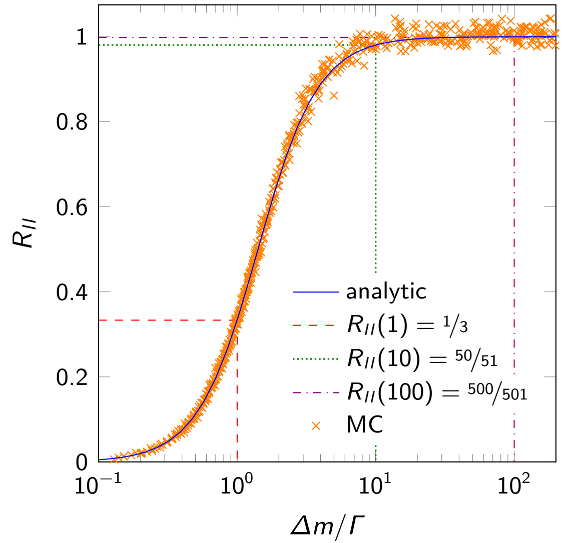


BM	$\Delta m/\mu\mathrm{eV}$	$c\tau_{\mathrm{osc}}/\mathrm{mm}$	R_{ll}
1	82.7	15	0.9729
2	207	6	0.9956
3	743	1.67	0.9997

Integrate oscillations from origin to infinity

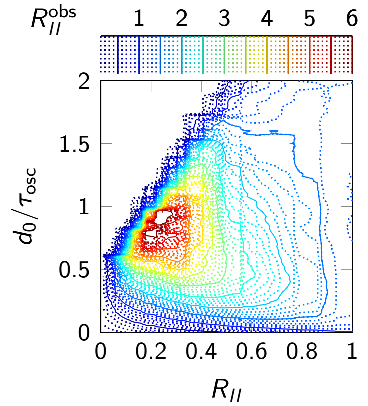
$$R_{ll} = \frac{N^{\mathrm{LNV}}}{N^{\mathrm{LNC}}} = \frac{\Delta m^2}{\Delta m^2 + 2\Gamma^2} .$$

R_{ll} simulation vs. calculation

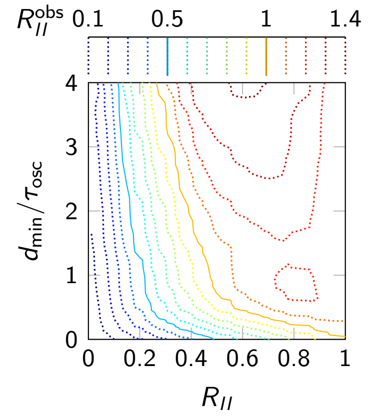


R_{ll}^{obs} with physical cuts

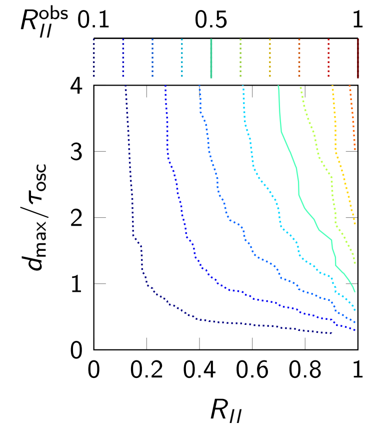
d_0 cut for $m = 10$ GeV



d_{min} cut for $m = 10$ GeV



d_{max} cut for $m = 10$ GeV

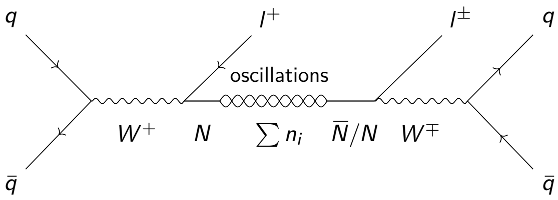


Transition to physical cuts d_0, d_{min}, d_{max} [2210.10738]

Changes feature but non-trivial pattern remains

Heavy neutrino-antineutrino oscillations at the LHC

Production, oscillation, and decay



Idea

Observe heavy neutrino-antineutrino oscillations in long-lived decays

Process

- Production of interaction eigenstates N or $\overline{N}$
- Oscillations between n_4 and n_5 due to Δm
- LNC decay into l^- or LNV decay into l^+

Simulation

- Model implementation in FeynRules
- Event generation in MadGraph
- CMS Detector simulation in Delphes

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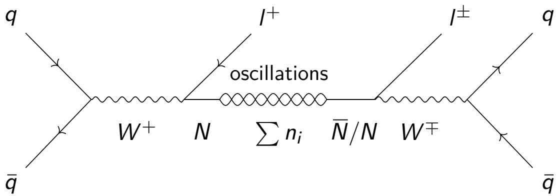
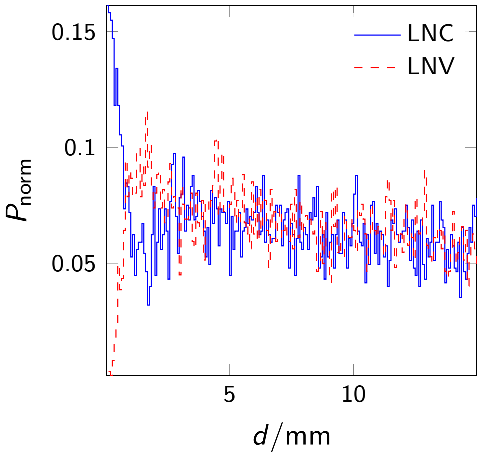


Lab frame



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Observations after MadGraph

- No oscillations in lab frame

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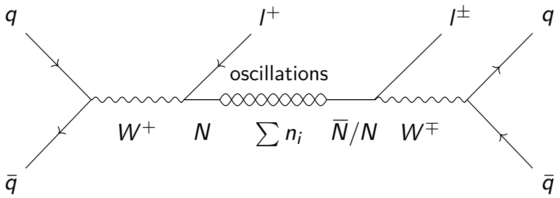
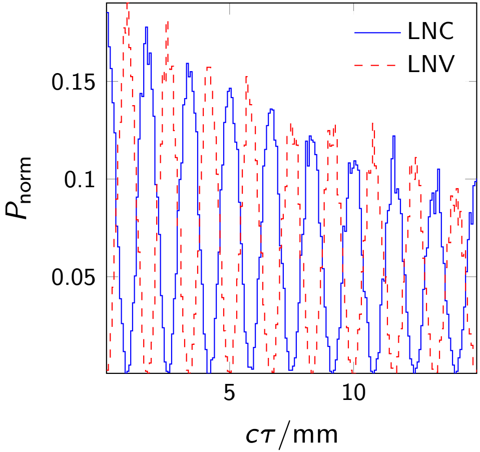


Proper time frame



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Simulation

- Model implementation in FeynRules
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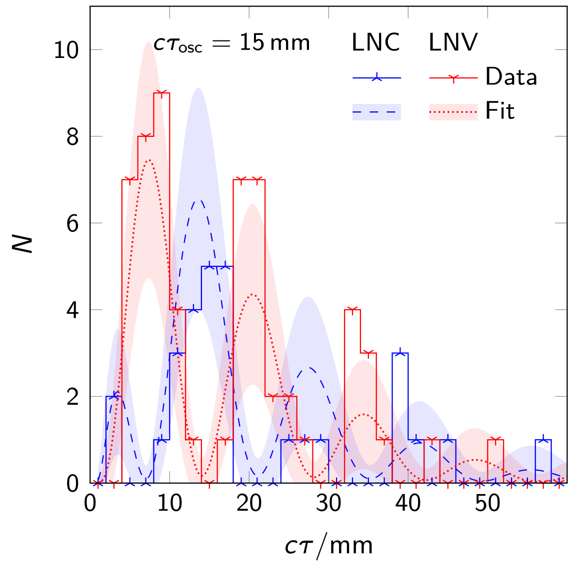
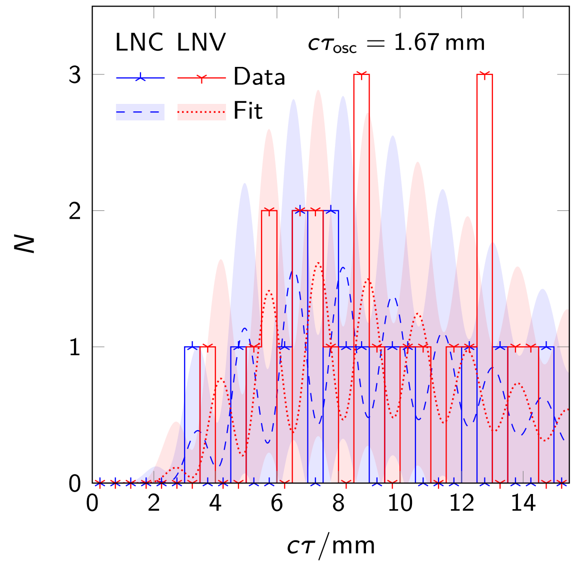
Observations after MadGraph

- No oscillations in lab frame
- Oscillations appear in proper time frame
- Crucial to reconstruct Lorentz factor γ
- Depends on final states without neutrinos

BM1 with $c\tau_{osc} = 15\,\mathrm{mm}$ and $Z = 6.66\,\sigma$



BM3 with $c\tau_{osc} = 1.67\,\mathrm{mm}$ and $Z = 0.67\,\sigma$



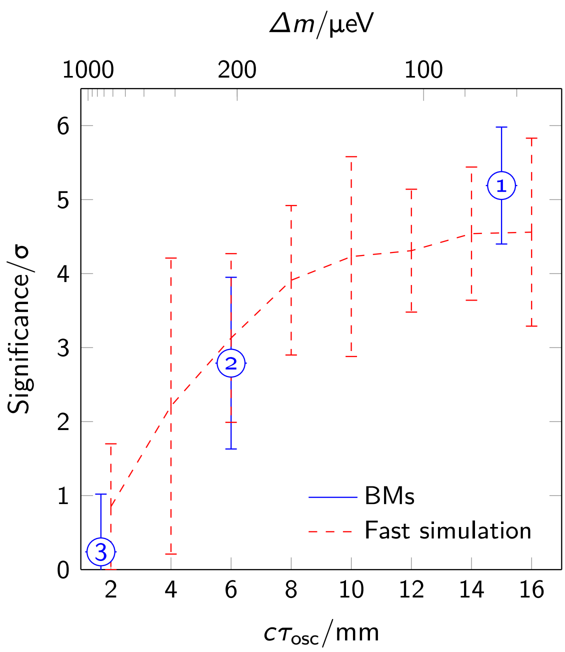
Results

- Large parts of accessible parameter space excluded by LHC
- HL-LHC can measure oscillations in some BMs with $5\,\sigma$

Detector simulation results

[2212.00562]

Discovery potential



HL-LHC

discovery possible

Large mass splitting hard to resolve
Neutrino Lorentz factor reconstruction crucial

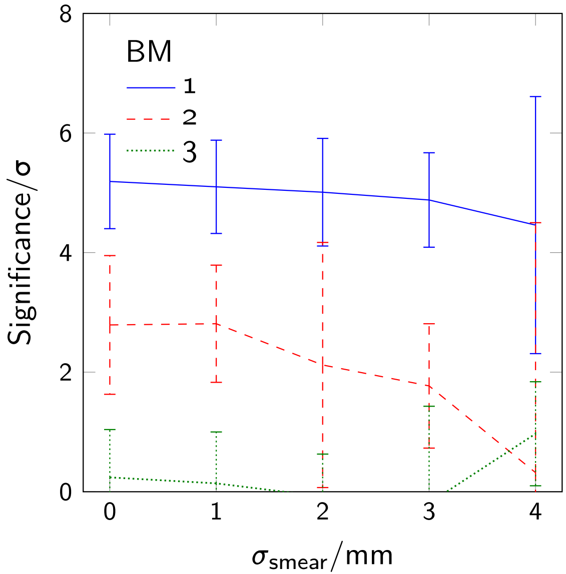
Future work

Properly simulate secondary vertex smearing
Improve Lorentz factor reconstruction

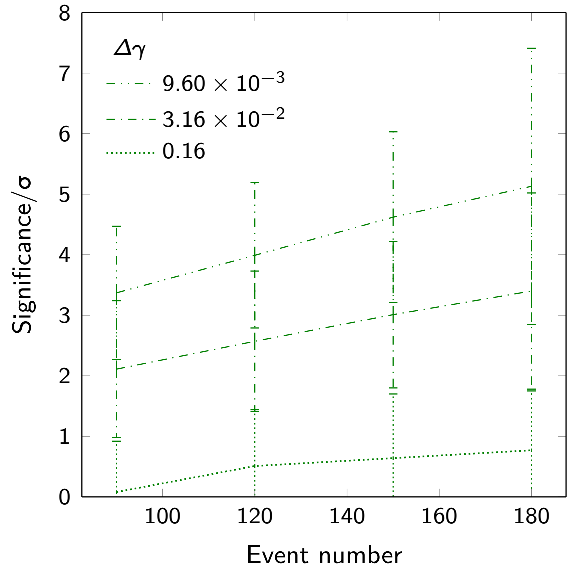
Detector simulation results

[2212.00562]

Secondary vertex smearing dependency



Lorentz factor reconstruction dependency



$$\Delta\gamma = \frac{|\gamma_{gen}^2 - \gamma_{reco}^2|}{\gamma_{gen}^2 - 1},$$

Angular dependence of the transverse impact parameter

Transverse impact parameter

$$d_0 = \frac{\boldsymbol{d}'_T \wedge \boldsymbol{p}'_T}{p'_T} = \frac{\epsilon_{ij} x'_i p'_j}{p'_T} = \frac{x' p'_y - y' p'_x}{p'_T}$$

Point with the smallest distance to the z-axis

$$\boldsymbol{d}'_T = (x', y')$$

Transverse momentum at $\boldsymbol{d}'_T$

$$\boldsymbol{p}'_T = (p'_x, p'_y)$$

Small B fields and going to secondary vertex

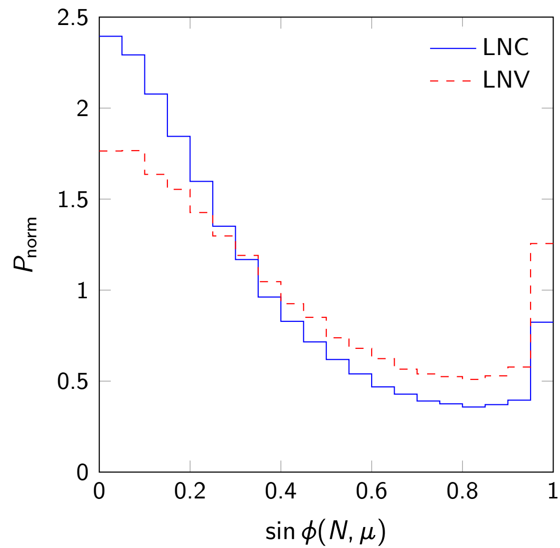
$$\boldsymbol{d}'_T \to \boldsymbol{d}_T = d_T \frac{\boldsymbol{p}_T^N}{p_T^N}$$

Approximation

$$d_0 \approx d_T \frac{\boldsymbol{p}_T^N \wedge \boldsymbol{p}_T^\mu}{p_T^N p_T^\mu} = d_T \sin\phi(N, \mu)$$

Angular dependence

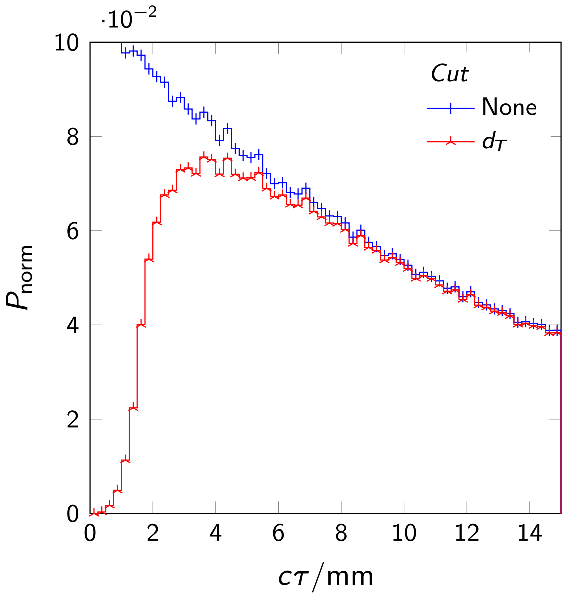
$$\sin\phi(N, \mu) = \frac{\boldsymbol{p}_T^N \wedge \boldsymbol{p}_T^\mu}{p_T^N p_T^\mu}$$



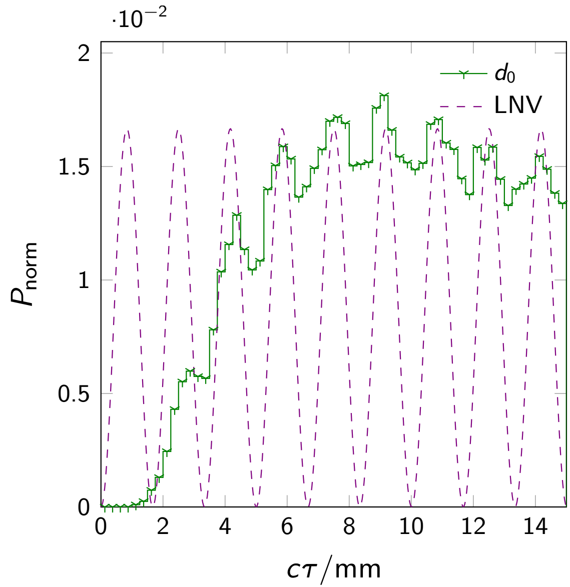
d_0 cuts introduce angular (spin) dependency

Total number of events $N = N_{\mathrm{LNC}} + N_{\mathrm{LNV}}$

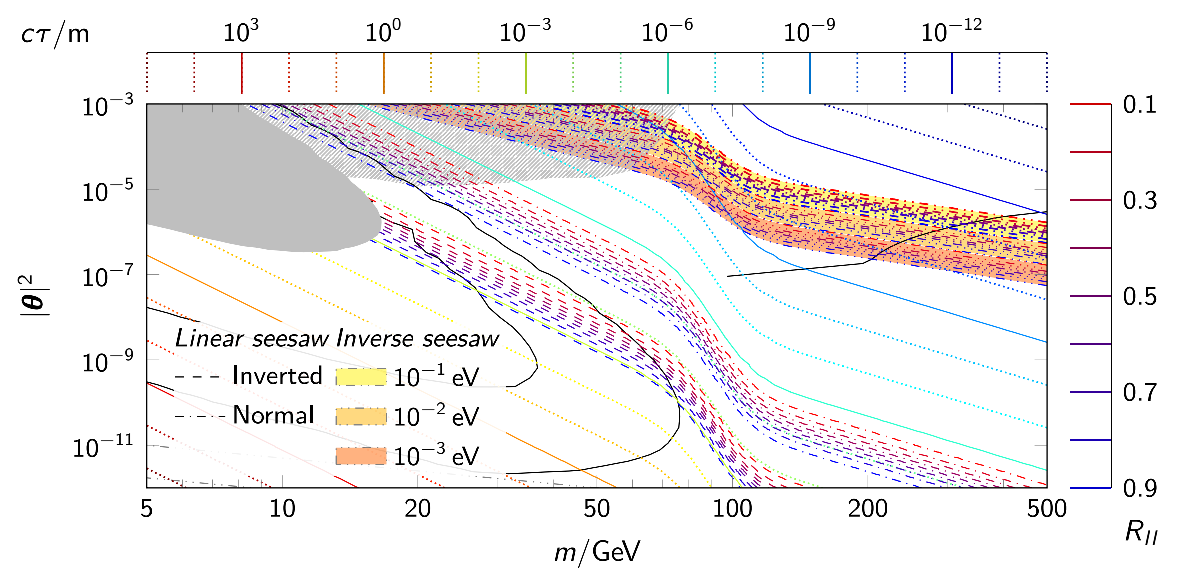
Impact of a d_T cut



LNV oscillation pattern after a d_0 cut



Reinterpretation of HNL searches as exclusion on low-scale seesaw models



Displaced searches

- Dirac HNLs good approximation when integrating over oscillations

Prompt LNV searches

- Majorana HNLs miss a factor of 2
- Model dependency governed by Δm
- Inconsequential above R_{ll} band

Conclusion

- Low-scale seesaw models predict pseudo-Dirac HNLs
- Pseudo-Dirac HNLs oscillate between LNC and LNV decays
- The symmetry protected seesaw scenario captures the relevant physics in a simple model
- We have implemented and published the necessary tools to simulate these oscillations
- Displaced HNL oscillations are resolvable at the HL-LHC
- Care has to be taken when measuring R_{ll}

References

S. Antusch, J. Hajer, and J. Rosskopp (Oct. 2022a). ‘Simulating lepton number violation induced by heavy neutrino-antineutrino oscillations at colliders’. arXiv: 2210.10738 [hep-ph]

J. Hajer and J. Rosskopp (Oct. 2022). ‘pSPSS: Phenomenological symmetry protected seesaw scenario’. FeynRules model file. DOI: 10.5281/zenodo.7268362. GitHub: heavy-neutral-leptons/pSPSS. URL: feynrules.irmp.ucl.ac.be/wiki/pSPSS

S. Antusch, J. Hajer, and J. Rosskopp (Dec. 2022b). ‘Beyond lepton number violation at the HL-LHC: Resolving heavy neutrino-antineutrino oscillations’. arXiv: 2212.00562 [hep-ph]