Probing freeze-in via invisible **Hease**

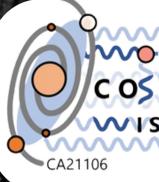
Oleg Lebedev, António P. Morais, Vinícius Oliveira, Roman Pasechnik

Aveiro University - Portugal 2nd Training School of the COST Action: Cosmic WISPers (CA21106)



universidade de aveiro theoria poiesis praxis





COS MIC

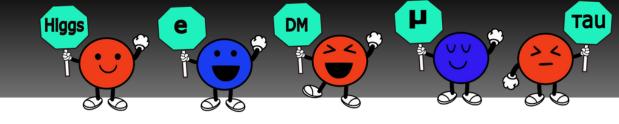


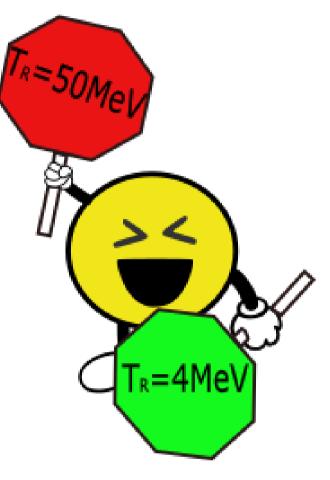
Fundacão para a Ciência Tecnologia

Goal

• Study DM production in a low-reheating temperature $T_R \sim O(4-50)MeV$ scenario;

- We are focus on Higgs portal scenario;
- We are interested in the case where $m_{DM} \gg T_R$;
- The DM is produced non-thermally due Boltzmann suppression.

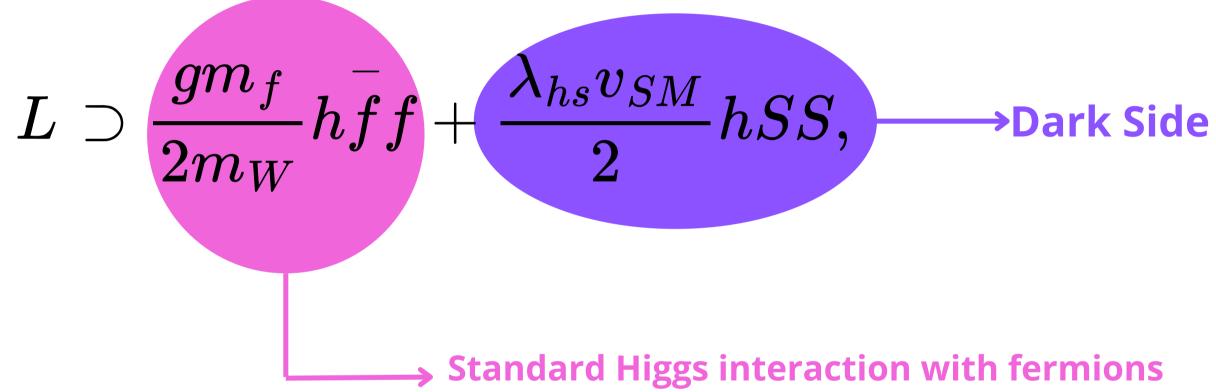






Lagrangia

• We assume that the scalar S DM candidate couples with SM particles via the Higgs portal, with the Lagrangian

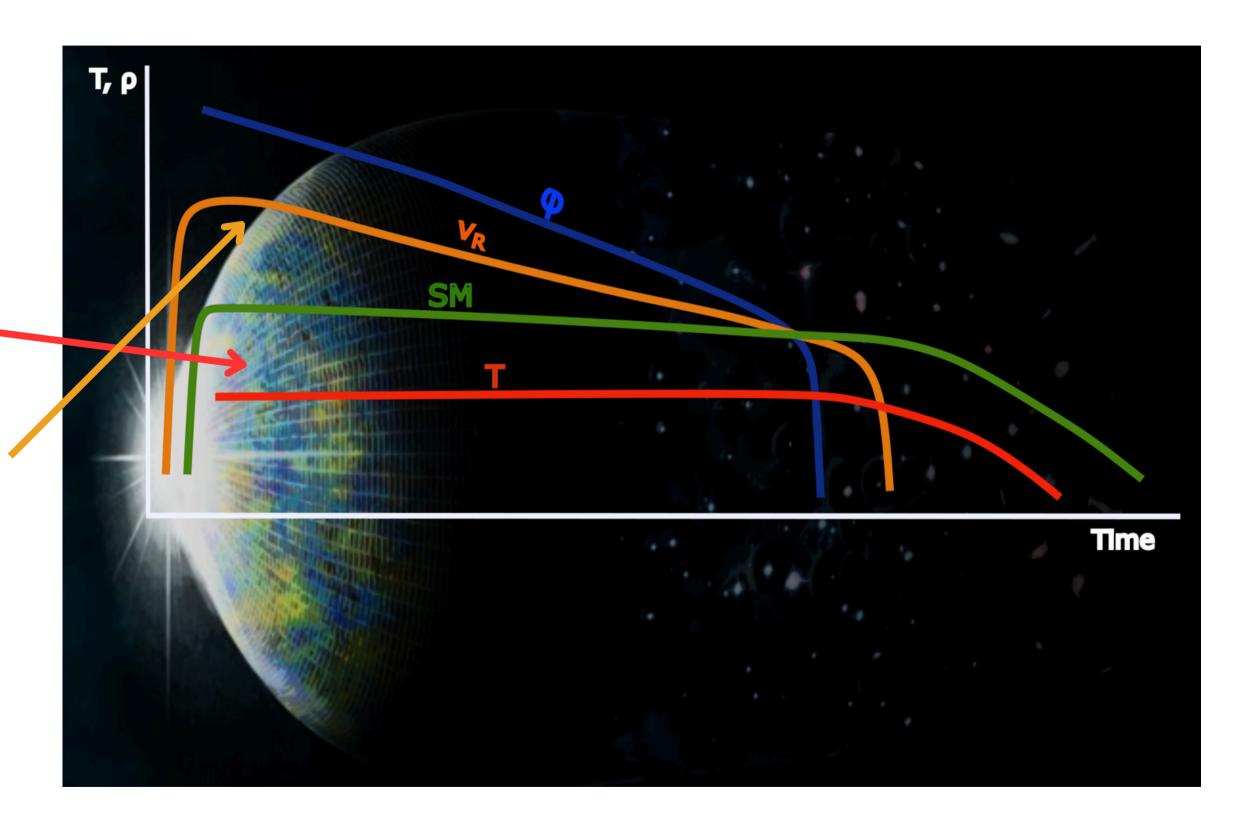


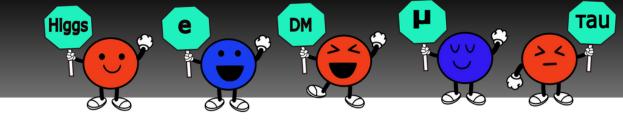
where f represents a fermionic SM particle.



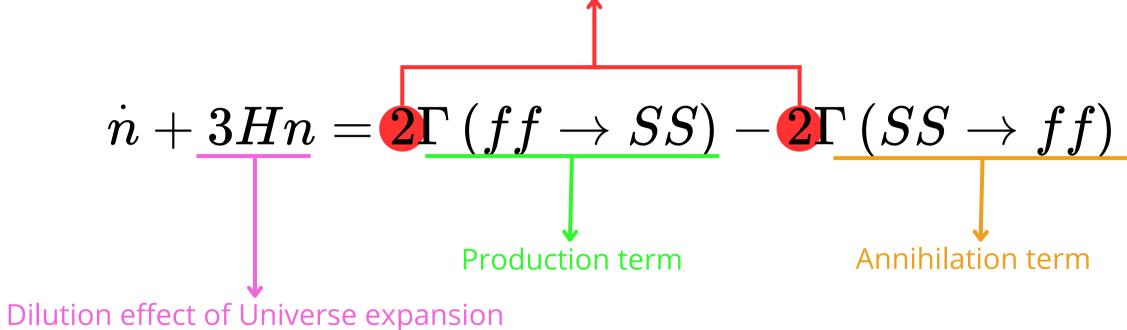
Cosmological History (Based on: arXiv2402.04743)

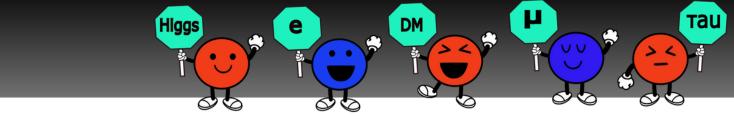
- We will assume that the maximum temperature of Universe is $T_{MAX} \sim T_R$.
- We can assume it for the scenario where the inflaton decays predominantly in other particle species (e.g. sterile neutrinos).



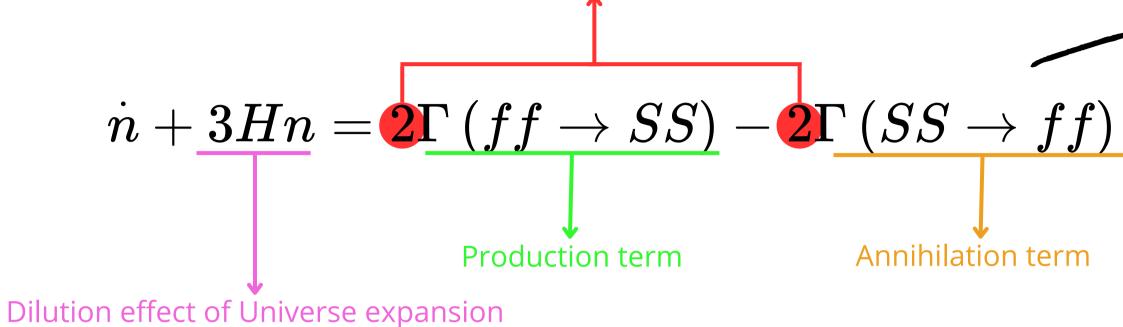


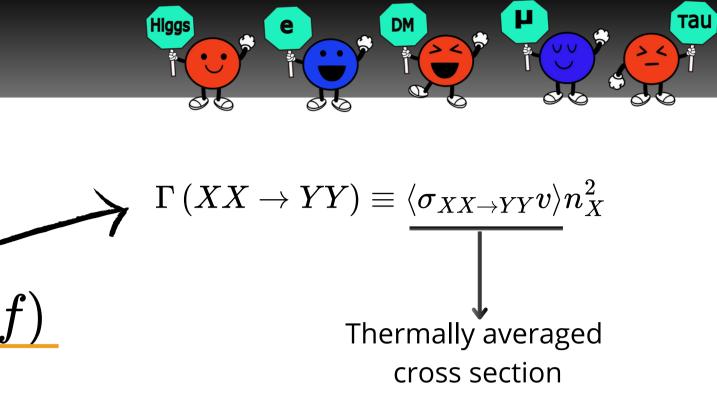
• Boltzmann equation for number density **Counts production / annihilation of 2 identical particles**





• Boltzmann equation for number density **Counts production / annihilation of 2 identical particles**



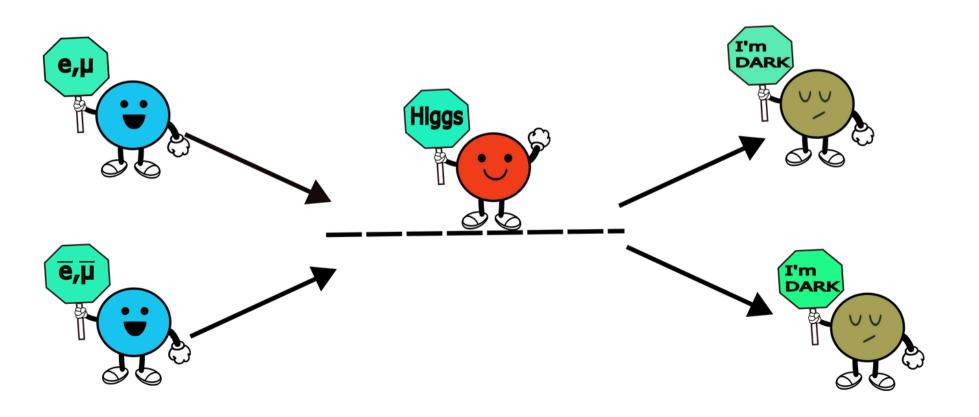


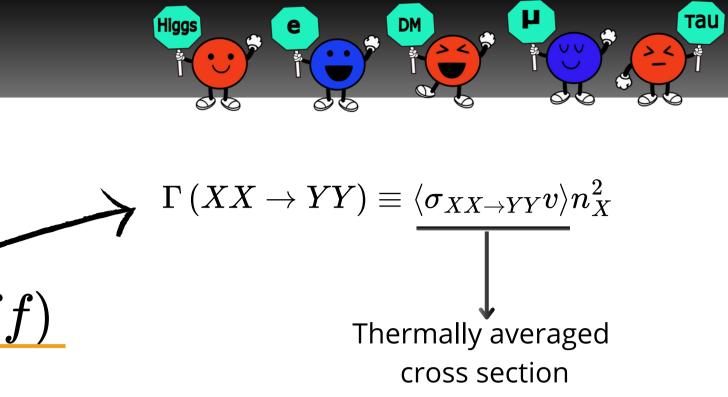
Boltzmann equation for number density
 Counts production / annihilation of 2 identical particles

$$\dot{n} + 3Hn = 2\Gamma (ff \rightarrow SS) - 2\Gamma (SS \rightarrow f)$$

Production term Annihilation ter
Dilution effect of Universe expansion

• Most important channels:





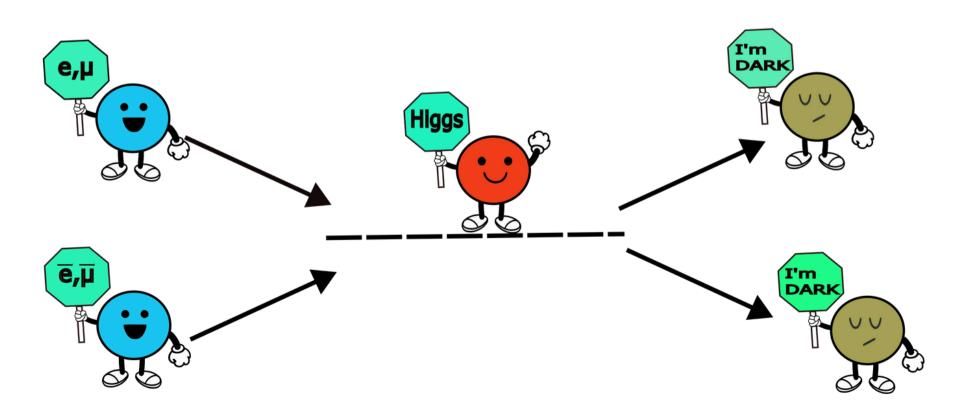
m

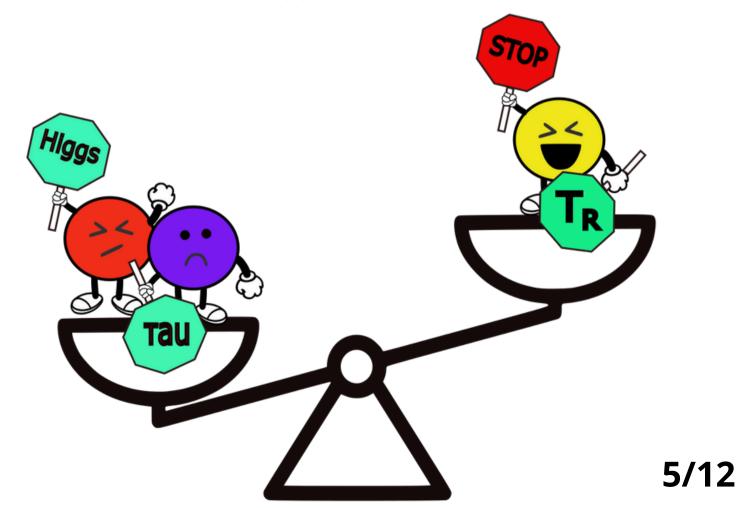
Boltzmann equation for number density
 Counts production / annihilation of 2 identical particles

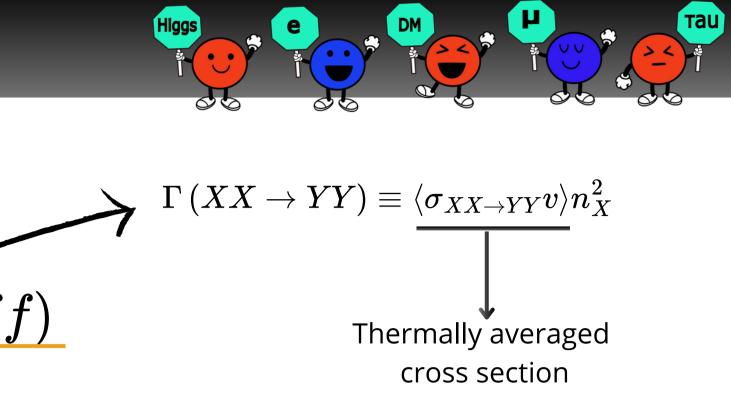
$$\dot{n} + 3Hn = 2\Gamma (ff \rightarrow SS) - 2\Gamma (SS \rightarrow f)$$
Production term
Annihilation ter
Dilution effect of Universe expansion

• Processes with $\boldsymbol{\mathcal{T}}$ and Higgs are forbidden, since $m_{ au,h} \gg T_R$.

• Most important channels:







m

• The production cross section is given by

$$\sigma_{ff \to SS} = 2 \times 4 \times \frac{1}{4} \times \frac{g^2 \lambda^2 m_f^2}{32\pi m_W^2 m_h^4} \sqrt{1 - \frac{4}{2}} + \frac{1}{32\pi m_W^2 m_H^4} \sqrt{1 - \frac{4}{32\pi m_W^2 m_H^4}} + \frac{1}{32\pi m_W^2 m_H^4} \sqrt{1 - \frac{4}{32\pi m_W^2 m_H^4}} + \frac{1}{32\pi m_W^2 m_H^4} \sqrt{1 - \frac{4}{32\pi m_W^2 m_H^4}} + \frac{1}{32\pi m_W^2 m_H^4} + \frac{1}{32\pi$$





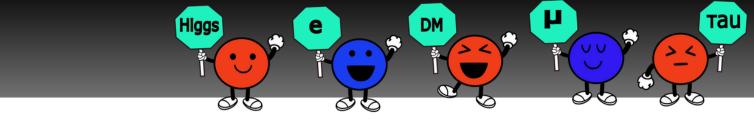
s ng

l state

• The production cross section is given by

$$\sigma_{ff \to SS} = \frac{1}{2} \times \frac{1}{4} \times \frac{g^2 \lambda^2 m_f^2}{32\pi m_W^2 m_h^4} \sqrt{1 - \frac{4}{4}}$$
• We can estimate the rate

$$\Gamma_e \approx \frac{g^2 \lambda^2 m_s^3 m_e^2 T^3}{256\pi^4 m_h^4 m_W^2} e^{-\frac{2m_s}{T}} \qquad \Gamma_\mu \approx \frac{3g^2 \lambda^2 m_W^2 m_h^4}{512\pi^4 m_W^4}$$



 $4m_f^2$

 \boldsymbol{S}

ng l state $m_{\mu}^2 T^4$ $2m\mu$ $\mathbf{2}$

T

 $n_h^4 m$

• The production cross section is given by

$$\sigma_{ff \to SS} = \frac{1}{2} \times \frac{1}{4} \times \frac{g^2 \lambda^2 m_f^2}{32\pi m_W^2 m_h^4} \sqrt{1 - \frac{4}{4}}$$
• We can estimate the rate

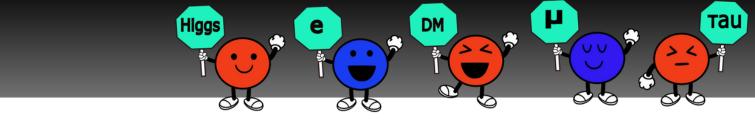
$$\Gamma_e \approx \frac{g^2 \lambda^2 m_s^3 m_e^2 T^3}{256\pi^4 m_h^4 m_W^2} e^{-\frac{2m_s}{T}} \qquad \Gamma_\mu \approx \frac{3g^2 \lambda^2 r}{512\pi^4 m_H^4}$$

• The ratio of the muon and electron rates for $T = T_R$ is

$$rac{\Gamma_\mu}{\Gamma_e}pprox rac{3}{2}rac{m_\mu^4 T_R}{m_e^2 m_s^3}e^{-2(m_\mu-m_s)/T_R}$$

for $m_S = 50 MeV$ we obtain that Γ_{μ} dominate the contribution for

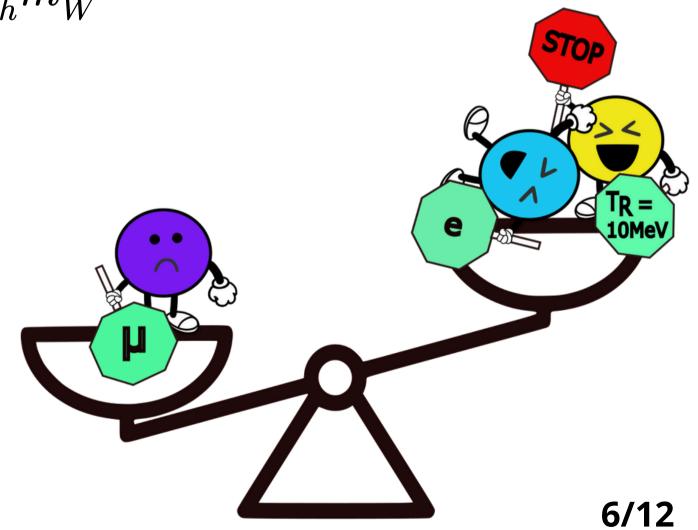
 $T_R > 10 MeV$



 $4m_f^2$

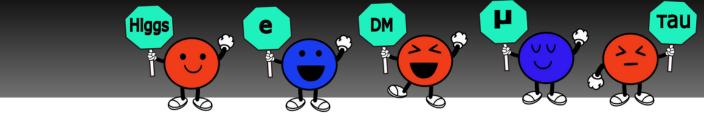
 \boldsymbol{S}

ng l state



• For the *pure-freezing* scenario $\Gamma(SS \rightarrow ff) = 0$

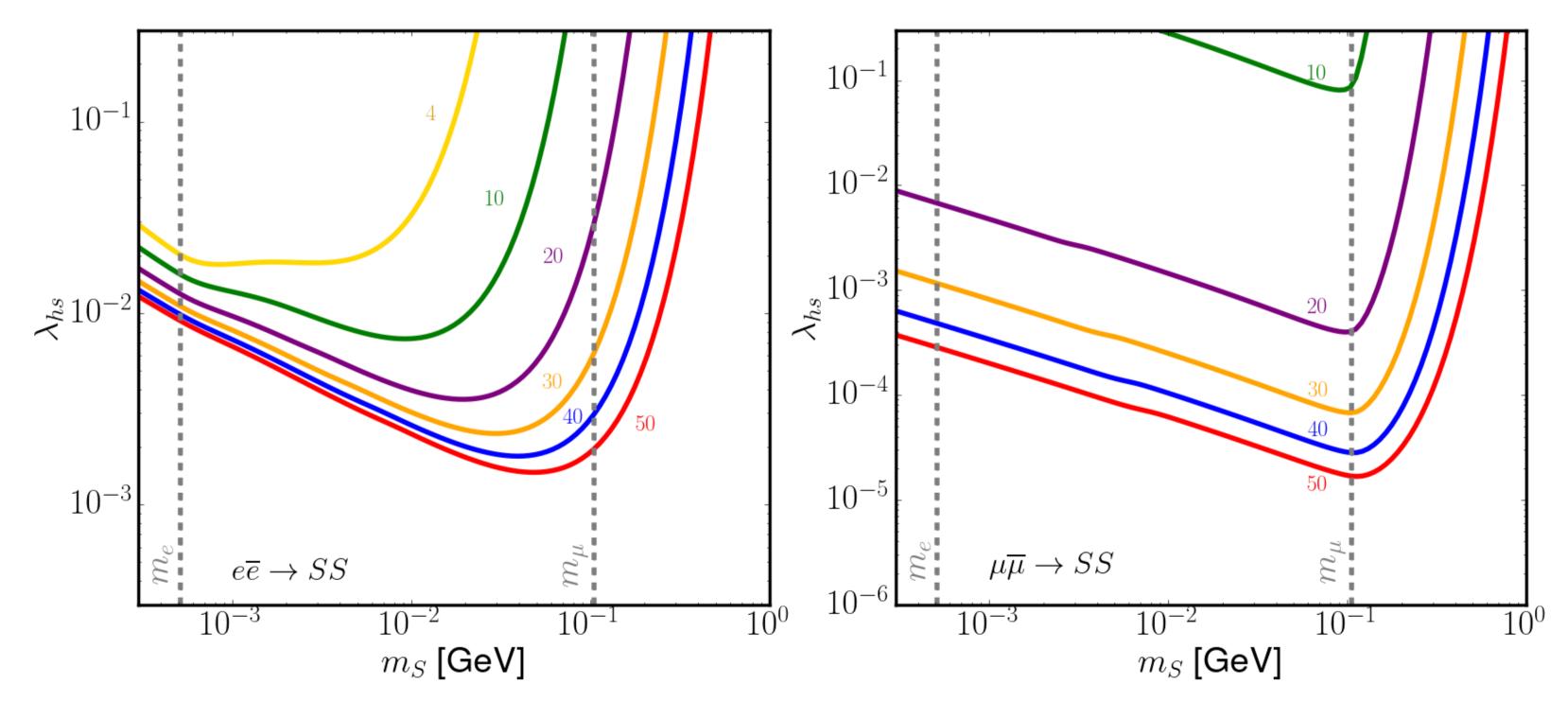
$$rac{dY_S}{dx} = 2 \sqrt{rac{8 \pi^2 M_{Pl}^2}{45}} rac{g_*^{1/2} m_s}{x^2} \sum_{f=e,\mu} \langle \sigma_{ff
ightarrow S}$$



 $_{
ightarrow SS} v
angle Y_{f}^{(eq)2}$

• For the *pure-freezing* scenario $\Gamma (SS \rightarrow ff) = 0$

$$rac{dY_S}{dx} = 2 \sqrt{rac{8 \pi^2 M_{Pl}^2}{45}} rac{g_*^{1/2} m_s}{x^2} \sum_{f=e,\mu} \langle \sigma_{ff
ightarrow S}$$





 $_{
m a}{}_{SS}v
angle Y_{f}^{(eq)2}$

Constraints

• Invisible Higgs decay

$$BR_{inv} = rac{\Gamma_{h
ightarrow ss}}{\Gamma_{SM}^{Tot} + \Gamma_{h
ightarrow ss}}$$

where

$$\Gamma_{h
ightarrow ss} = rac{\lambda^2}{32\pi m_h} \sqrt{1-rac{4m_s^2}{m_h^2}}$$

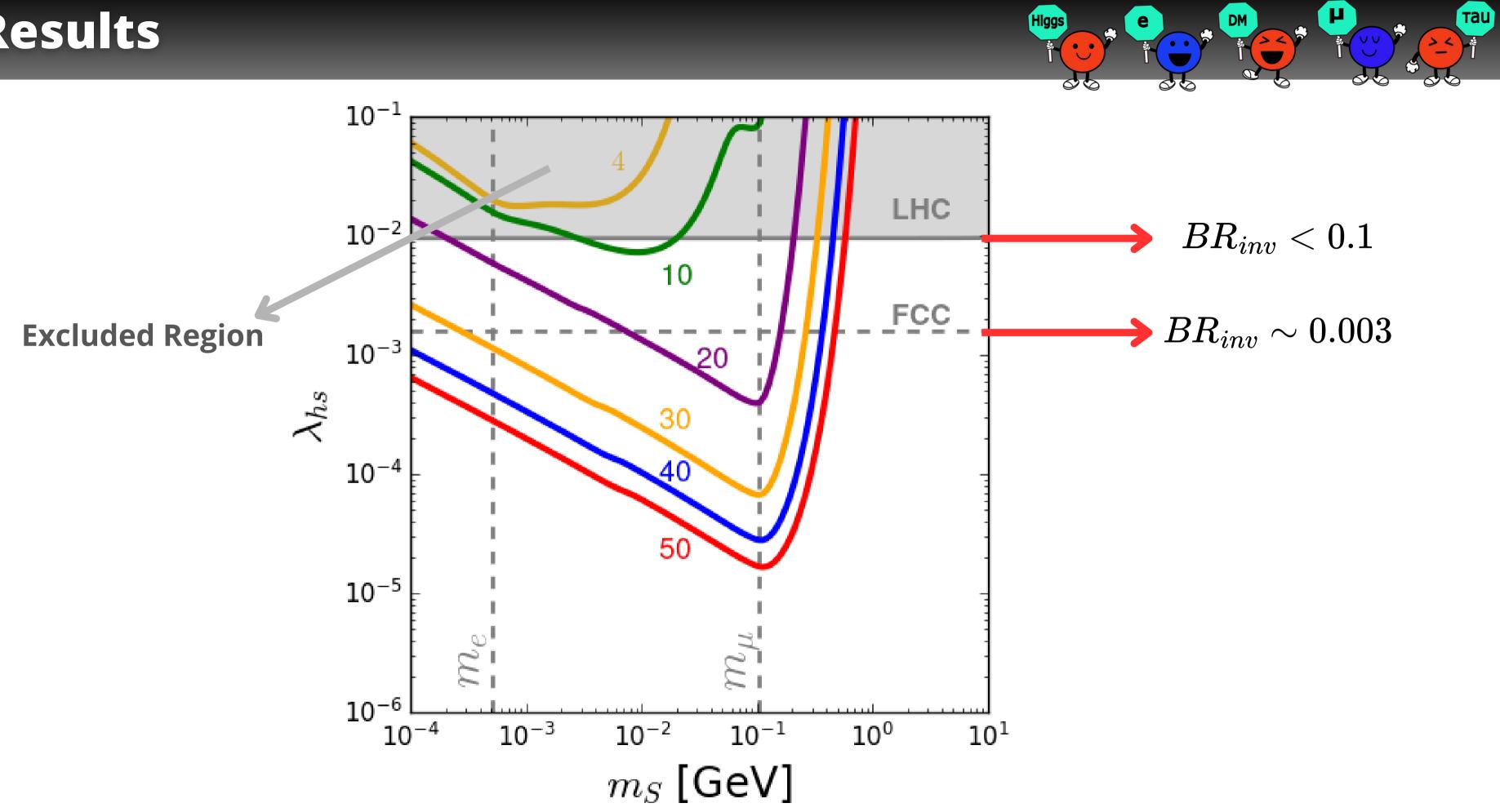
• We impose $BR_{inv} < 0.1$.

• The future collider FCC tuned as a Higgs factory can reach the limit of

$$BR_{inv}\sim 0.003$$



Results



In progress...

• We will also consider the case of fermionic DM

$$L \supset rac{gm_f}{2m_W} har{f}f + rac{\lambda_{hs} v_{SM}}{2} h\chi ar{\chi}, \hspace{0.5cm} L \supset rac{gm}{2m_W}$$



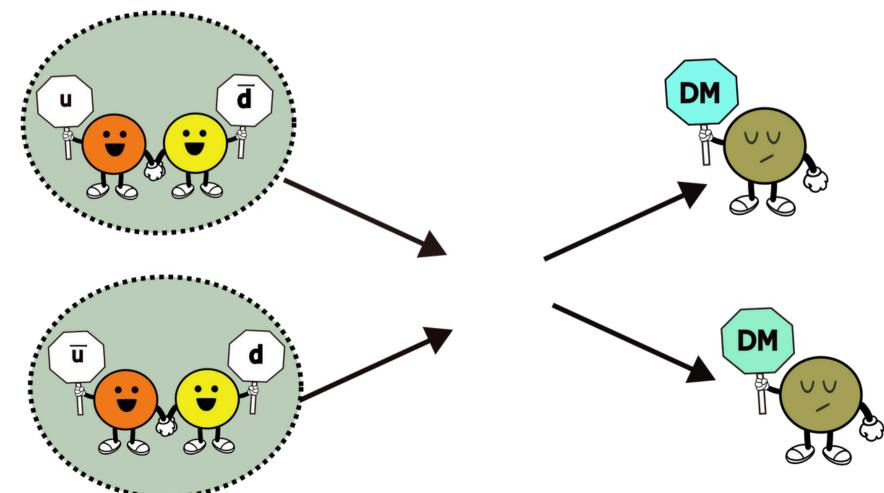
 $rac{n_f}{n_W} har{f}f + rac{\lambda_{hs} v_{SM}}{2} ih\chi\gamma^5 \chi$

In progress...

• We will also consider the case of fermionic DM

$$L \supset rac{gm_f}{2m_W} har{f}f + rac{\lambda_{hs} v_{SM}}{2} h\chi ar{\chi}, \hspace{0.5cm} L \supset rac{gm_f}{2m_W}$$

• We will also consider the mesons contributions





 ${n_f\over n_W} h {ar f} f + {\lambda_{hs} v_{SM}\over 2} ih \chi \gamma$

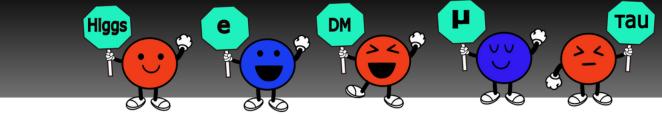
Conclusions

• Freeze-in Mechanism:

- Dark Matter is Boltzmann suppressed if its mass exceeds the reheating temperature.
- Correct DM abundance achieved with lower mass at low reheating temperatures.

• Viability and Implications:

- Scenarios with low reheating temperatures are viable.
- Future Implications
 - Significant contribution from mesons after the Quantum Chromodynamics (QCD) phase transition.



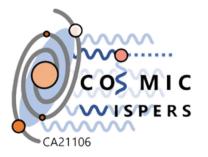


The Early Universe PUB ••• 2/24 J 0 • • Ve 3











SCAN ME



Fundação para a Ciência e a Tecnologia