

# Binary pulsars as fuzzy dark matter detectors

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# Team



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# Abstract

Ultra-light dark matter (ULDM) perturbs the orbital motion of binary pulsars, inducing variations in the pulses' times-of-arrival. Binary pulsars are thus promising detectors of ULDM.

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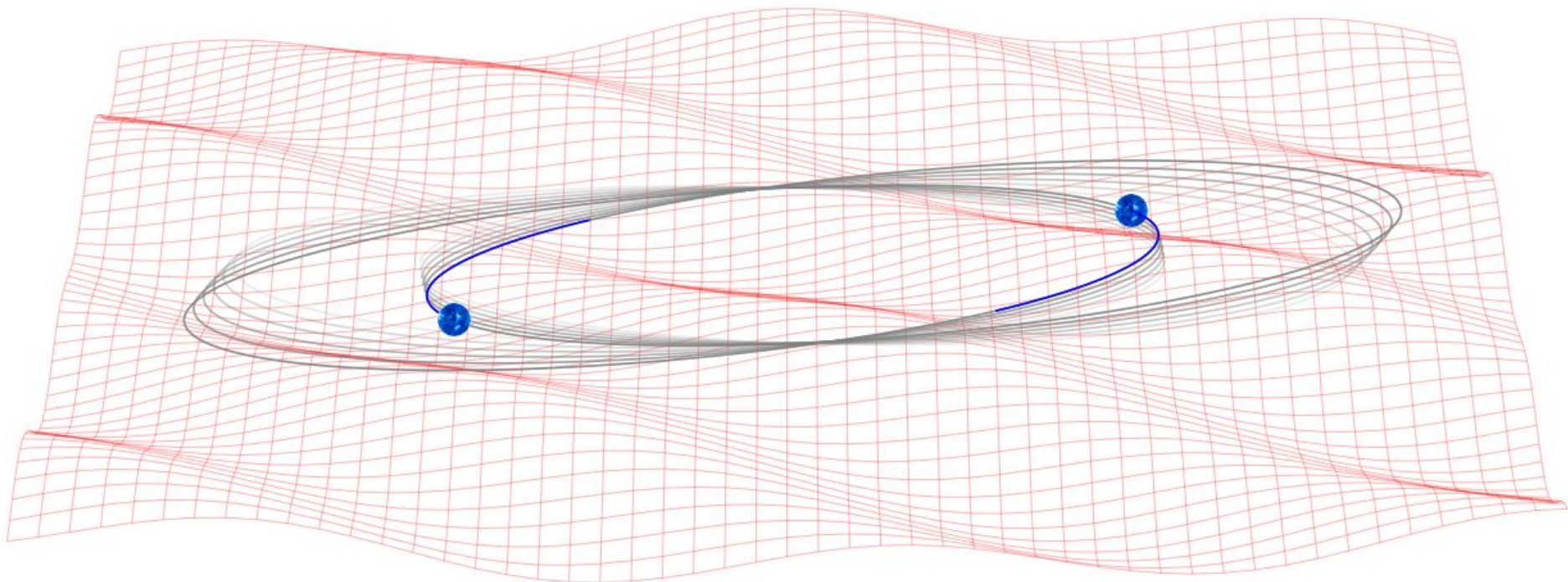


Figure comes from <https://arxiv.org/pdf/2107.04063.pdf>

# 1. Dark matter interaction

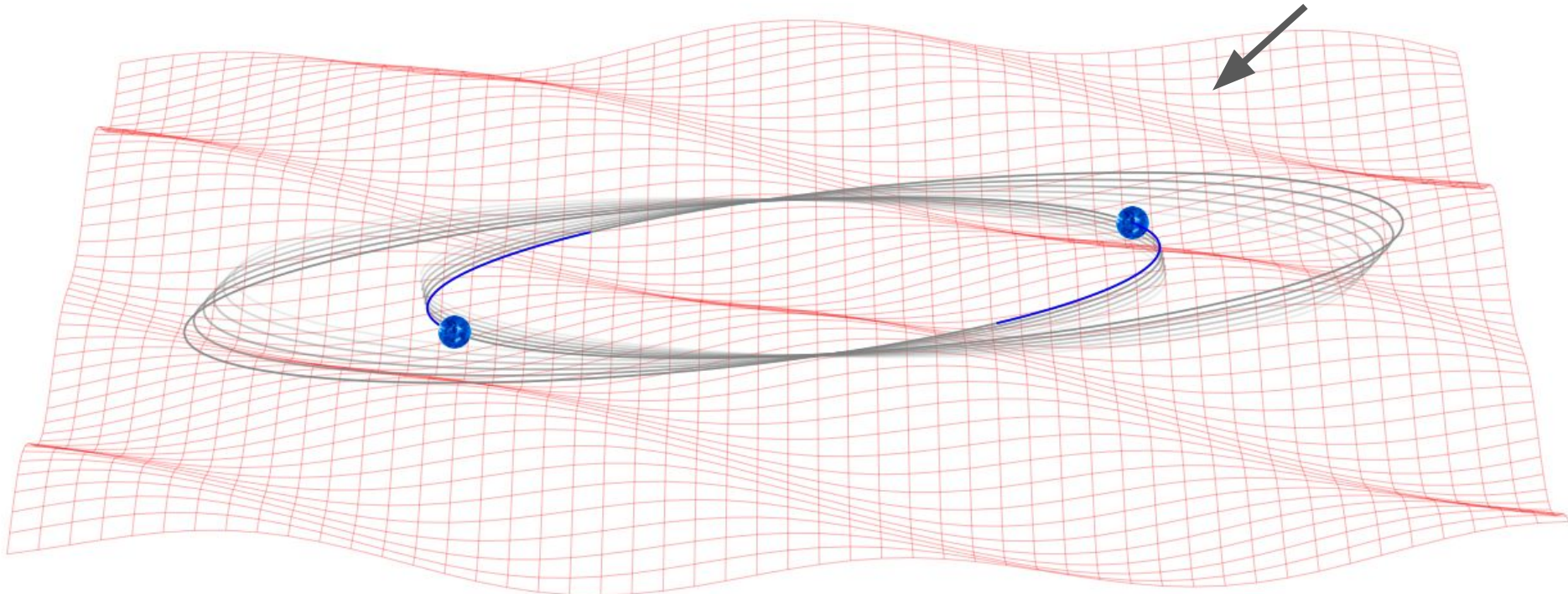
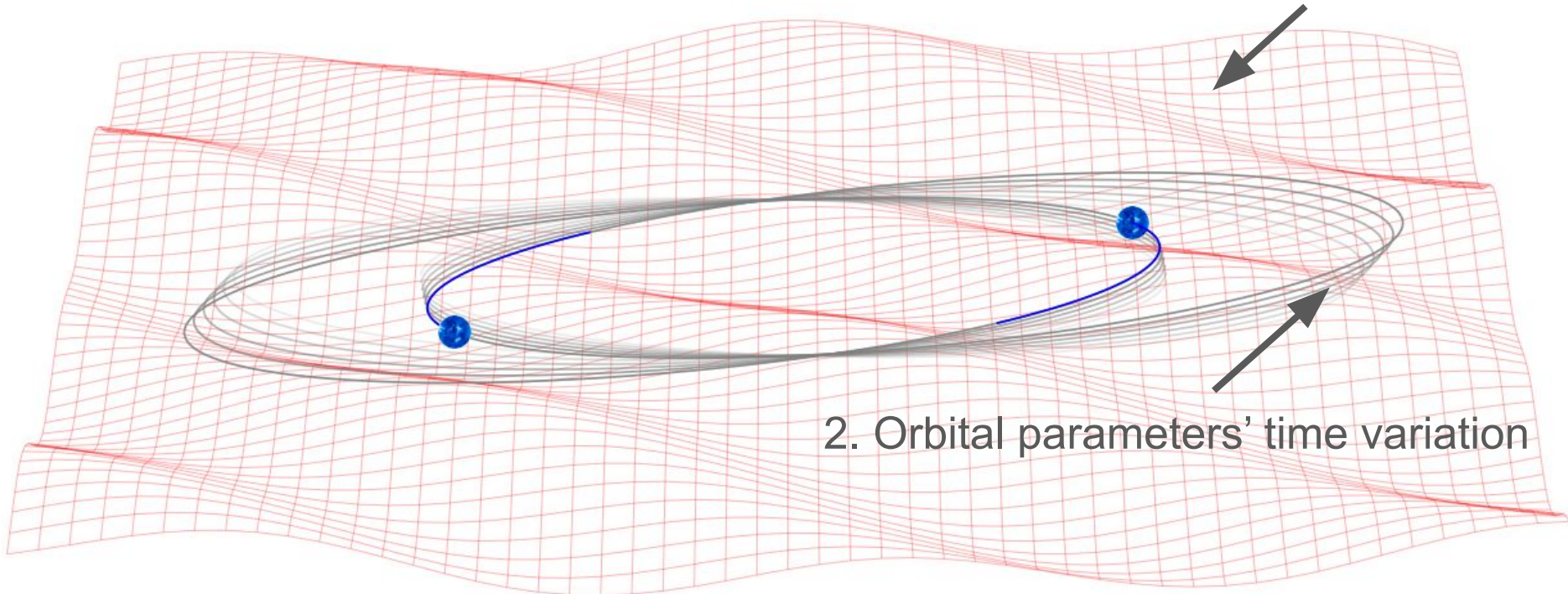


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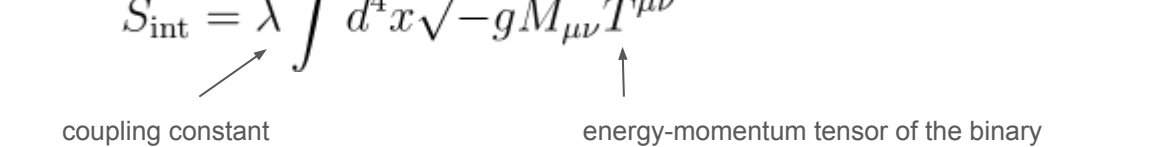


2. Orbital parameters' time variation

# ULDM models and Binaries

Interaction: gravitational and/or **direct**

Probed masses:  $m = 10^{-23} \text{ eV} \div 10^{-18} \text{ eV}$  (back-reaction negligible)

↗	Spin 0 $\Phi$	$M_{1,2}^{(\alpha)}(\Phi) = M_{1,2} (1 + \alpha \Phi)$ or $M_{1,2}^{(\beta)}(\Phi) = M_{1,2} \left( 1 + \frac{\beta}{2} \Phi^2 \right)$	arXiv:1612.06789, arXiv:1910.08544
→	Spin 1 $\vec{A}$	$L_{\text{int}} = q_1 \dot{\vec{r}}_1 \cdot \vec{A} + q_2 \dot{\vec{r}}_2 \cdot \vec{A},$	arXiv:1807.10491
↘	Spin 2 $M_{\mu\nu}$	$S_{\text{int}} = \lambda \int d^4x \sqrt{-g} M_{\mu\nu} T^{\mu\nu}$ 	arXiv:1909.13814



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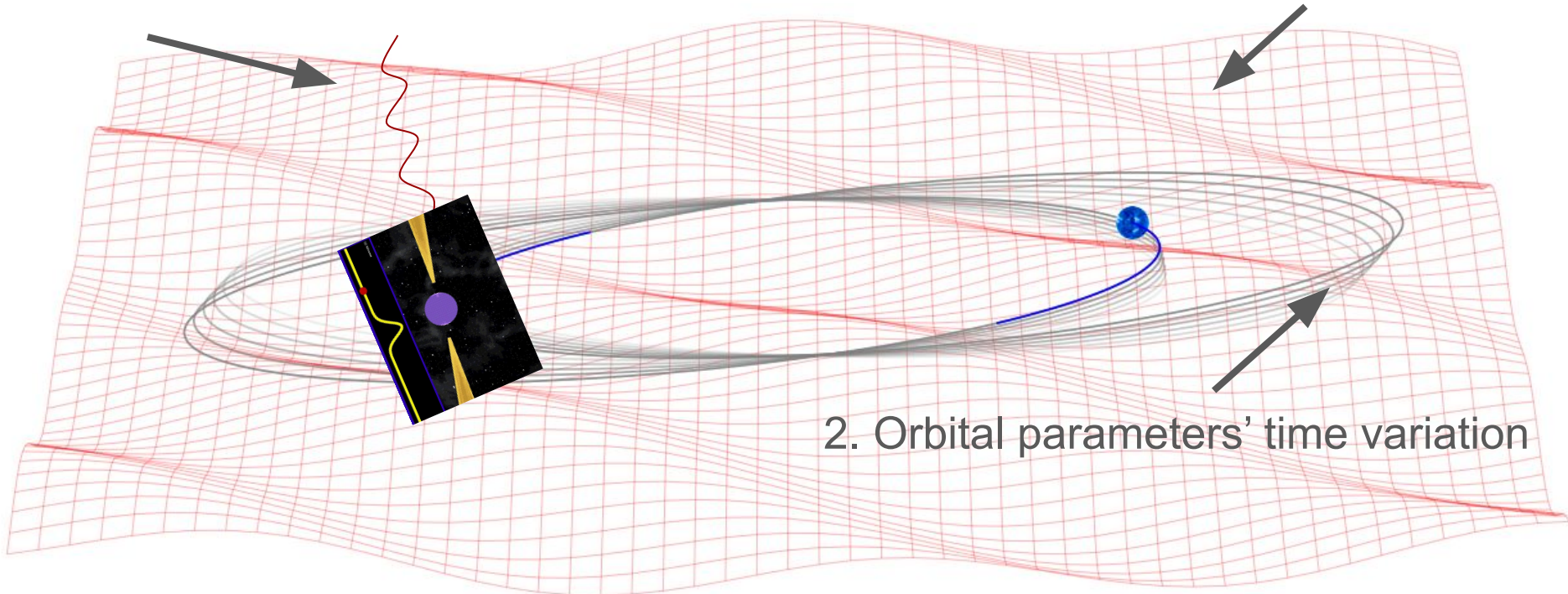
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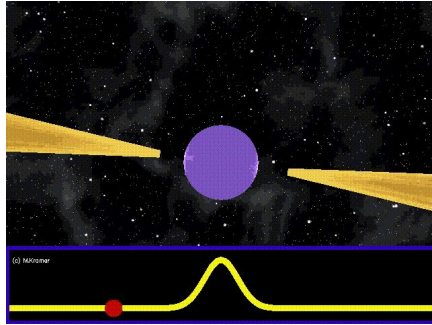
3. Pulse emission  
(imprint of 1. through 2.)

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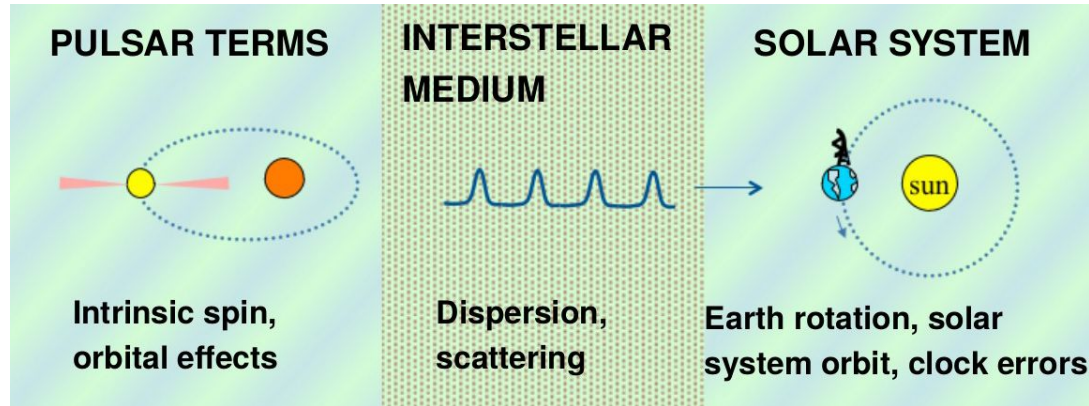
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# Pulsar timing



pulse propagation



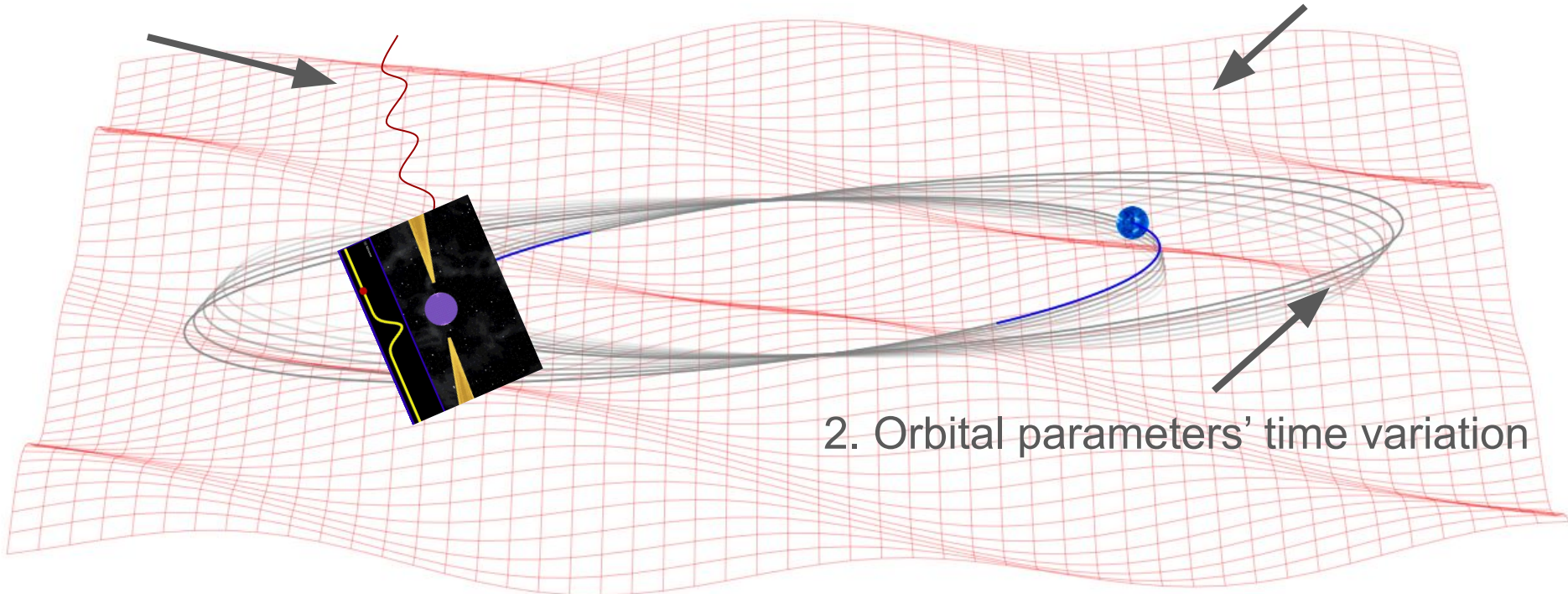
Timing model = description of the propagation of pulses from the pulsar to the Earth

“*binary pulsar*” = binary star with a pulsar as one component

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# Bayes' Theorem

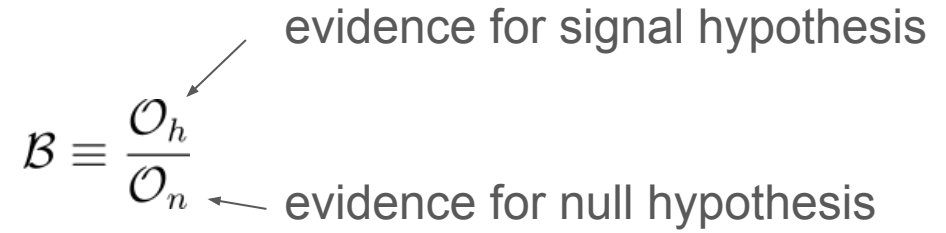
$$\overbrace{P(\theta | y)}^{\text{Posterior}} \propto \prod_i^N \underbrace{P(y_i | \theta)}_{\text{Data Likelihood}} \overbrace{P(\theta)}^{\text{Prior}}$$

# A bit of theory

$$\mathcal{B} \equiv \frac{\mathcal{O}_h}{\mathcal{O}_n}$$

evidence for signal hypothesis

evidence for null hypothesis



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$$\mathcal{O}_n = \int \mathcal{L}(\Xi) d\vec{\Xi}$$

nuisance parameters (errors, secular effects)

$$\mathcal{O}_h = \int P(X, Y) P(\alpha, m) \mathcal{L}(\Xi, X, Y, \alpha, m) d\vec{\Xi} dX dY d\alpha dm, P(\alpha, m) = \delta(\alpha - \alpha_f) \delta(m - m_f).$$

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$$\bar{\mathcal{B}} \equiv \int d\mathbf{n} P(\mathbf{n}) \mathcal{B}$$

averaging over noise

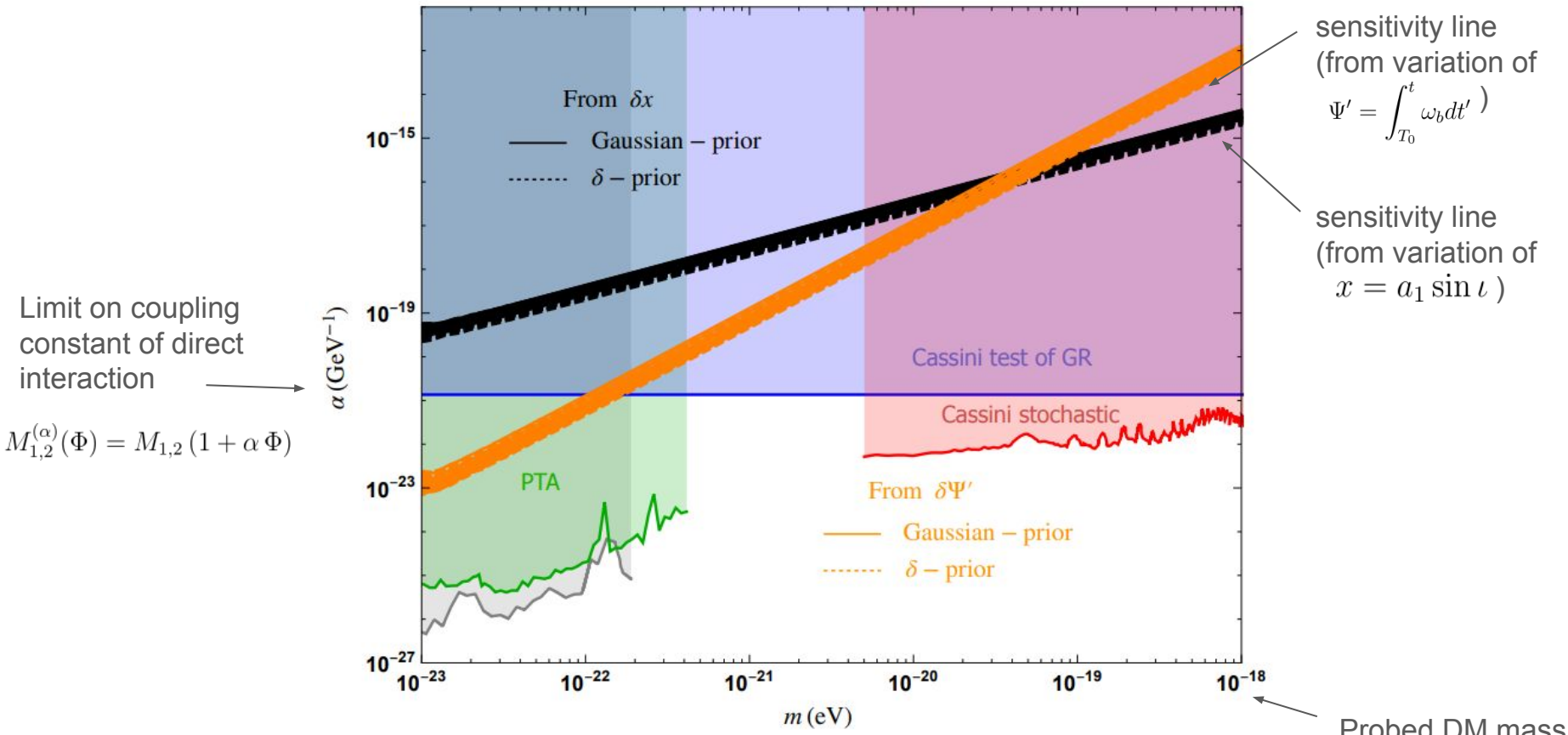
$$\bar{\mathcal{B}}^C = \prod_{p=1}^{N_p} \bar{\mathcal{B}}_p = 1000 \quad (\text{depends on } \alpha)$$

combination of all pulsars

# Abstract

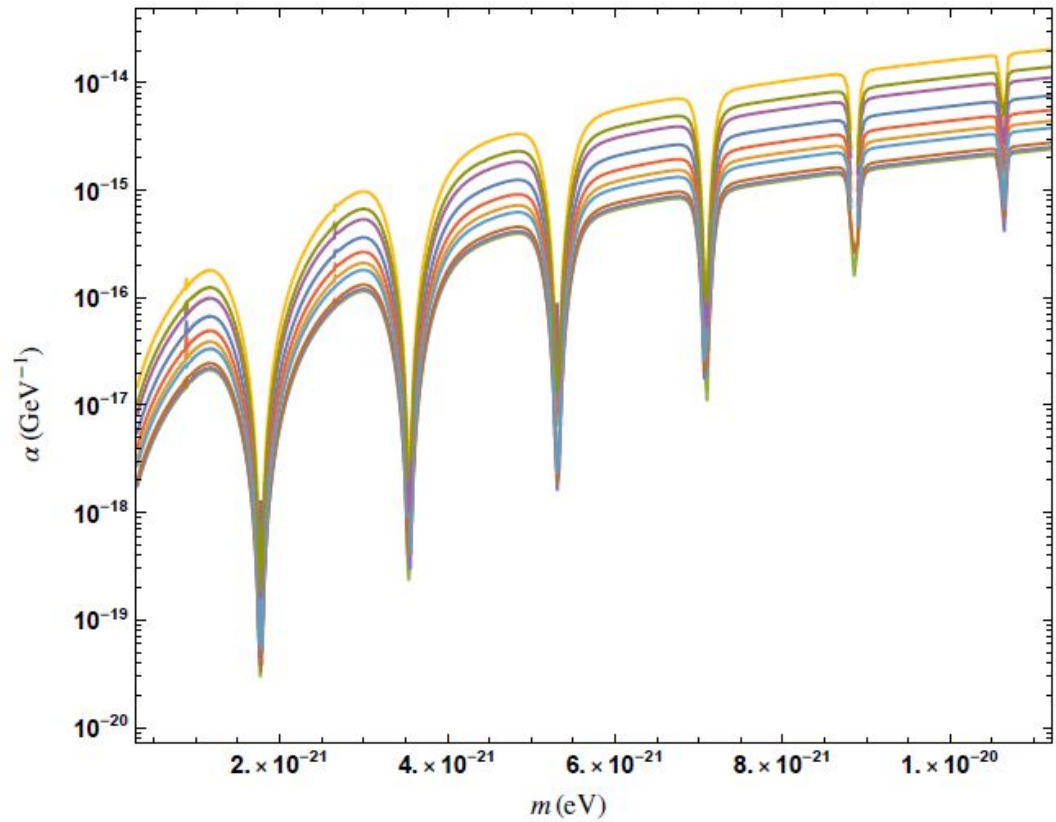
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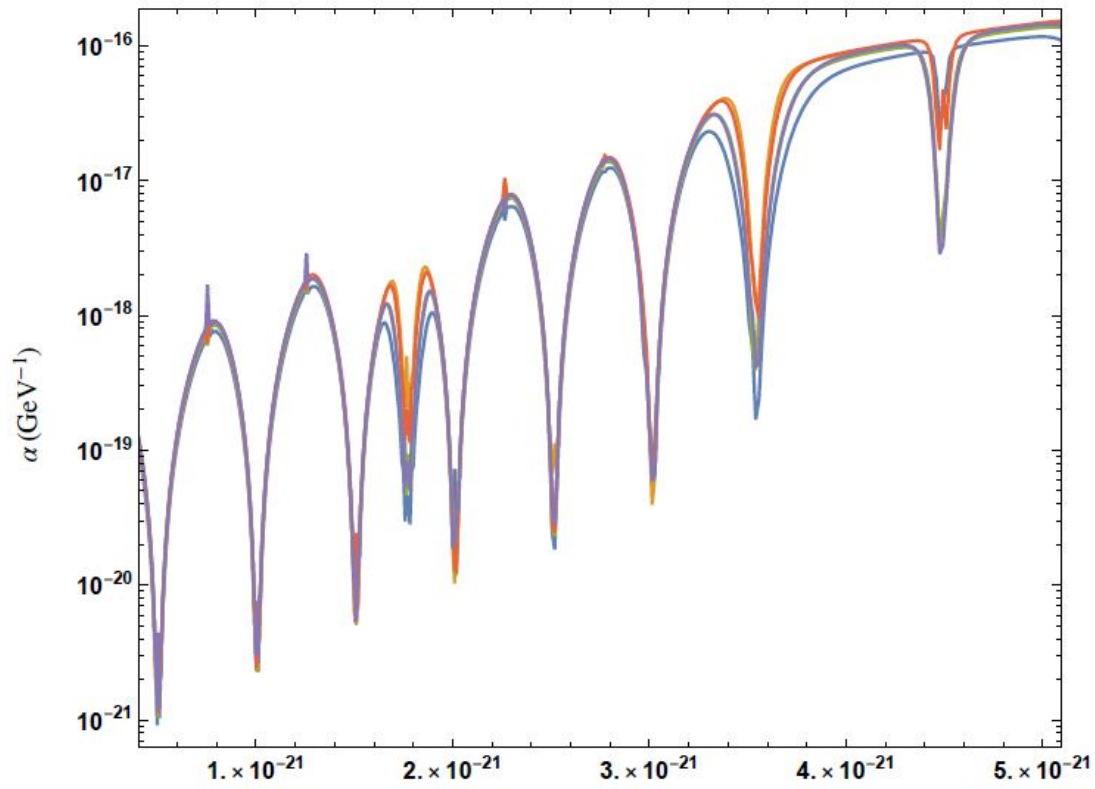


111 binary pulsars from the ATNF catalogue and assuming next-generation radio-telescope precision (future SKA)





J1946+3417, NANOGrav



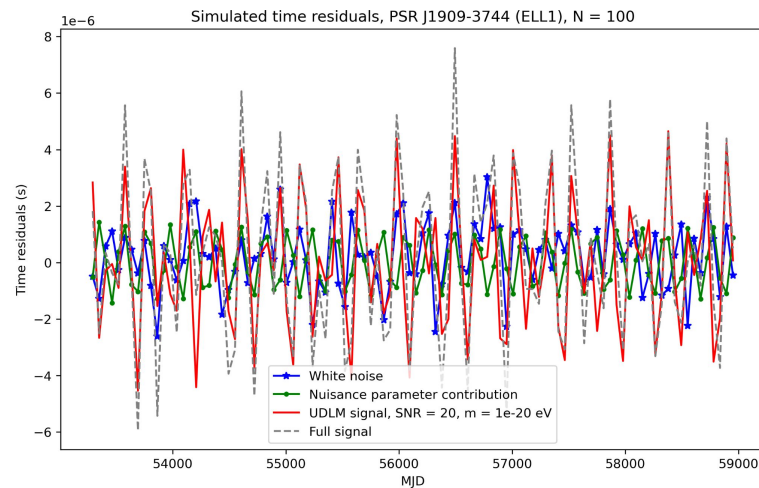
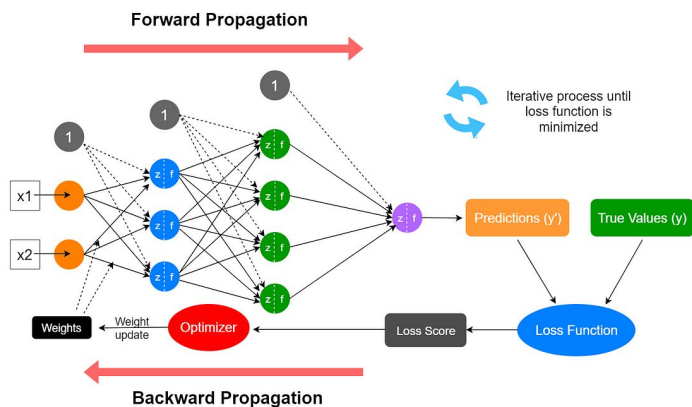
J1946+3417, J2234+0611, J1903+0327, NANOGrav

# Summary

1. binary pulsars are *UDLM detectors*
2. new method → *Bayesian sensitivity lines*
3. applied to *scalar ULDM*, large portion of DM *phase space constrained*
4. *forecast* of the constraining power of data, not an actual data analysis
5. each orbital parameter has a *different constraining power* (main  $x, \Psi'$ )
6. ELL1 binaries — even more restrictive than eccentric (BT) systems!

# Outlook

1. Generalization to spin 1 and 2
2. Application of Machine learning to search for UDLM signals in data



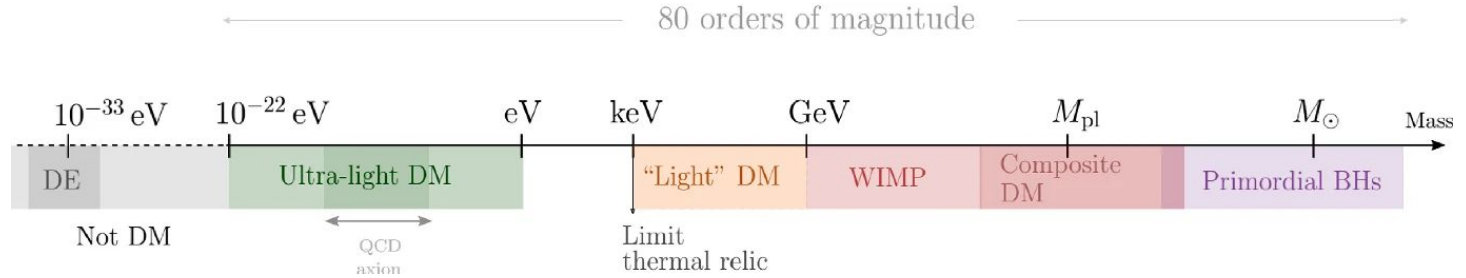
Thank you for your attention!

Pavel Kůs  
[pavel.kus@fzu.cz](mailto:pavel.kus@fzu.cz)

*Backup slides*



# ULDM



## Dark matter particle

$$m = 10^{-23} \text{ eV} \div 1 \text{ eV}$$

very light bosons

de Broglie wavelength of astrophysical scale

$$\lambda_{\text{dB}} \sim 1.3 \times 10^{12} \text{ km} \left( \frac{10^{-3}}{V_0} \right) \left( \frac{10^{-18} \text{ eV}}{m} \right)$$

↑  
typical DM velocity in the  
Milky Way halo

## Within the halo of typical galaxy

classical field theory description applicable

Schrödinger–Poisson system of eqs

wave-like behavior, interference patterns

$$\Phi = \Phi_0 \varrho \cos(mt + \Upsilon) \quad \Phi_0 \equiv \frac{\sqrt{2\rho_{\text{DM}}}}{m}$$

↑ Local DM amplitude / ↑ Local DM phase

↑ DM mass (= oscillation frequency)

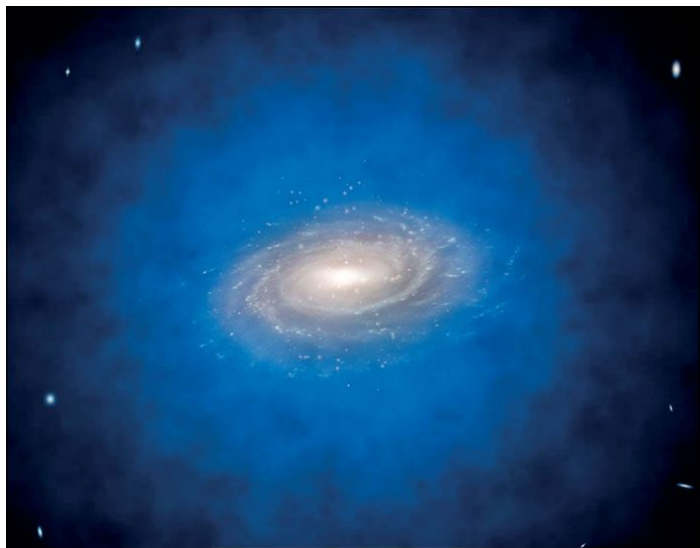


Granular structure  
→  
ULDM

$\varrho_7, \Upsilon_7$	$\varrho_8, \Upsilon_8$	$\varrho_9, \Upsilon_9$
$\varrho_4, \Upsilon_4$	$\varrho_5, \Upsilon_5$	$\varrho_6, \Upsilon_6$
$\varrho_1, \Upsilon_1$	$\varrho_2, \Upsilon_2$	$\varrho_3, \Upsilon_3$

$\sim \lambda_{dB}/2$

$$P(\varrho) = 2\varrho e^{-\varrho^2}, \quad P(\Upsilon) = \frac{1}{2\pi}$$



Granular structure  
 →  
 ULDM

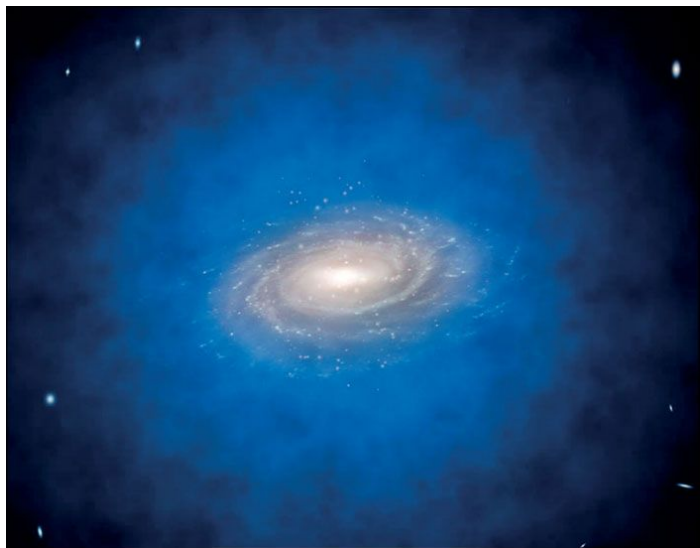
ULDM homogeneous within one patch

$\varrho_7, \Upsilon_7$	$\varrho_8, \Upsilon_8$	$\varrho_9, \Upsilon_9$
$\varrho_4, \Upsilon_4$	$\varrho_5, \Upsilon_5$	$\varrho_6, \Upsilon_6$
$\varrho_1, \Upsilon_1$	$\varrho_2, \Upsilon_2$	$\varrho_3, \Upsilon_3$

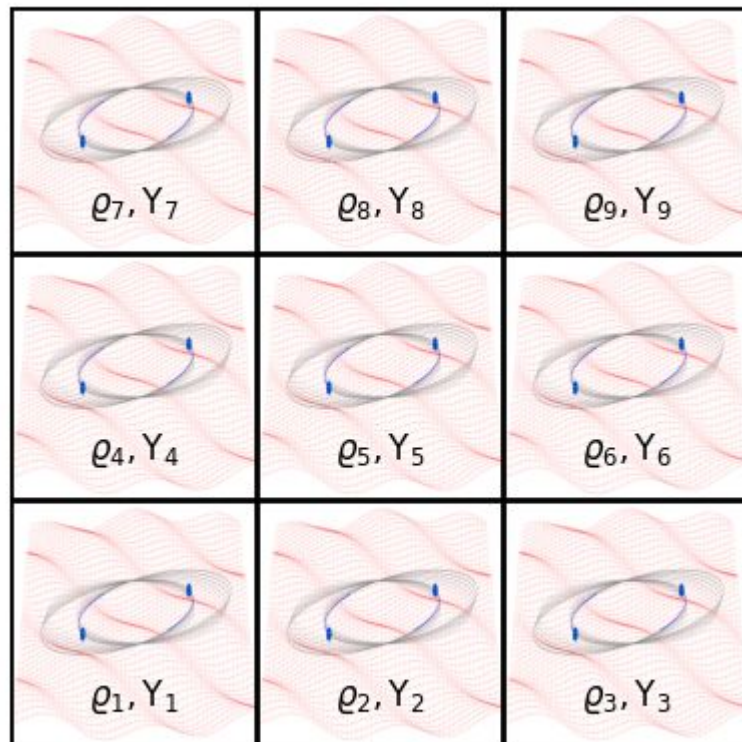
← O(1) fluctuations

$\sim \lambda_{dB}/2$

$$P(\varrho) = 2\varrho e^{-\varrho^2}, \quad P(\Upsilon) = \frac{1}{2\pi}$$



Granular structure  
→  
ULDM



$\sim \lambda_{dB}/2$

Binary size  $\ll \lambda_{dB}$

### 3 key assumptions

1. orbital resonance on  $P_b$

$$m = k\omega_b, k \in \mathbb{N}$$

2.  $e \gg 0$

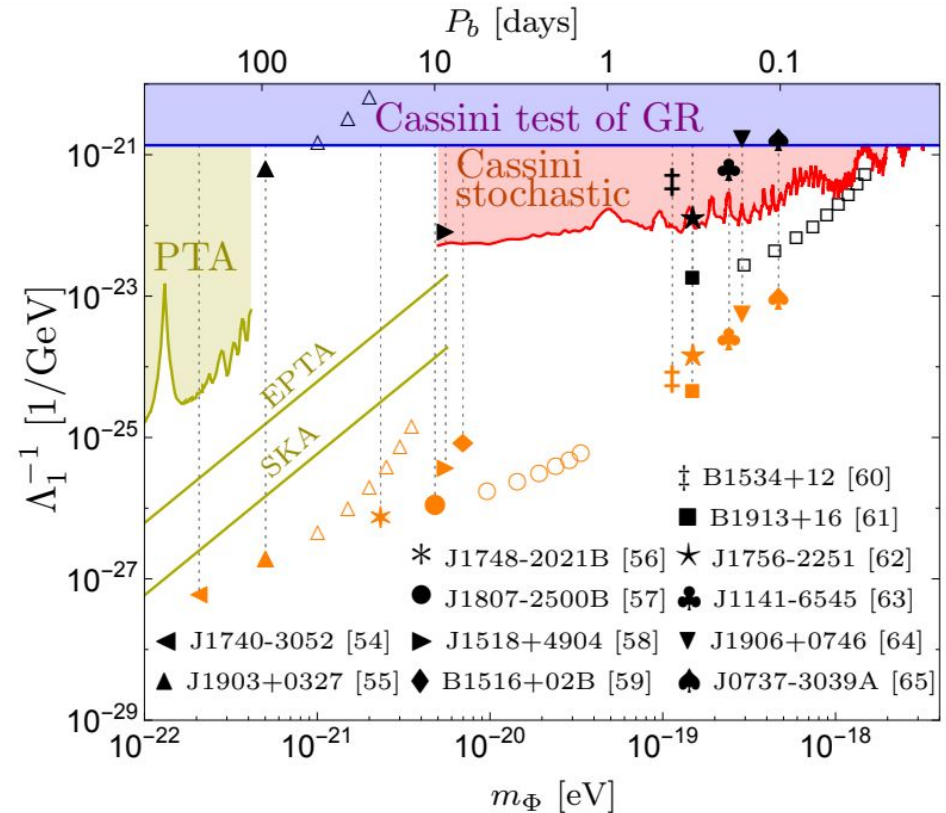
small subset of all pulsars

3.  $\langle (\dot{P}_b)_{\text{UDLM}} \rangle \leq (\dot{P}_b)_{\text{error}}$

ignoring the time behaviour  
of the ULDM perturbations

### New method:

1. beyond resonance (all masses)
2. any eccentricity
3. inclusion of the time behaviour of ULDM perturbations
4. combining all pulsars together to form one global sensitivity line.



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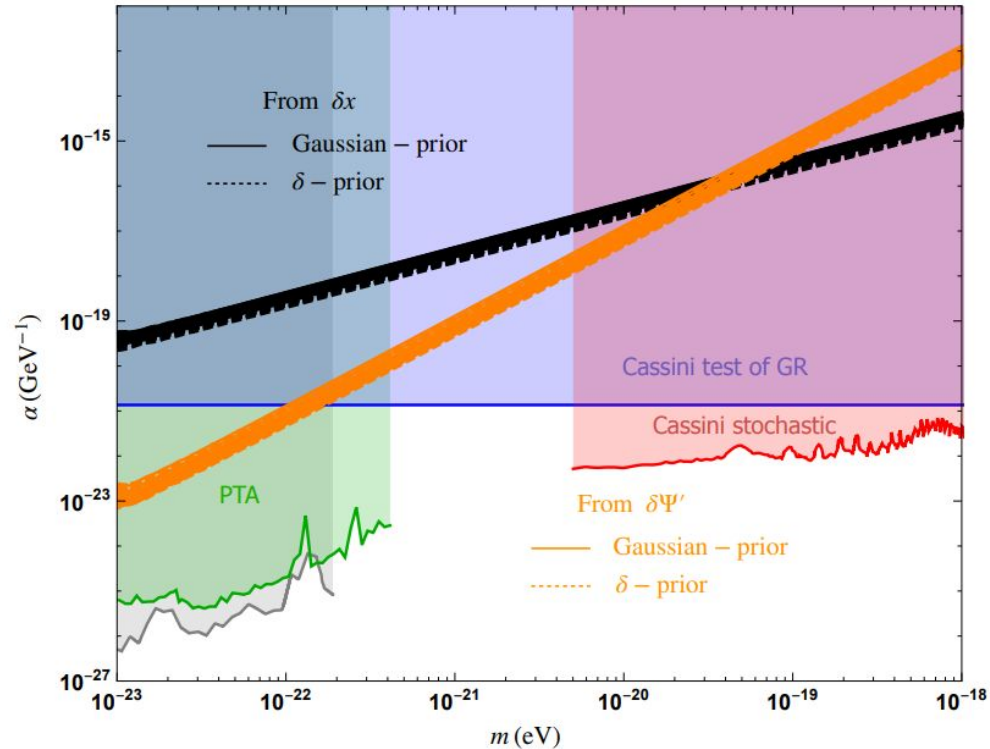
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# Pulsar timing

$$N = N_0 + \nu T + \frac{1}{2} \frac{d\nu}{dT} T^2 + \frac{1}{6} \frac{d^2\nu}{dT^2} T^3$$

$T$  ... proper time  
 $t$  ... infinite-frequency barycenter arrival time  
 $T(t)$  ... timing model

$$N = N_0 + \nu t + \frac{\dot{\nu}}{2} t^2 + \frac{\ddot{\nu}}{6} t^3 - \dot{\nu} x g(\Psi') t - \nu g(\Psi') + \nu \omega_b g(z) \frac{dg(z)}{dz} \Big|_{z=\Psi'},$$

where  $g(z) = -x \frac{\eta}{2} \cos 2z - x \frac{3\eta}{2} + x \frac{\kappa}{2} \sin 2z + x \sin z$  (binary terms)



ELL1, non-relativistic systems, low  $e$

BT, non-relativistic systems, larger  $e$

# Time residuals

$\{N_0^{(1)}, \nu^{(1)}, \dots\}$  ... "first guess"

true values

estimated values

$$-\nu^{(1)} R(t) \equiv N(t, N_0, \nu, \dots) - N(t, N_0^{(1)}, \nu^{(1)}, \dots),$$

definition of time residuals

variation of orbital parameters

$$R(t) = \delta K - \frac{\partial N}{\partial x} \Big|_1 \frac{\delta x}{\nu} - \frac{\partial N}{\partial \eta} \Big|_1 \frac{\delta \eta}{\nu} - \frac{\partial N}{\partial \kappa} \Big|_1 \frac{\delta \kappa}{\nu} - \frac{\partial N}{\partial \Psi'} \Big|_1 \frac{\delta \Psi'}{\nu}$$

$$\delta K \equiv -\frac{\delta N_0}{\nu} - \frac{\delta \nu}{\nu} t - \frac{\delta \dot{\nu}}{2\nu} t^2 - \frac{\delta \ddot{\nu}}{6\nu} t^3$$

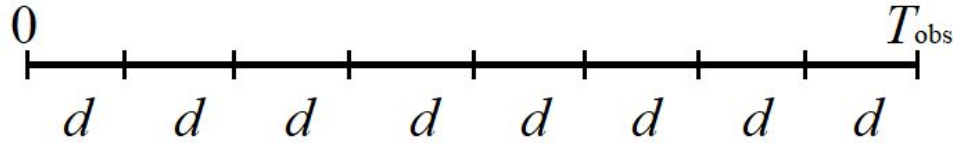
Post-Keplerian formalism:

$$\frac{\dot{x}}{x} = -2\alpha \dot{\Phi}, \quad \delta x(t) = \overset{\text{Error}}{\delta x_{\text{asc}}} + \overset{\text{Secular effect}}{\dot{x}(t - T_{\text{asc}})} - \overset{\text{ULDM signal}}{2\alpha x \Phi_0 \varrho [\cos(mt + \Upsilon) - \cos(mT_{\text{asc}} + \Upsilon)]}$$

# Two-step method

Step 1: variances

$$\delta\mathbf{S}_{\text{ELL1}} \equiv \{\delta K, \delta x, \delta\eta, \delta\kappa, \delta\Psi'\}$$



$n_c$  ... cadence

$n_c d$  ... number of observations

$$\frac{1}{\epsilon^2} \sum_{i=1}^{n_c d} \mathbf{M}^i \delta\mathbf{S} = \mathbf{D}$$

$$\sum_{i=1}^{n_c d} \mathbf{M}^i \simeq \frac{n_c d}{2\pi} \int_0^{2\pi} \mathbf{M}^i d\Theta' = n_c d \bar{\mathbf{M}}$$

$$\mathbf{C} = \frac{\epsilon^2}{n_c d} \bar{\mathbf{M}}^{-1}$$

# Two-step method

Step 2: time-dependence

$$\mathcal{B} \equiv \frac{\mathcal{O}_h}{\mathcal{O}_n}$$

evidence for signal hypothesis

evidence for null hypothesis

$$\mathcal{O}_n = \int \mathcal{L}(\Xi) d\vec{\Xi}$$

nuisance parameters (errors, secular effects)

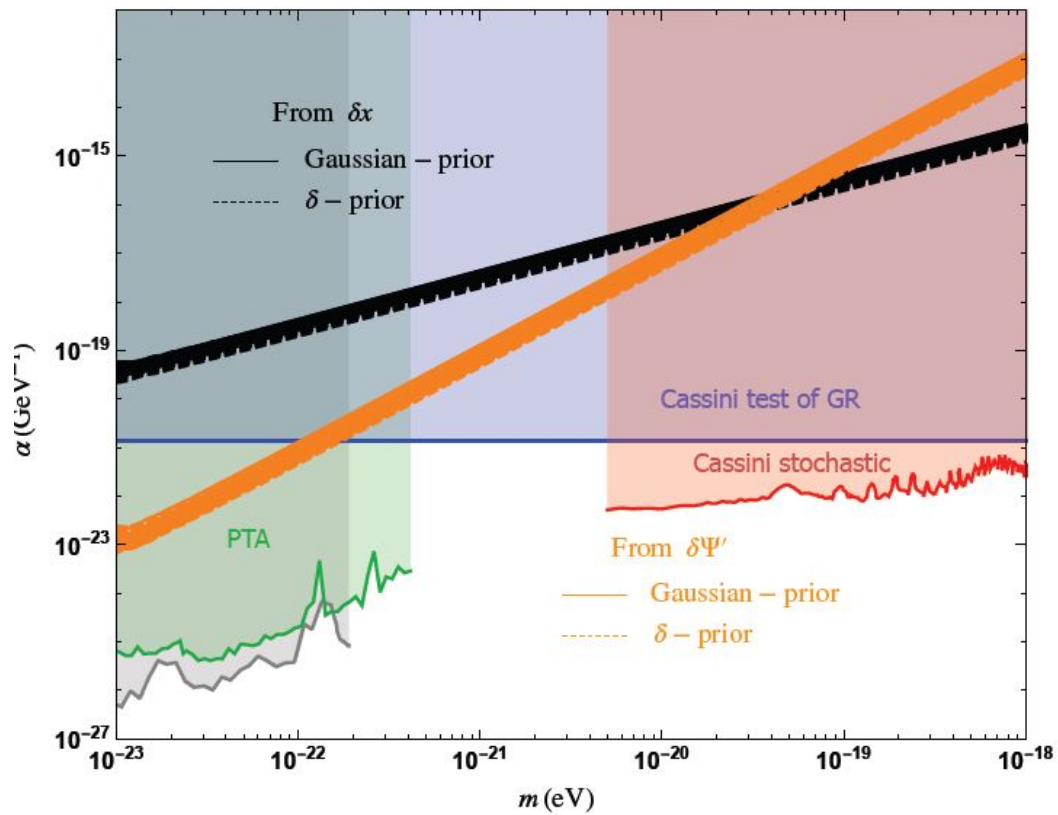
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22 ELL1, NANOGrav