



Istituto Nazionale di Fisica Nucleare
SEZIONE DI BOLOGNA



LRSM: Search for Doubly Charged Higgs bosons with the ATLAS detector + LNV Higgs decays

Bled 2024: International Workshop on Lepton Number Violation

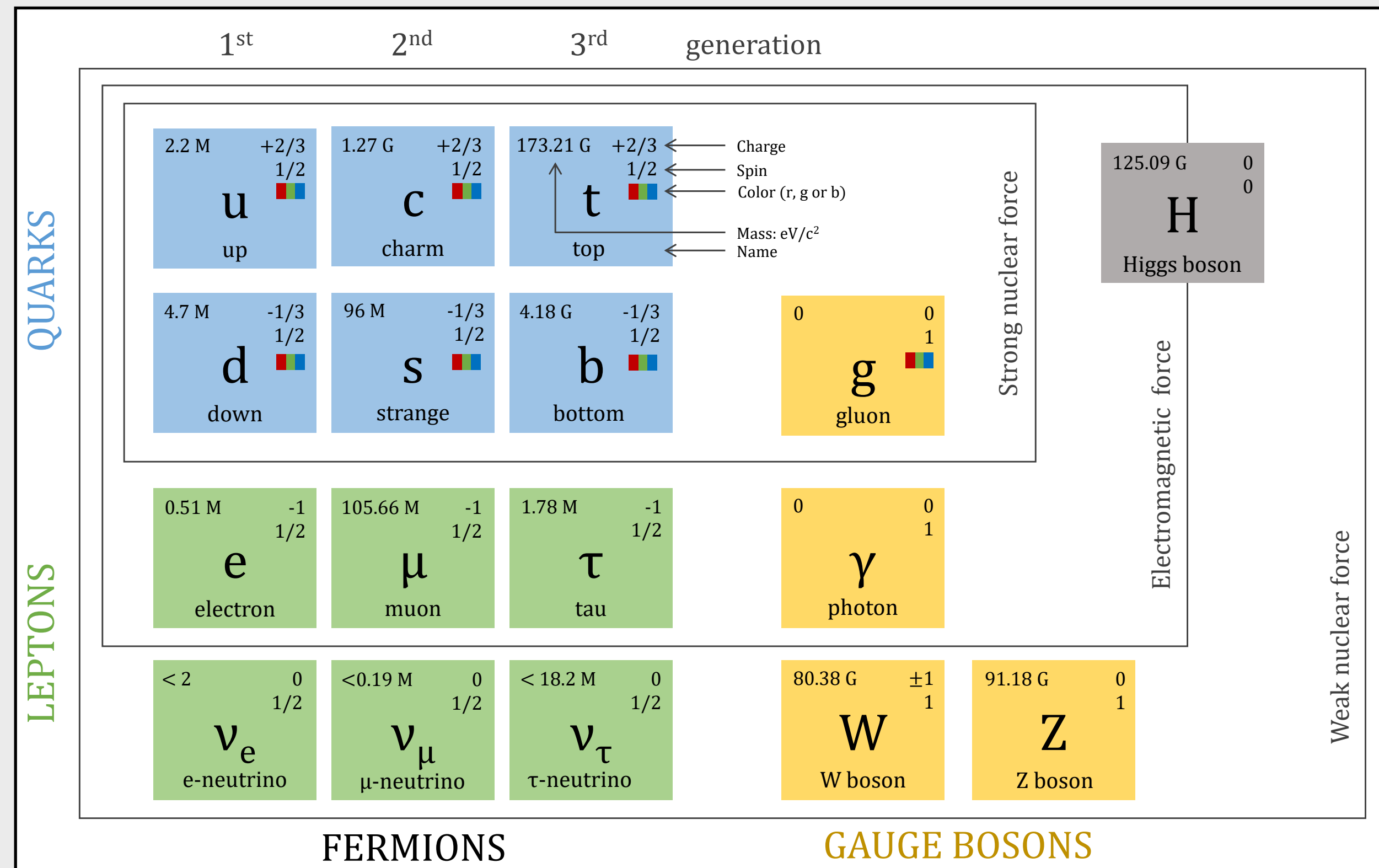
Blaž Leban



BLED 2024

- Introduction to the Standard Model (SM)
- Beyond the Standard Models (BSMs):
 - Seesaw mechanism,
 - Left-Right Symmetric Model (LRSM),
 - Zee-Babu model.
- Search for heavy doubly charged Higgs (DCH) bosons:
 - Analysis strategy,
 - Background estimation,
 - Statistical Analysis and Results.
- Lepton Number Violating Higgs decays:
 - Test of the KNN UFO model → signal sample validation.

- Well-established quantum field theory of elementary particles and their interactions with high predictive power.
- Gives accurate predictions, which (mostly) very well agree with the experiments.



Shortcomings:

- Many free parameters,
- Dark matter and dark energy,
- Matter - antimatter asymmetry,
- Hierarchy problem and naturalness,
- **Neutrino masses,**
- **Weak interaction is completely asymmetric between left- and right-handed fields.**

- The SM cannot explain non-zero neutrino masses at the re-normalisable level. Neutrino masses can only be obtained by:

- Introducing a new degree of freedom, i.e. a right-handed neutrino to build the **Dirac mass** term:

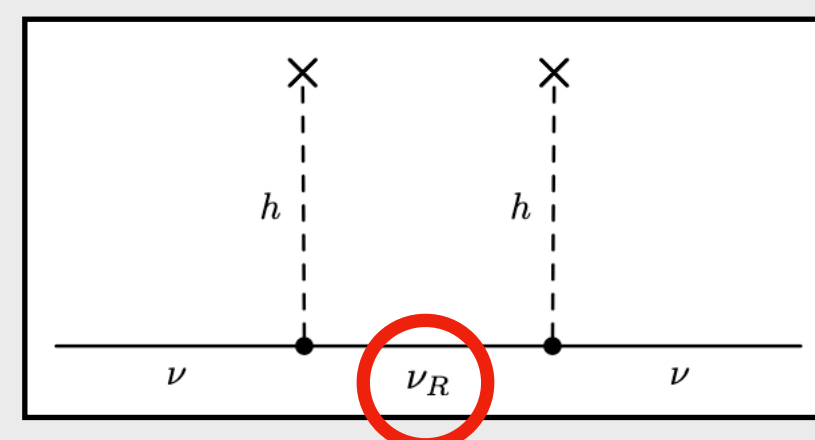
$$\mathcal{L}_D = -m_D (\bar{\nu}_R \nu_L + \bar{\nu}_L \nu_R), \quad \text{where} \quad m_D = y_D v / \sqrt{2}.$$

- Using non-renormalisable *dimension-5 Weinberg operator* to get **Majorana masses**:

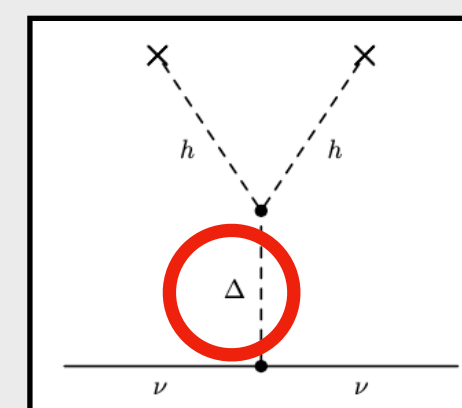
$$\mathcal{L}_M = -\frac{1}{2} m_M (\nu_L^T C \nu_L + h.c.), \quad \text{where} \quad m_M = y_M \frac{v^2}{\Lambda}.$$

- The observations of small neutrino masses indicate that either Λ is very large, $\Lambda \gg v$, or y must be small.
- The **seesaw mechanism** offers an explanation of small neutrino masses compared to those of quarks and charged leptons. Three types:

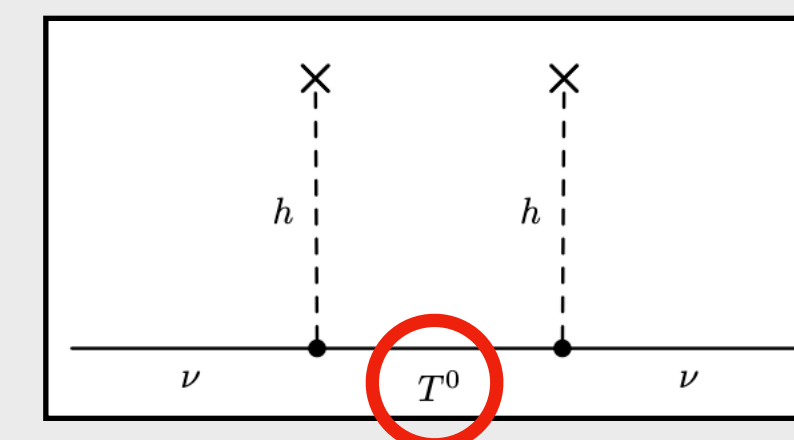
type-I: fermionic singlet



type-II: scalar triplet



type-III: fermionic triplet



- Parity violation remains unexplained within the SM, therefore we seek higher symmetry, which is then spontaneously broken at lower energy scales.
- Pati, Salam, Mohapatra, and Senjanović proposed the **left-right symmetric model**:

- Extend SM gauge group by an $SU(2)_R$ local symmetry:

$$SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}.$$

- Quarks and leptons become completely symmetric:

$$Q_{L,R} = \begin{pmatrix} u \\ d \end{pmatrix}_{L,R}, \quad L_{L,R} = \begin{pmatrix} \nu \\ e \end{pmatrix}_{L,R}.$$

- Interaction Lagrangian is written as $\mathcal{L} = \bar{f}i\gamma^\mu D_\mu f$, with the covariant derivative containing coupling strengths:

$$D_\mu = \partial_\mu - i \left(g_S A_\mu^a T_{ij}^a + g_L W_{\mu L}^a \frac{\sigma_a^L}{2} + g_R W_{\mu R}^a \frac{\sigma_a^R}{2} + g' B_\mu \frac{B-L}{2} \right).$$

- Higgs sector is extended with two SU(2) **triplets** $\Delta_{L,R}$ and a **bi-doublet** Φ :

$$\Delta_{L,R} = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}_{L,R}, \quad \Phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}.$$

- Spontaneous Symmetry Breaking (SSB) happens in two consecutive steps:

LRSM \rightarrow SM

- $\Delta_L(1, 3, 1, 2)$ and $\Delta_R(1, 1, 3, 2)$ triplets,
- $\langle \Delta_L \rangle = 0$ and $\langle \Delta_R \rangle = \begin{pmatrix} 0 & 0 \\ v_R & 0 \end{pmatrix}$.

SM \rightarrow EM

- $\Phi(1, 2, 2, 0)$ (includes SM Higgs boson),
- $\langle \Phi \rangle = \begin{pmatrix} v_1 & 0 \\ 0 & v_2 e^{i\alpha} \end{pmatrix}$,
- In turn, Δ_L develops a tiny induced vev $v_L = \langle \Delta_L \rangle \propto v^2/v_R$.

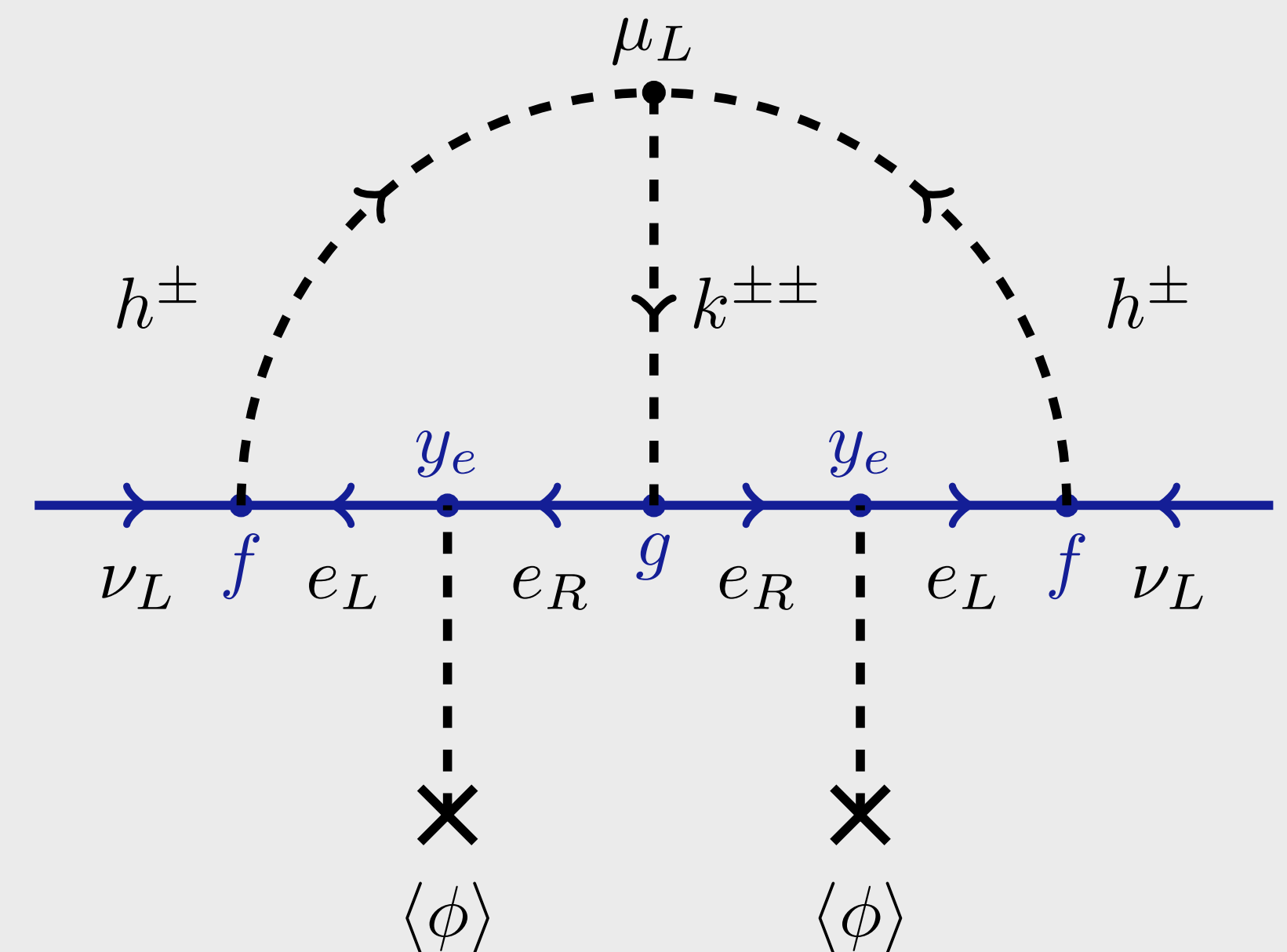
- Right-handed counterparts of W and Z bosons acquire masses as a result of the SSB.

- The existence of two scalar $SU(2)_L$ singlets k and h , which carry hypercharge, is hypothesised:

$$h^\pm \sim (1, 1, +1),$$

$$k^{\pm\pm} \sim (1, 1, +2).$$

- Since there are no right-handed neutrinos in the Zee–Babu model, Dirac mass terms can not be constructed. Furthermore, k and h cannot contract with lepton and Higgs doublets to generate Majorana masses at the tree level.
- Left-handed Majorana **neutrino masses** are generated **radiatively via two-loop diagrams**.
- Total and differential **production rates** for $H_R^{\pm\pm}$ and $k^{\pm\pm}$ **are identical** in shape and can only differ in normalisation.

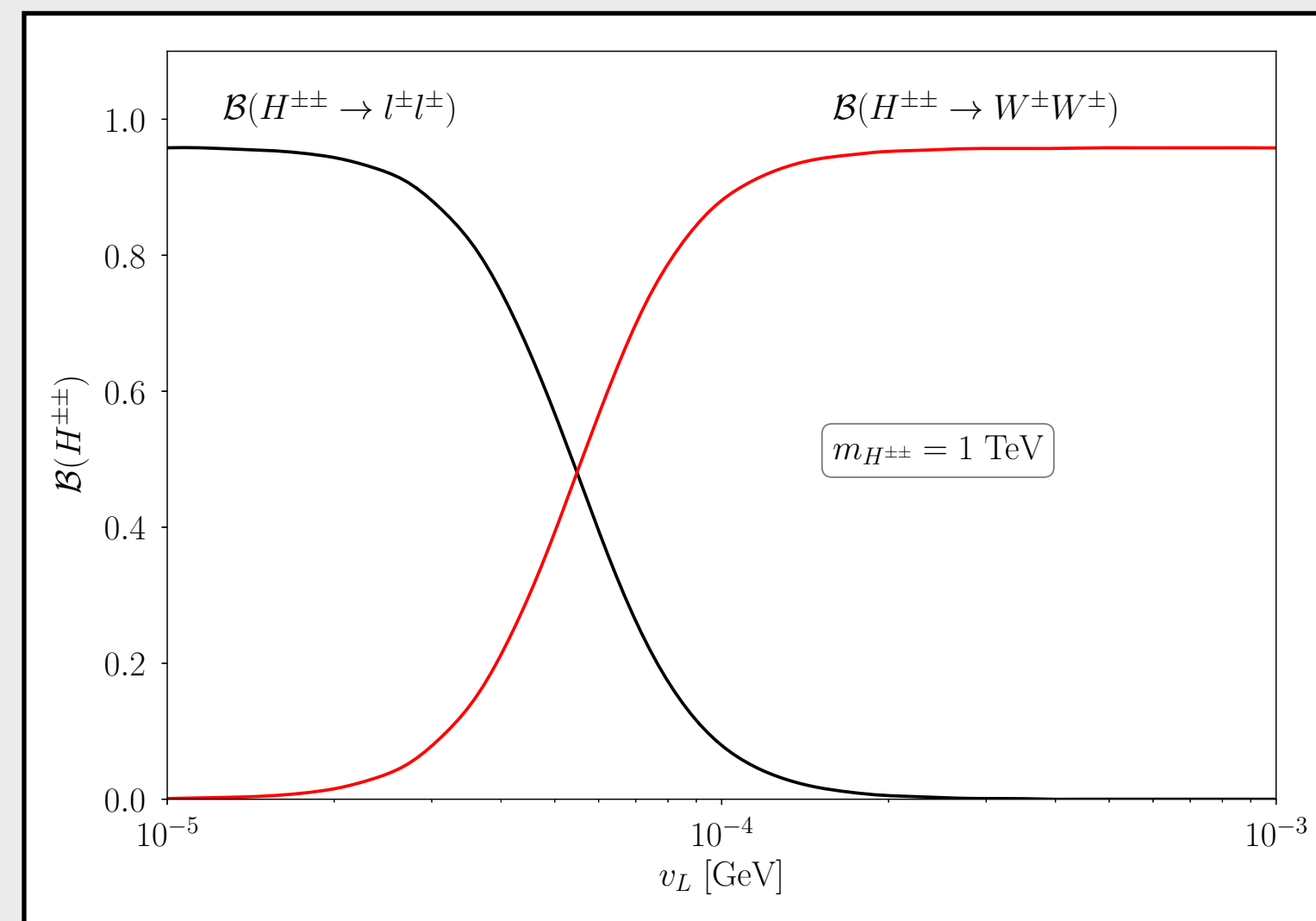
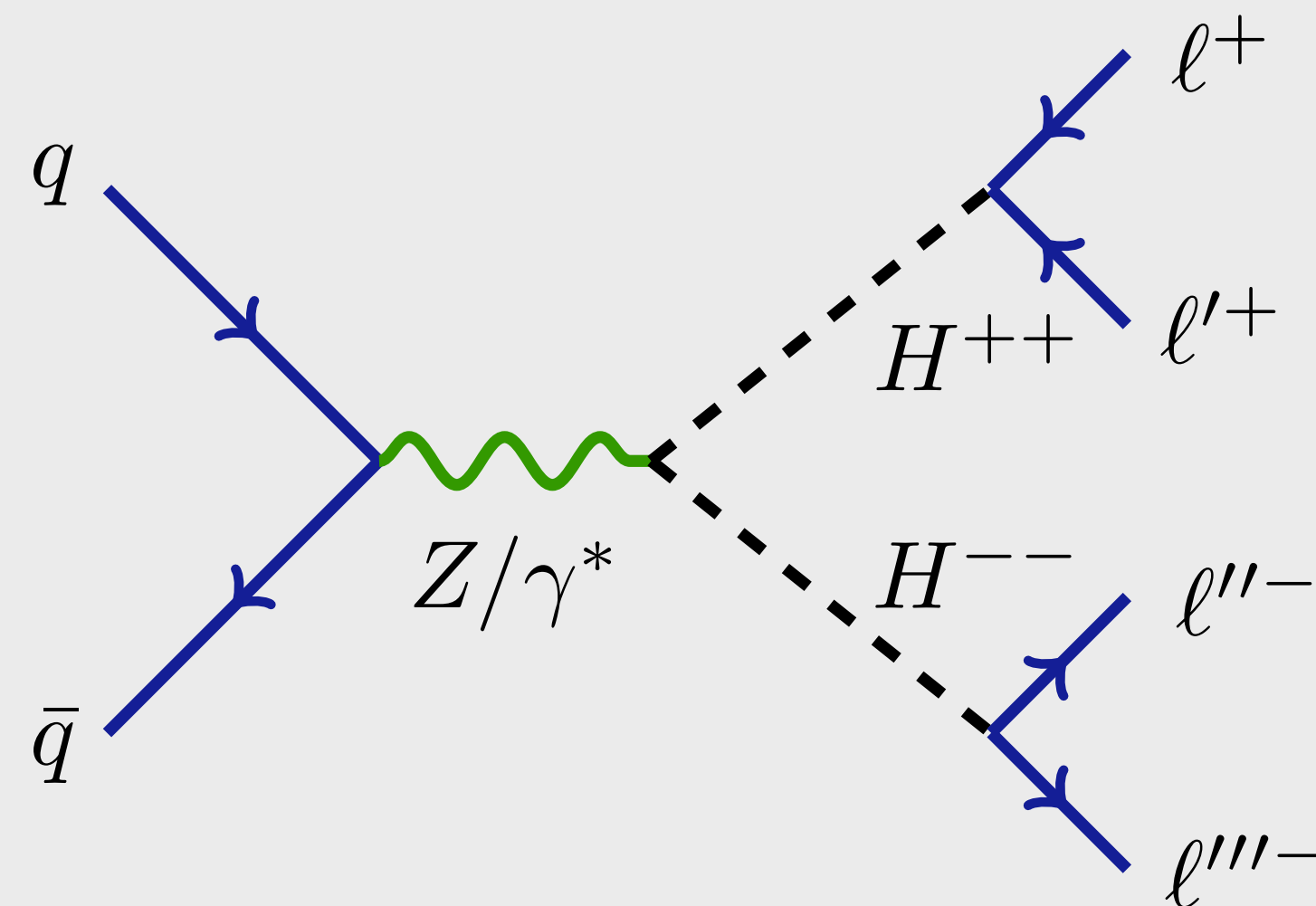


- Various BSM theories predict the existence of doubly charged bosons.

LRSM

- Two chiralities $H_L^{\pm\pm}$ and $H_R^{\pm\pm}$.
- Can couple to left-, right-handed leptons, vector bosons or scalars from the triplet. **Lepton flavour violation is allowed.**
- The dominant production mechanism is the Drell-Yan process through an s-channel photon or a Z boson exchange.

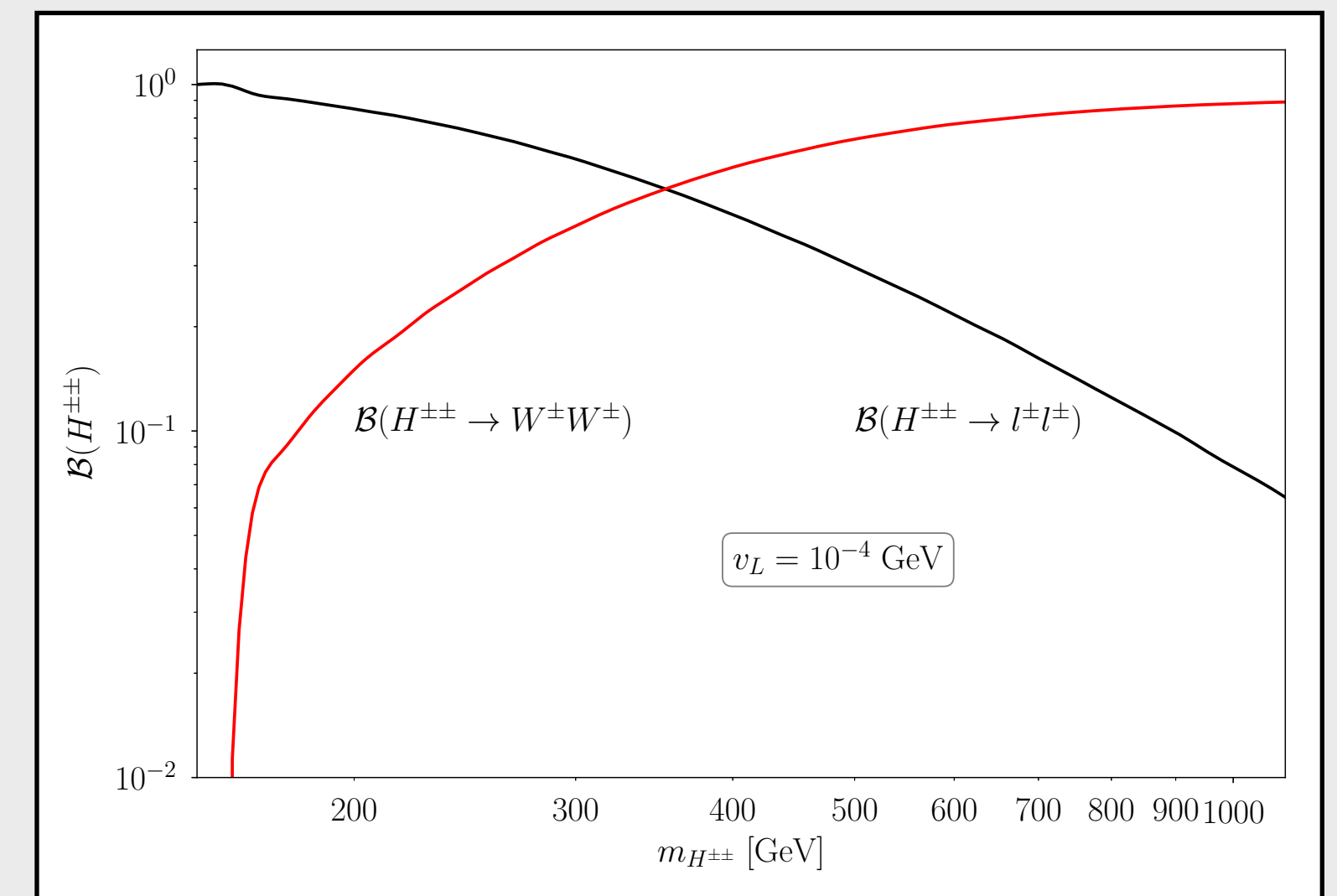
$$\Gamma(H^{\pm\pm} \rightarrow \ell^\pm \ell'^{\pm}) = \frac{1}{1 + \delta_{\ell\ell'}} \frac{h_{\ell\ell'}^2}{16\pi} m_{H^{\pm\pm}}$$



Zee-Babu model

- The $k^{\pm\pm}$ has the same quantum numbers as $H_R^{\pm\pm}$ so their electroweak production is the same.
- For the Drell-Yan production mechanism, cross-sections and differential scalar distributions in the Zee-Babu and type-II seesaw models differ at most by a normalisation factor.

$$\Gamma(k^{\pm\pm} \rightarrow \ell^\pm \ell'^{\pm}) = \frac{1}{1 + \delta_{\ell\ell'}} \frac{|g_{\ell\ell'}|^2}{4\pi} m_{k^{\pm\pm}}$$





Doubly charged Higgs analysis

- **Prompt, highly energetic, same-charge** lepton pairs represent a striking signature for BSM physics - such events are produced rarely in pp collisions by the SM processes.
- Considering only final states containing **light leptons**, including leptonic τ decays.
- Branching fraction to each possible leptonic final state is assumed equal:

$$\mathcal{B}(H^{\pm\pm} \rightarrow e^{\pm}e^{\pm} / e^{\pm}\mu^{\pm} / \mu^{\pm}\mu^{\pm} / e^{\pm}\tau^{\pm} / \mu^{\pm}\tau^{\pm} / \tau^{\pm}\tau^{\pm}) = 1/6$$

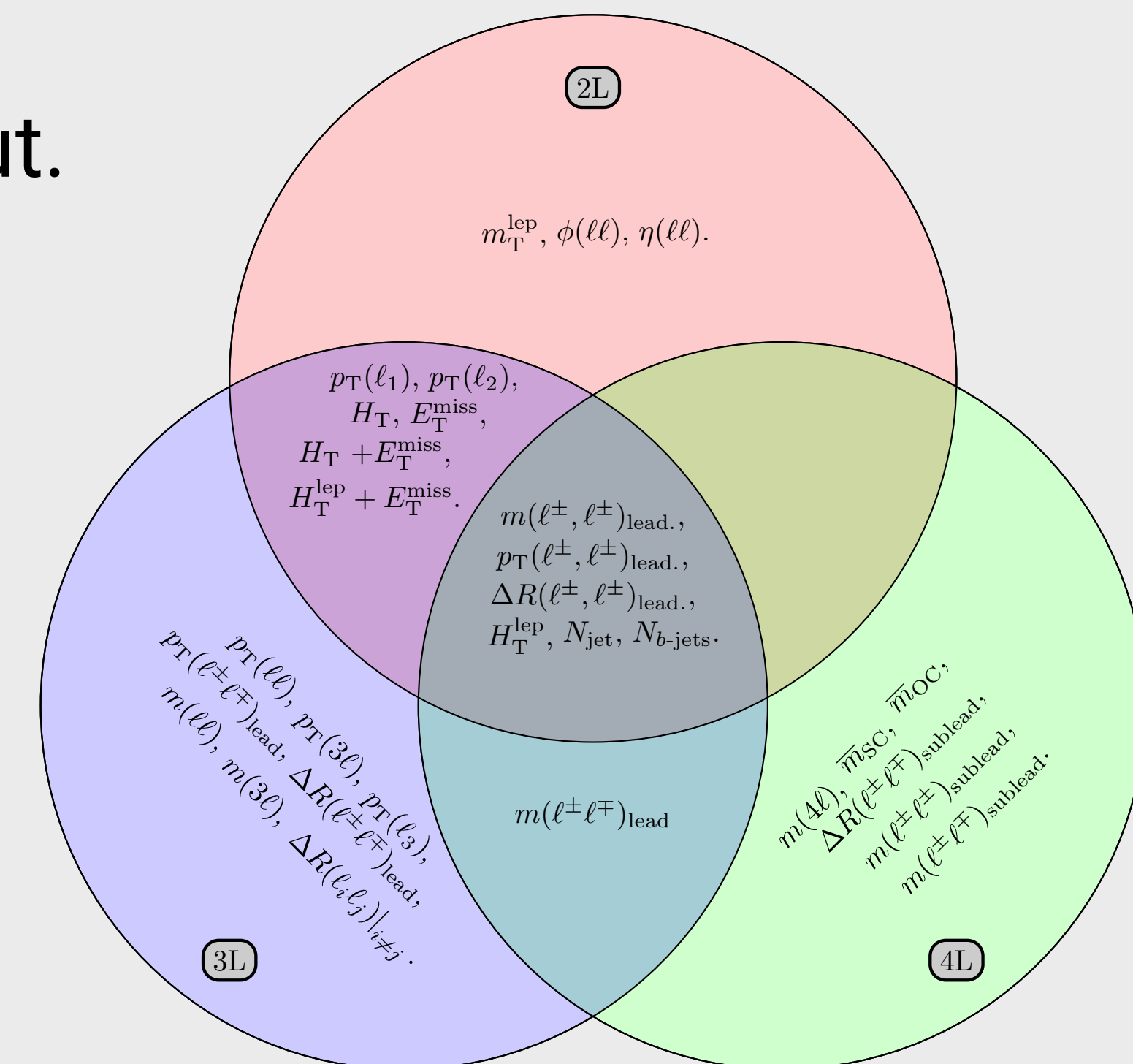
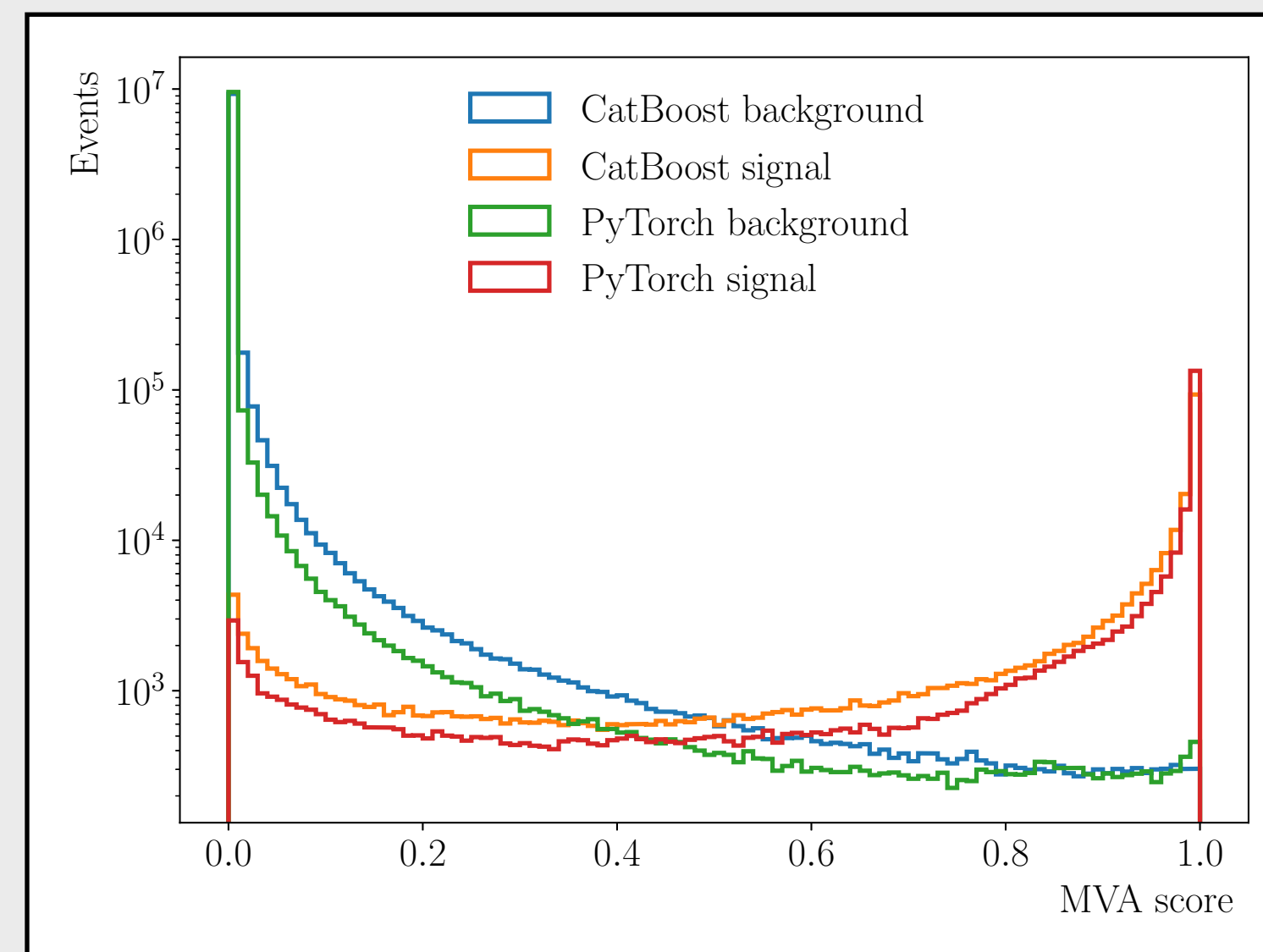
- Two analysis approaches were tested:
 - **Cut-based analysis (CB)**: selection requirements on kinematic variables are chosen orthogonally. Published in [Eur. Phys. J C 83 \(2023\) 605](#).
 - **Multivariate analysis (MVA)**: a functional dependence on a combination of observables is implemented.
- A binned maximum-likelihood fit to data is performed on:
 - $m(\ell^{\pm}\ell'^{\pm})_{\text{lead}}$ and event yield in the CB analysis,
 - logit MVA score in MVA analysis.

} Fit the dominant Drell-Yan and diboson backgrounds

- Two different binary classification algorithms are tested:
 - a gradient boosting on decision trees (BDT) implemented in CatBoost,
 - a deep neural network (DNN) provided by an ML library PyTorch.
- Due to the different topologies, an individual model is trained for each lepton multiplicity.
- Samples are split in 50% (training) : 25% (test) : 25% (validation) ratio, using stratified sampling.
- Features selected based on the BDT feature importance output.

- Neural Networks expectedly outperformed the BDTs.
- Signal to background was separated well, so it was challenging to construct sensible region definitions:

➔ use logit MVA score



- Six **control regions (CR)** for:
 - Constraining nuisance parameters related to systematic uncertainties,
 - Fitting the dominant SM backgrounds: DY (2ℓ) and diboson ($2\ell, 3\ell, 4\ell$).
- Five **validation regions (VR)** are used to cross-check the background modelling.
- Five **signal regions (SR)** as parts of phase space where a signal model predicts a significant excess of events over the expected background level.

	Control regions				Signal regions			Validation regions		
	DYCR	DBCR2L	DBCR3L	CR4L	SR2L	SR3L	SR4L	VR2L	VR3L	VR4L
Channel	e^+e^-	$e^\pm e^\pm$ $e^\pm \mu^\pm$ $\mu^\pm \mu^\pm$	$l^\pm l^\pm l^\mp$	$l^+ l^+ l^- l^-$	$e^\pm e^\pm$ $e^\pm \mu^\pm$ $\mu^\pm \mu^\pm$	$l^\pm l^\pm l^\mp$	$l^+ l^+ l^- l^-$	$e^\pm e^\pm$ $e^\pm \mu^\pm$ $\mu^\pm \mu^\pm$	$l^\pm l^\pm l^\mp$	$l^+ l^+ l^- l^-$
$m(l^\pm, l^\pm)_{\text{lead.}}$ [GeV]*	≥ 300	[200, 300)	≥ 300	[100, 200)	≥ 300	≥ 300	≥ 300	≥ 300	[100, 300)	[200, 300)
$p_T(l^\pm, l^\pm)_{\text{lead.}}$ [GeV]	-	-	-	-	≥ 300	≥ 300	-	[200, 300)	-	-
$\Delta R(l^\pm, l^\pm)_{\text{lead.}}$	-	-	-	-	< 3.5	-	-	< 3.5	-	-
\bar{m}_{SC} [GeV]	-	-	-	-	-	-	≥ 300	-	-	-
E_T^{miss} [GeV]	-	> 30	-	-	-	-	-	> 30	-	-
$ \eta(l, l') $	-	< 3.0	-	-	-	-	-	< 3.0	-	-
Z-boson veto	-	-	inverted	-	-	✓	✓	-	✓	-
logit MVA result	-	-	-	-	≥ 0	≥ 3	≥ 0	< 0	< 3	< 0
$m(l^\pm, l^\pm)_{\text{lead.}}$ [GeV]*	≥ 300	[200, 300)	-	-	≥ 300	-	-	≥ 300	-	-
$p_T(l^\pm, l^\pm)_{\text{lead.}}$ [GeV]	-	-	-	-	≥ 300	-	-	-	-	-
\bar{m}_{SC} [GeV]	-	-	-	< 200	-	-	≥ 200	-	-	≥ 200
$ \eta(l, l') $	-	-	-	-	-	-	< 3.0	-	-	-
Z-boson veto	-	-	inverted	-	-	✓	-	-	✓	-

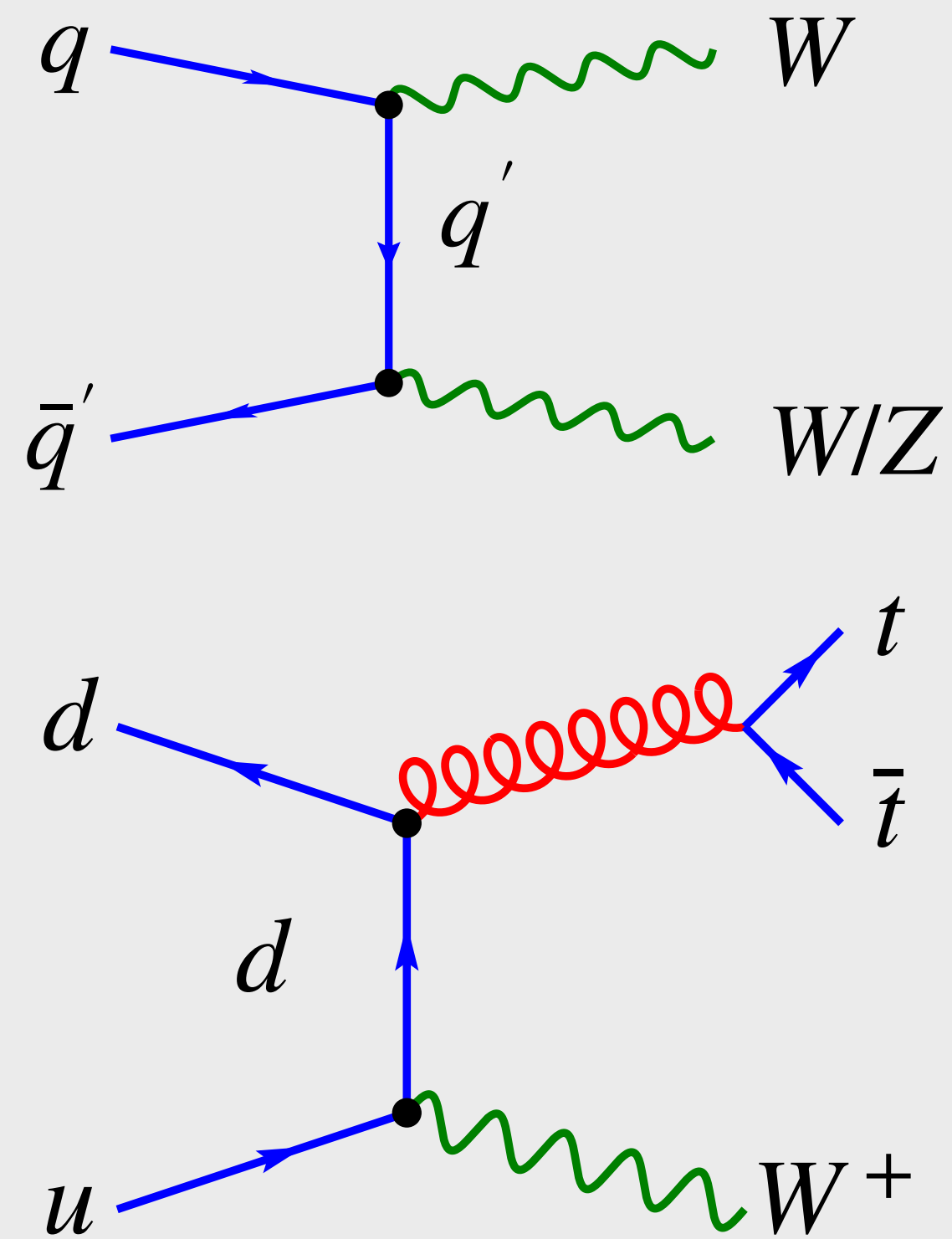
*Note that the cut on the invariant mass of the leading **opposite-charge** lepton pair is applied in DYCR region.

DiLepton Triggers

	2015	2016	2017+2018
ee	HLT_2e12_lhloose_L12EM10VH	HLT_2e17_lhvloose_nod0	HLT_2e17_lhvloose_nod0 OR HLT_2e24_lhvloose_nod0
$e\mu$	HLT_e17_lhloose_mu14		HLT_e17_lhloose_nod0_mu14
$\mu\mu$	HLT_mu18_mu8noL1		HLT_mu22_mu8noL1

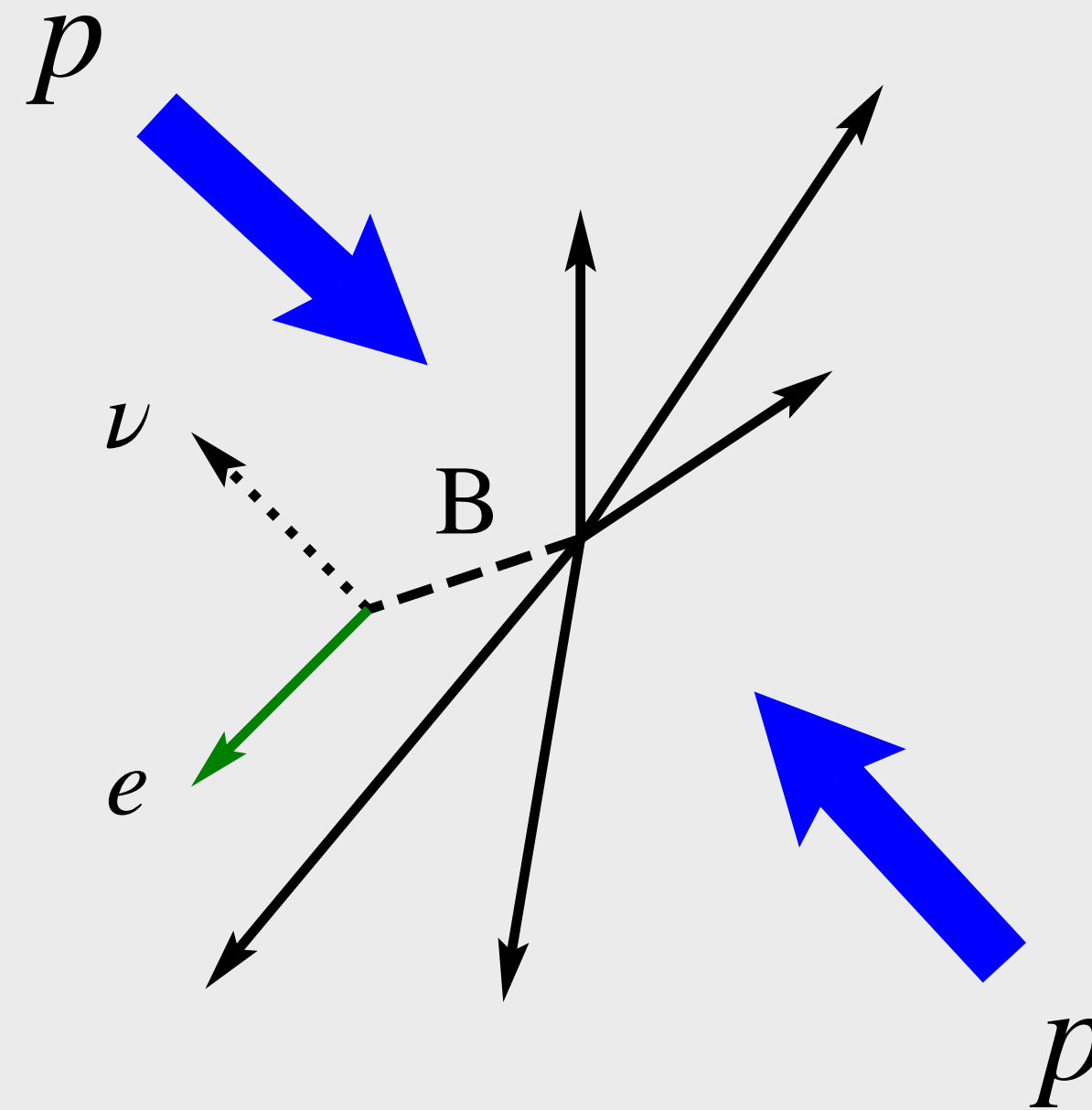
- To extract the signal, **dilepton triggers** are used in the analysis - main reason is to avoid trigger bias on fakes.
- Prescaled single lepton triggers are used for fake estimation (see next slides):
 - Average prescales which are derived from luminosity blocks are used.

- Prompt



Produced in hard interaction, e.g. diboson, $t\bar{t}W$ backgrounds
estimated from MC simulations.

- Fake/non-prompt

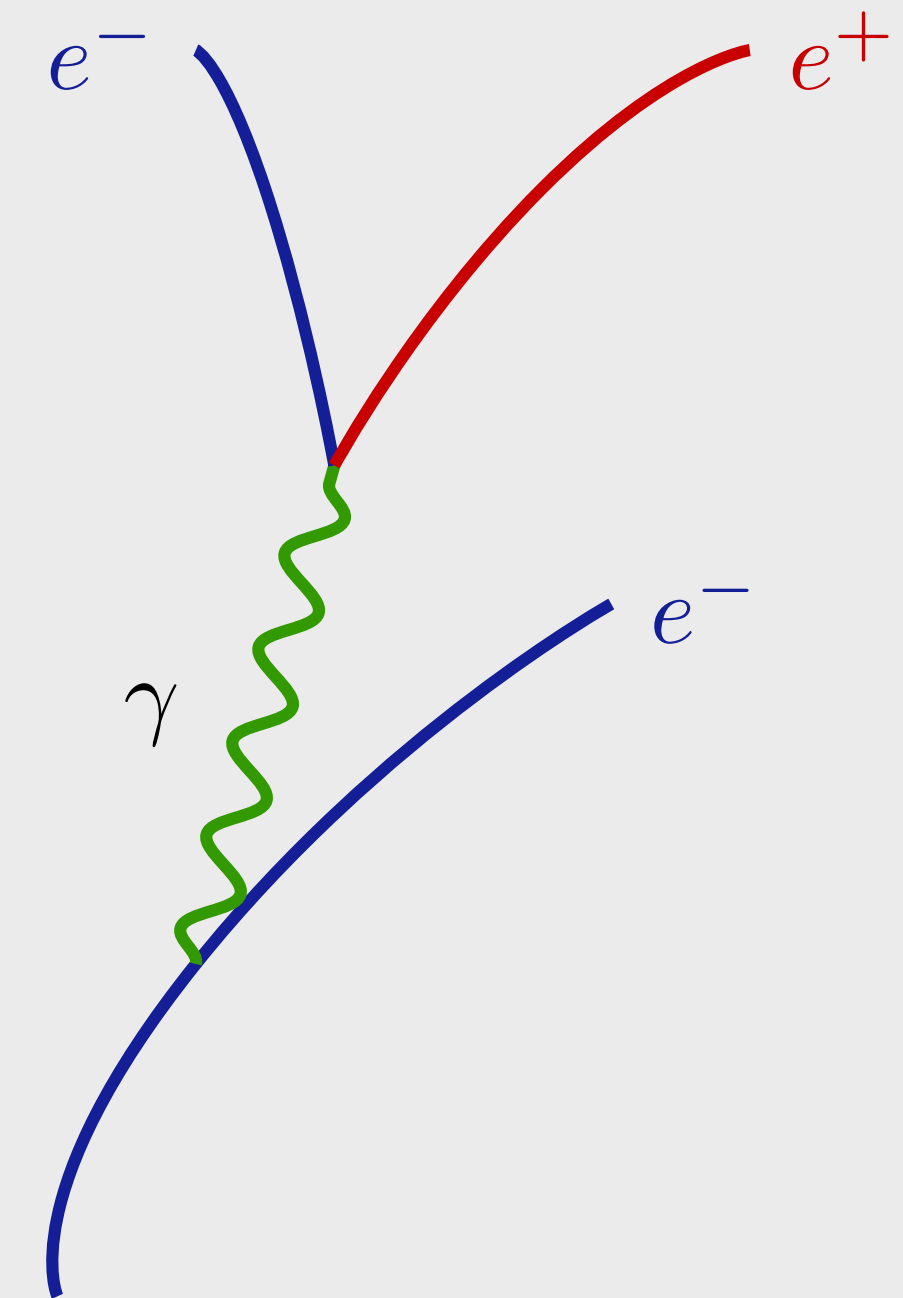


Any other object misidentified as lepton (jets, muons, hadrons, ...).

Leptons coming from heavy flavoured mesons or mis-identified jets.

Estimated with **data-driven fake factor method.**

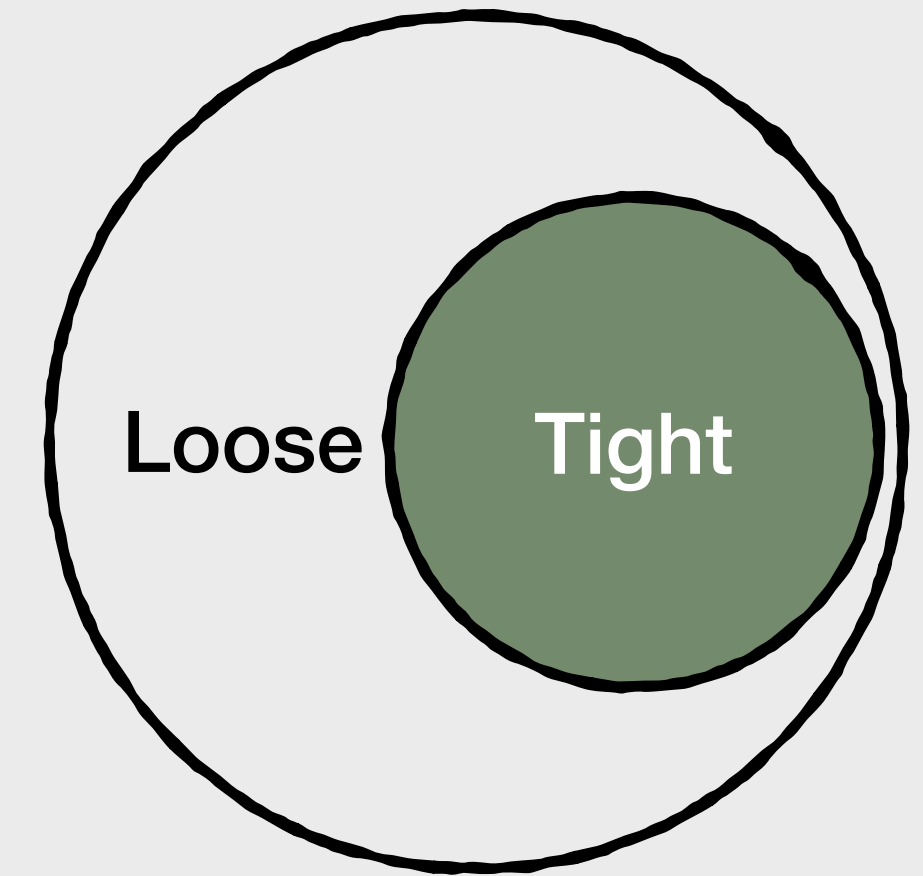
- Misidentified charge



Charge misidentification
estimated by comparing data to MC, simulated events corrected with scale factors.

- Fake-factor method is a simplified matrix method - we define loose (not tight) and tight regions that are orthogonal to each other.
- Two lepton definitions:
 - **tight**: nominal signal selection - leptons used in the analysis,
 - **loose**: relaxed identification, isolation cuts.
- Fake-factor is the ratio of tight and loose leptons: $F_\ell = \frac{N_{tight}}{N_{loose}}$.
- Fake estimation is done in **fake enriched regions**, see next slides.
- Fake factors are applied to each fake lepton individually:

Fake-factor Method

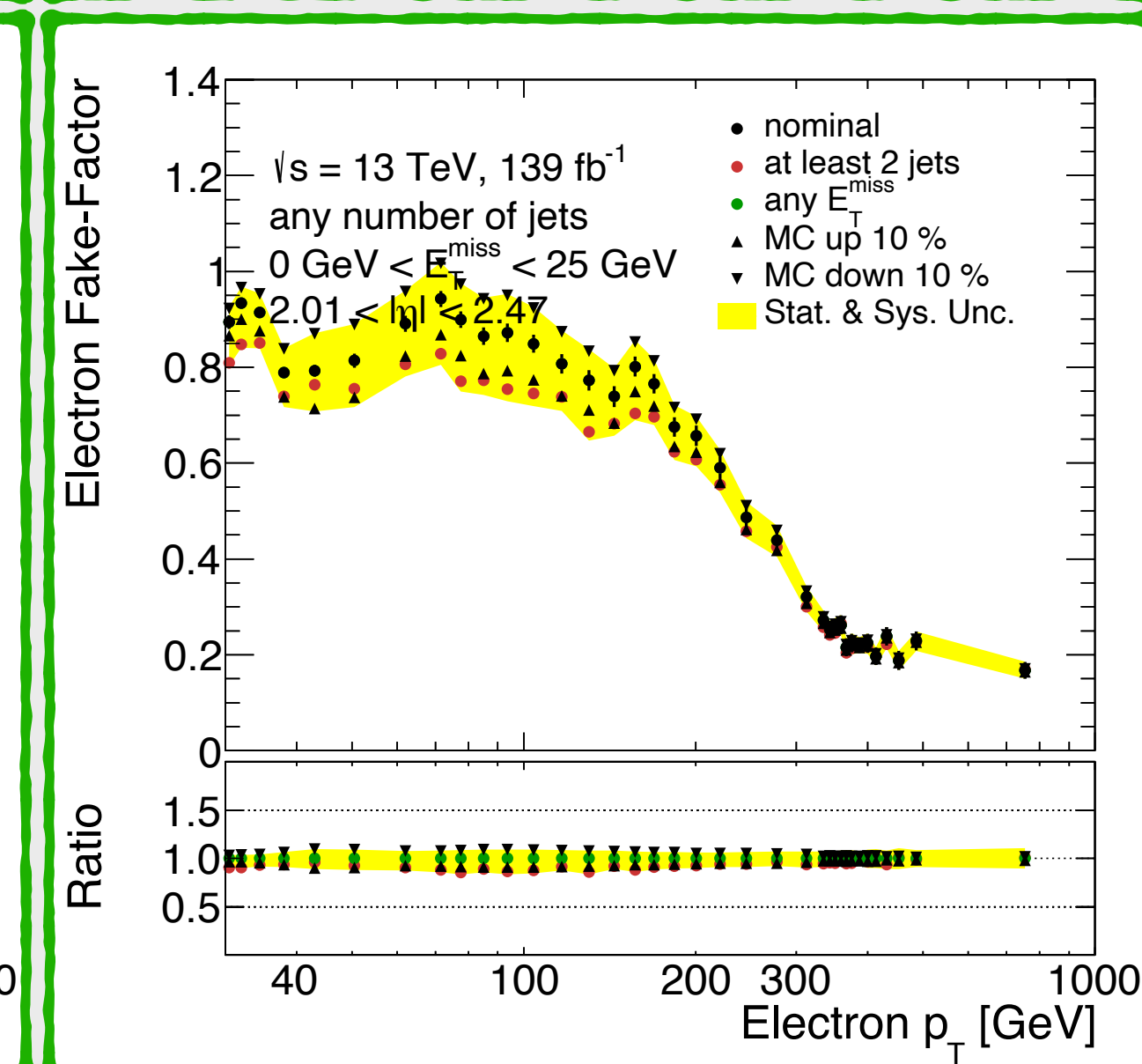
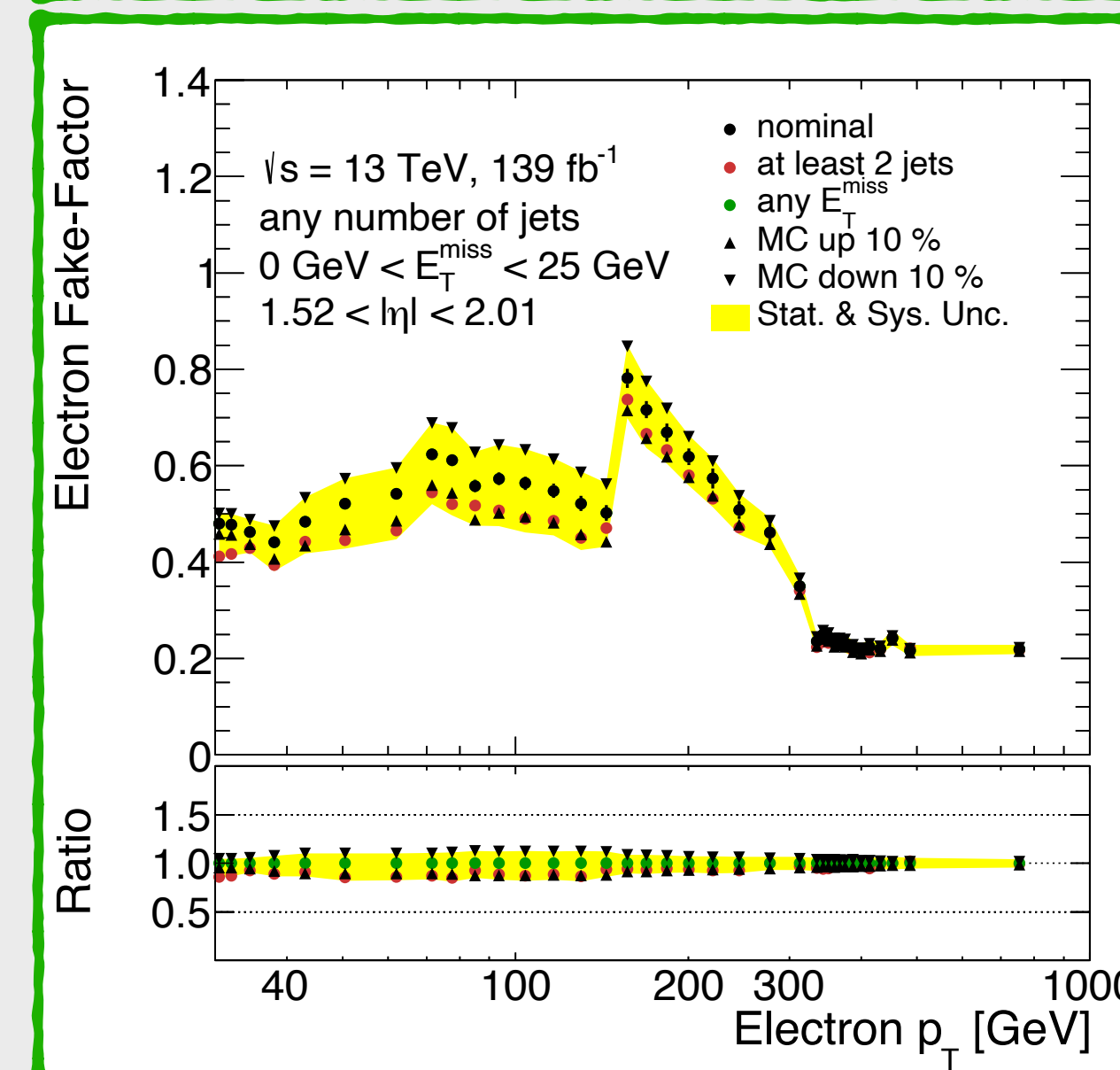
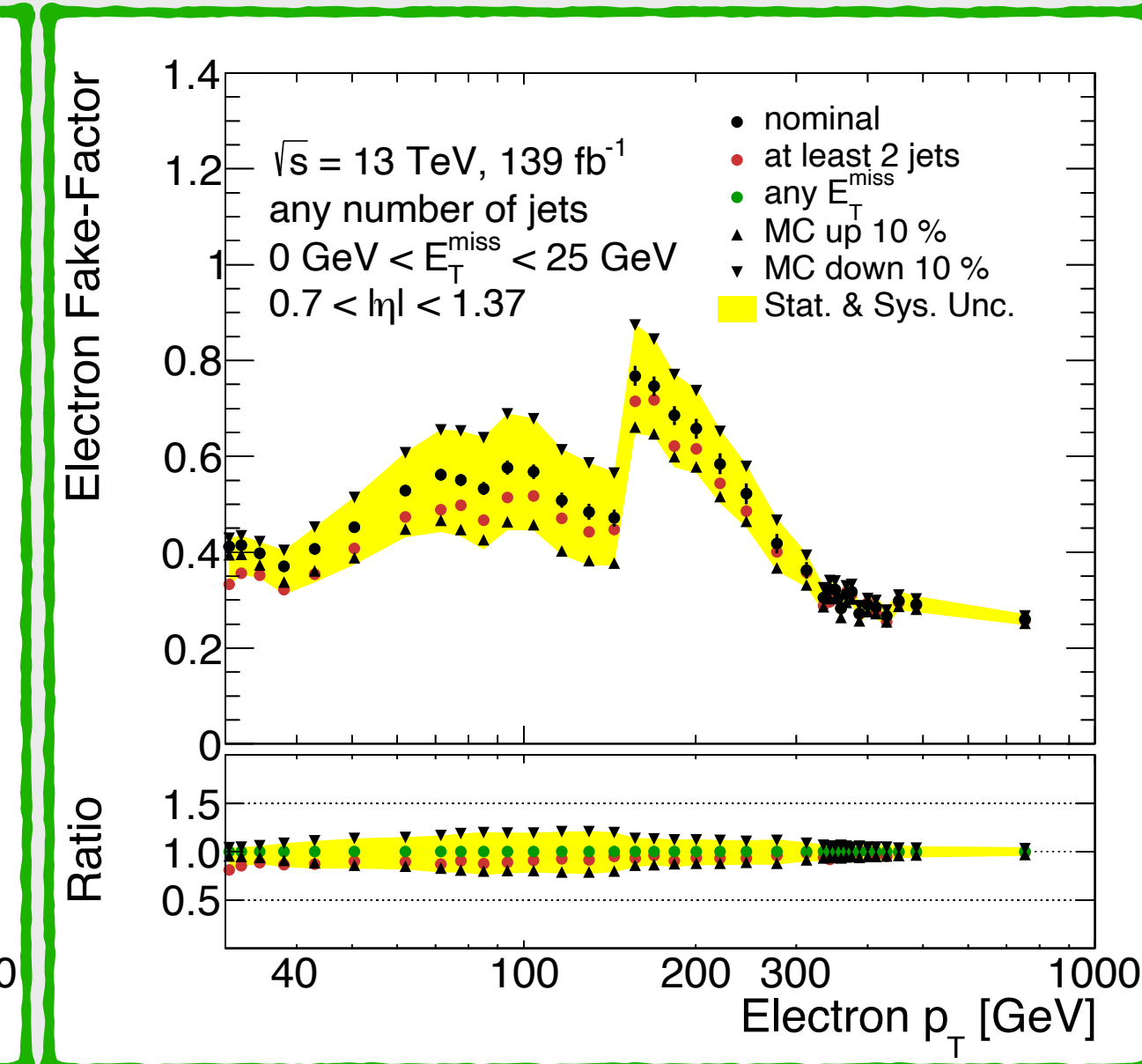
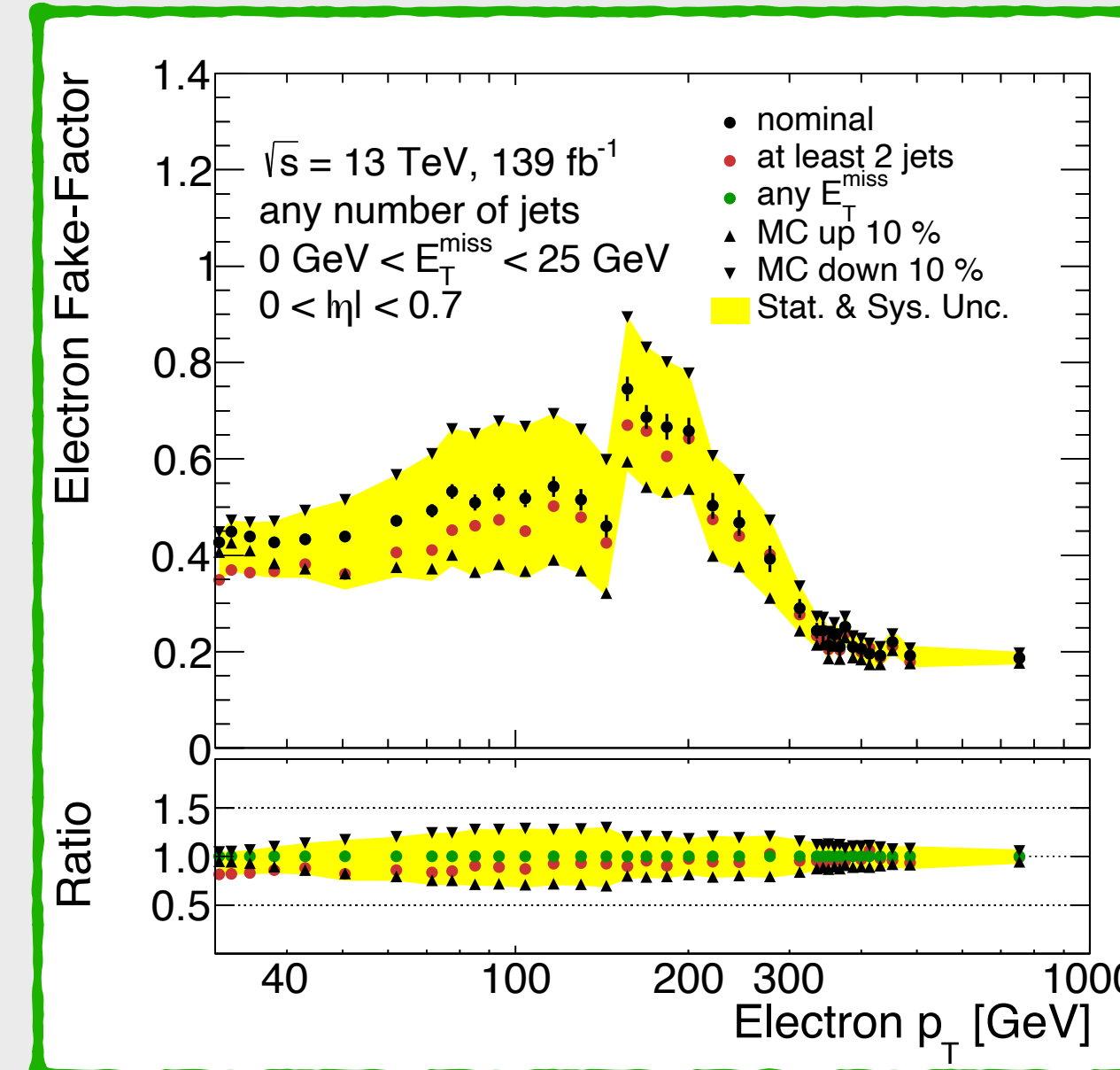


	L ₁	T ₁
L ₂	2 fake leptons: -F ₁ F ₂	1 fake lepton: +F ₂
T ₂	1 fake lepton: +F ₁	SR

$$N^{fake} = \sum_{i=1}^{N_{SB}^{data}} (-1)^{N_{L,i}+1} \prod_{l=1}^{N_{L,i}} F_l - \sum_{i=1}^{N_{SB}^{MC}} (-1)^{N_{L,i}+1} \prod_{l=1}^{N_{L,i}} F_l$$

$$\stackrel{l=2}{=} \left[\sum_{TL} F_2 + \sum_{LT} F_1 - \sum_{LL} F_1 F_2 \right]_{data} - \left[\sum_{TL} F_2 + \sum_{LT} F_1 - \sum_{LL} F_1 F_2 \right]_{prompt}$$

- Fake enriched region requires:
 - exactly 1 baseline electron,
 - any number of jets,
 - b - jet veto.
- Four η bins, two E_T^{miss} bins, variable p_T binning to ensure statistical error below a fixed threshold.



Trigger	Average prescale	Period
HLT_e26_lhvloose_nod0_L1EM20VH	111.2	2015-2016
HLT_e28_lhvloose_nod0_L1EM20VH	367.6	2017
HLT_e28_lhvloose_nod0_L1EM22VH	384.5	2018
HLT_e60_lhvloose_nod0	32.93	2015-2018
HLT_e70_lhvloose_nod0	64.13	2018
HLT_e80_lhvloose_nod0	40.43	2018
HLT_e100_lhvloose_nod0	19.45	2018
HLT_e120_lhvloose_nod0	12.15	2016, 2018
HLT_e140_lhvloose_nod0	2.637	2017-2018
HLT_e160_lhvloose_nod0	1.601	2017-2018
HLT_e200_etcut	-	2015
HLT_e300_etcut	-	2016-2018

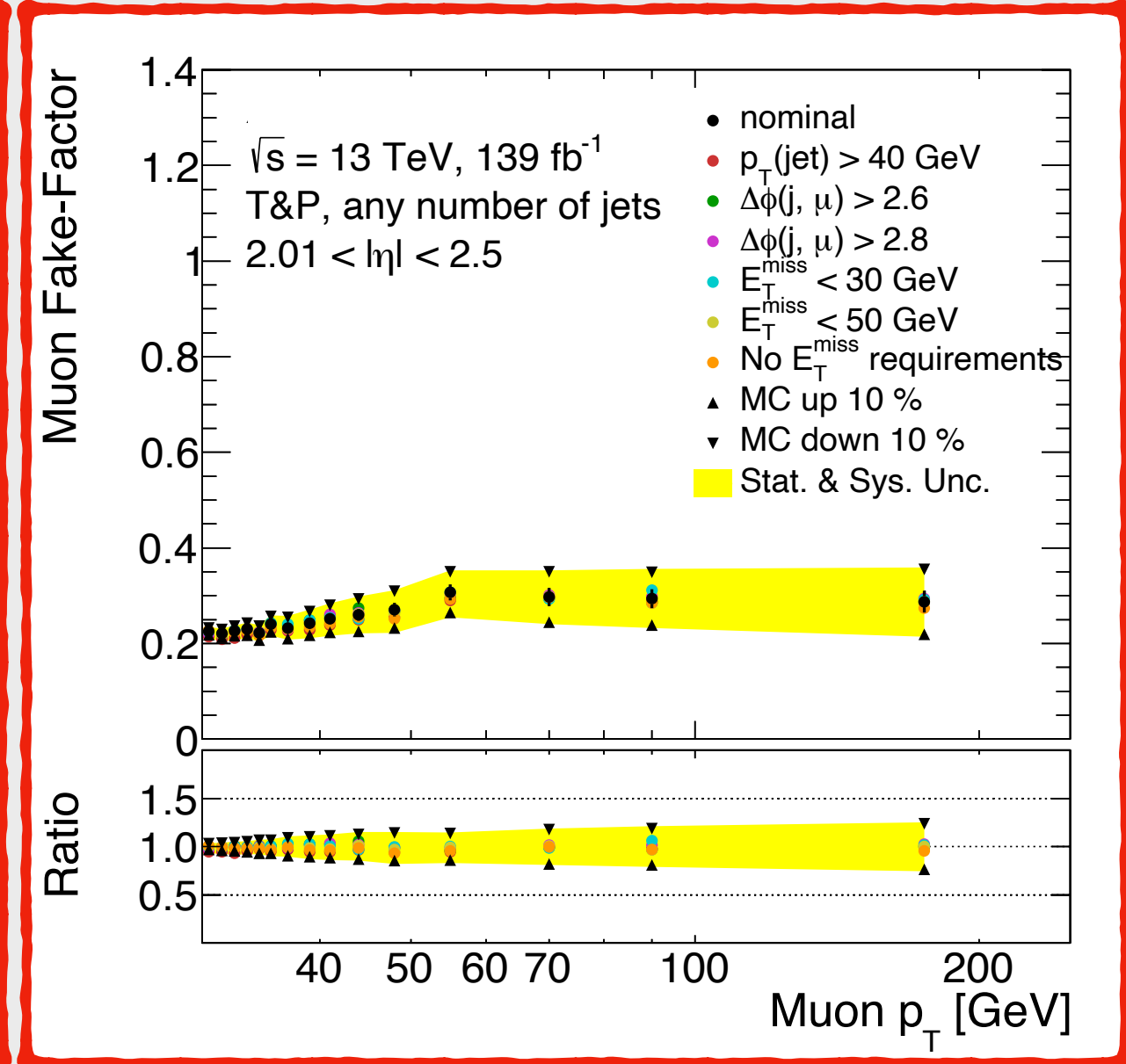
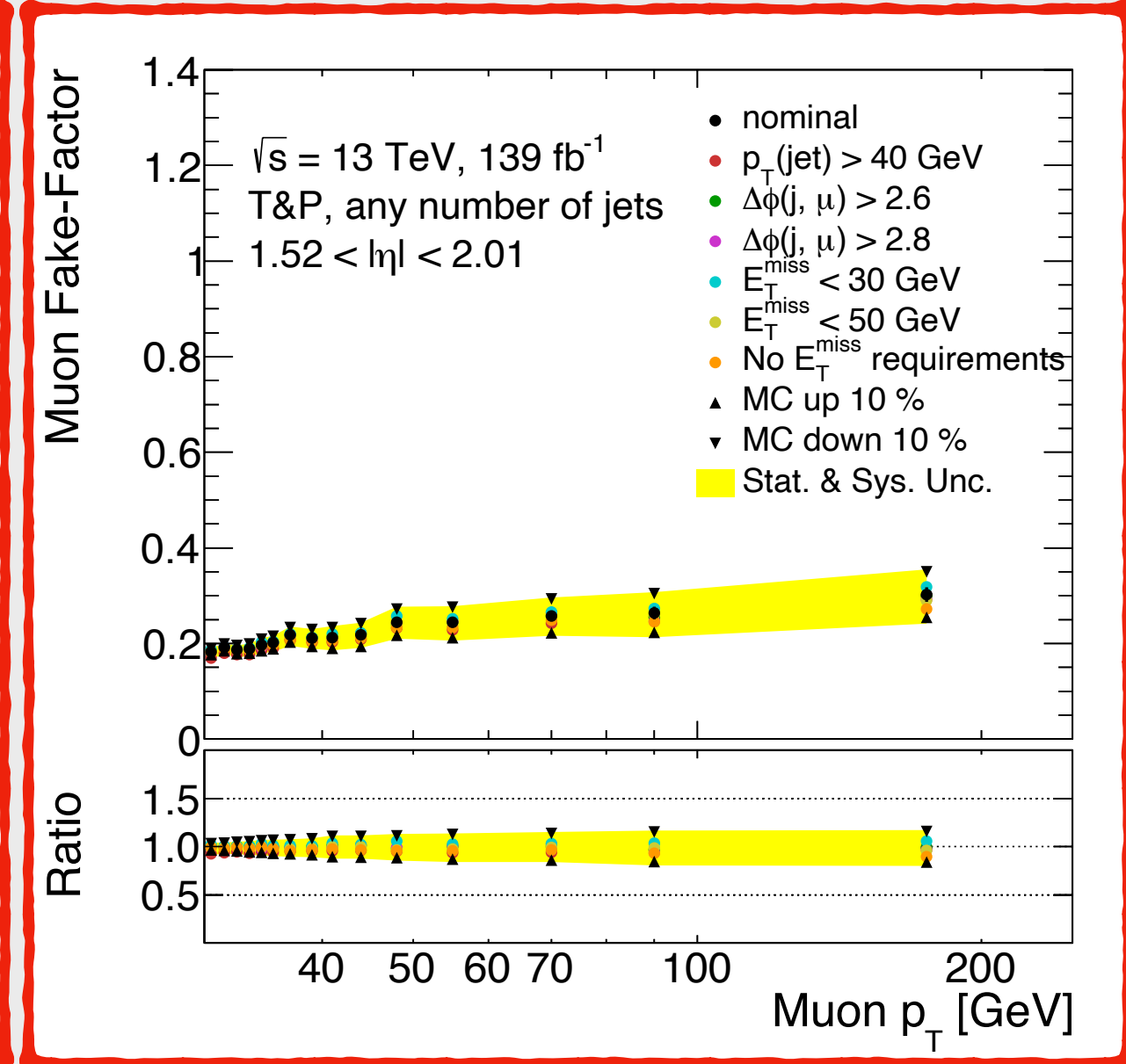
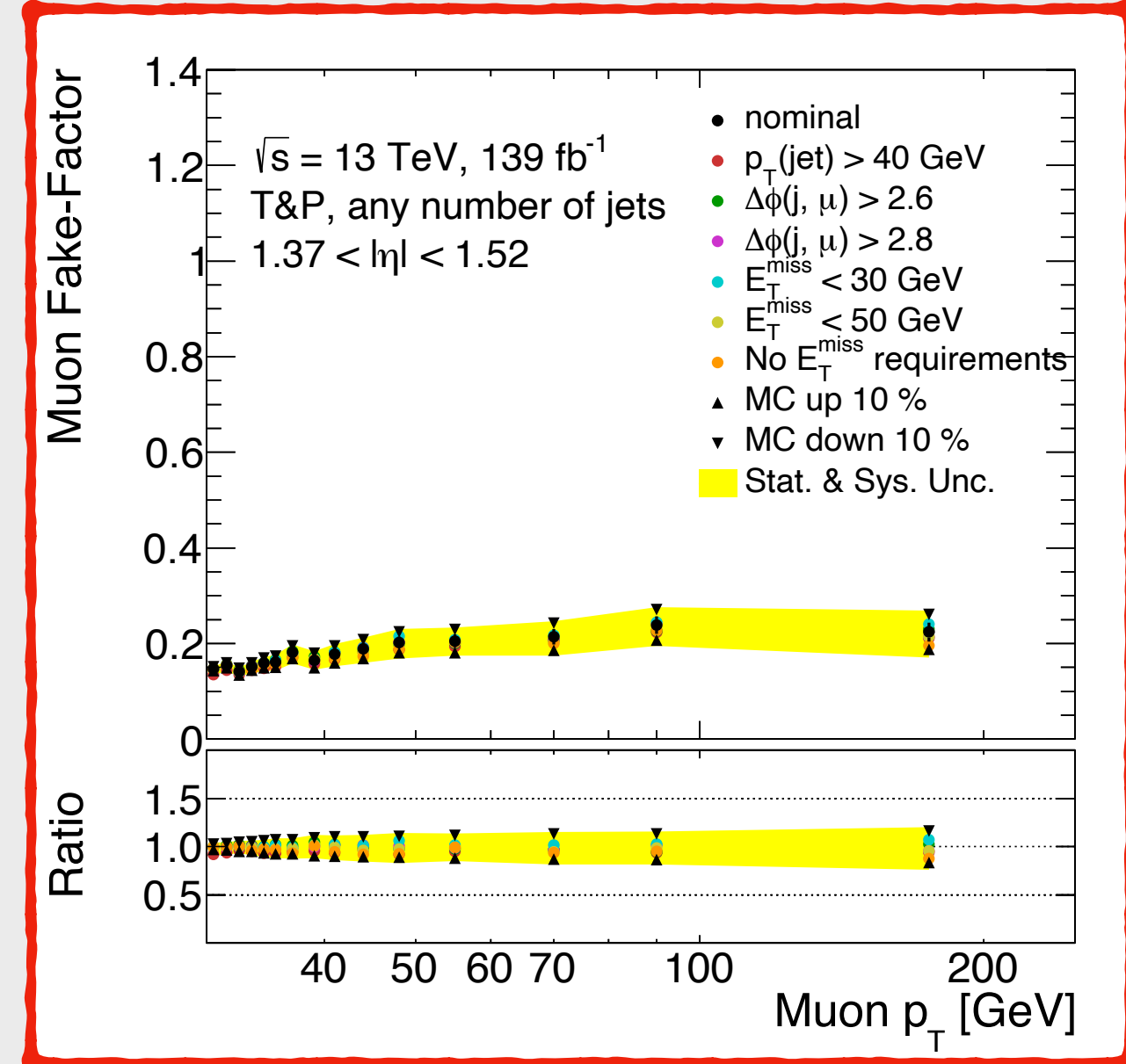
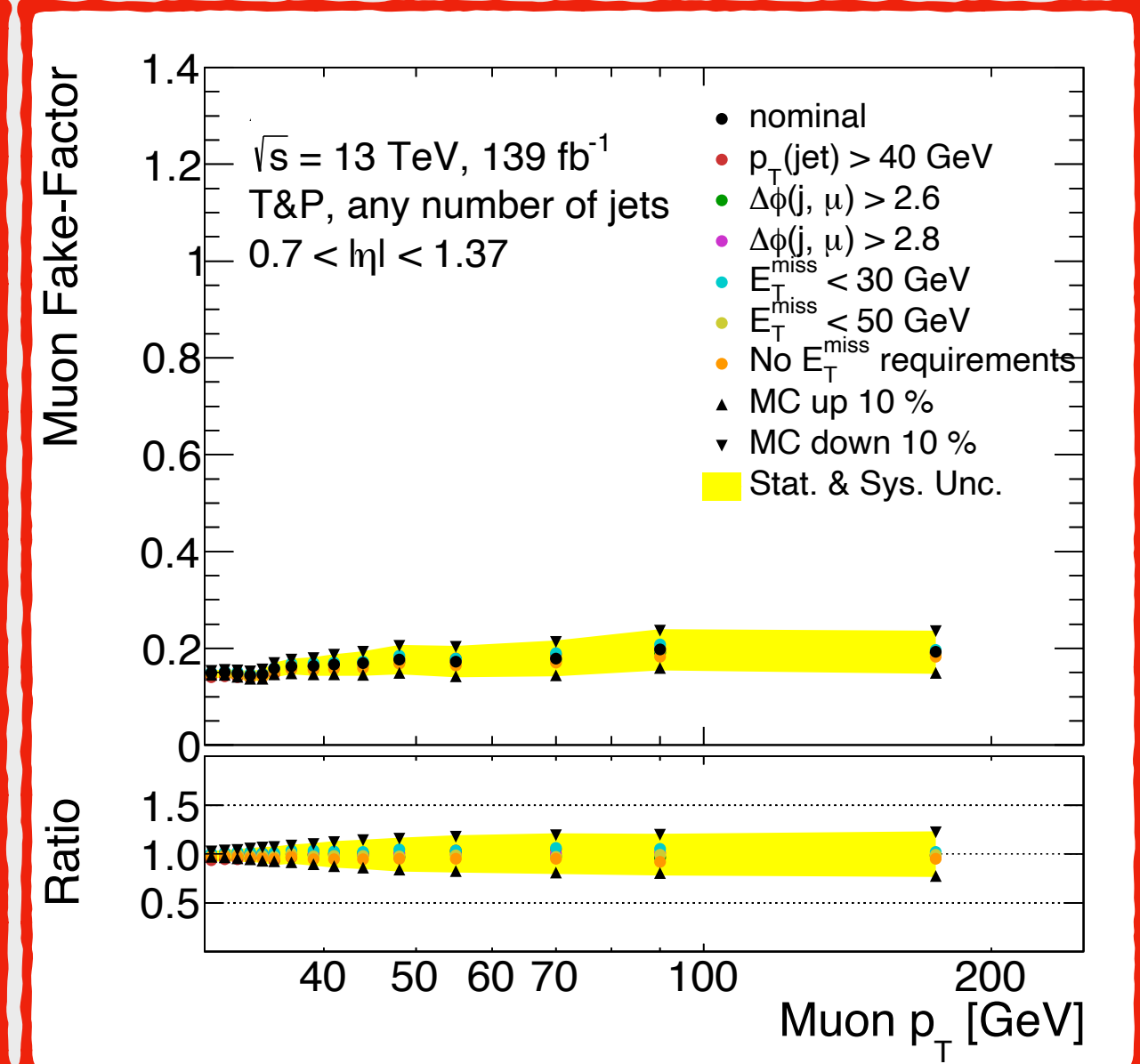
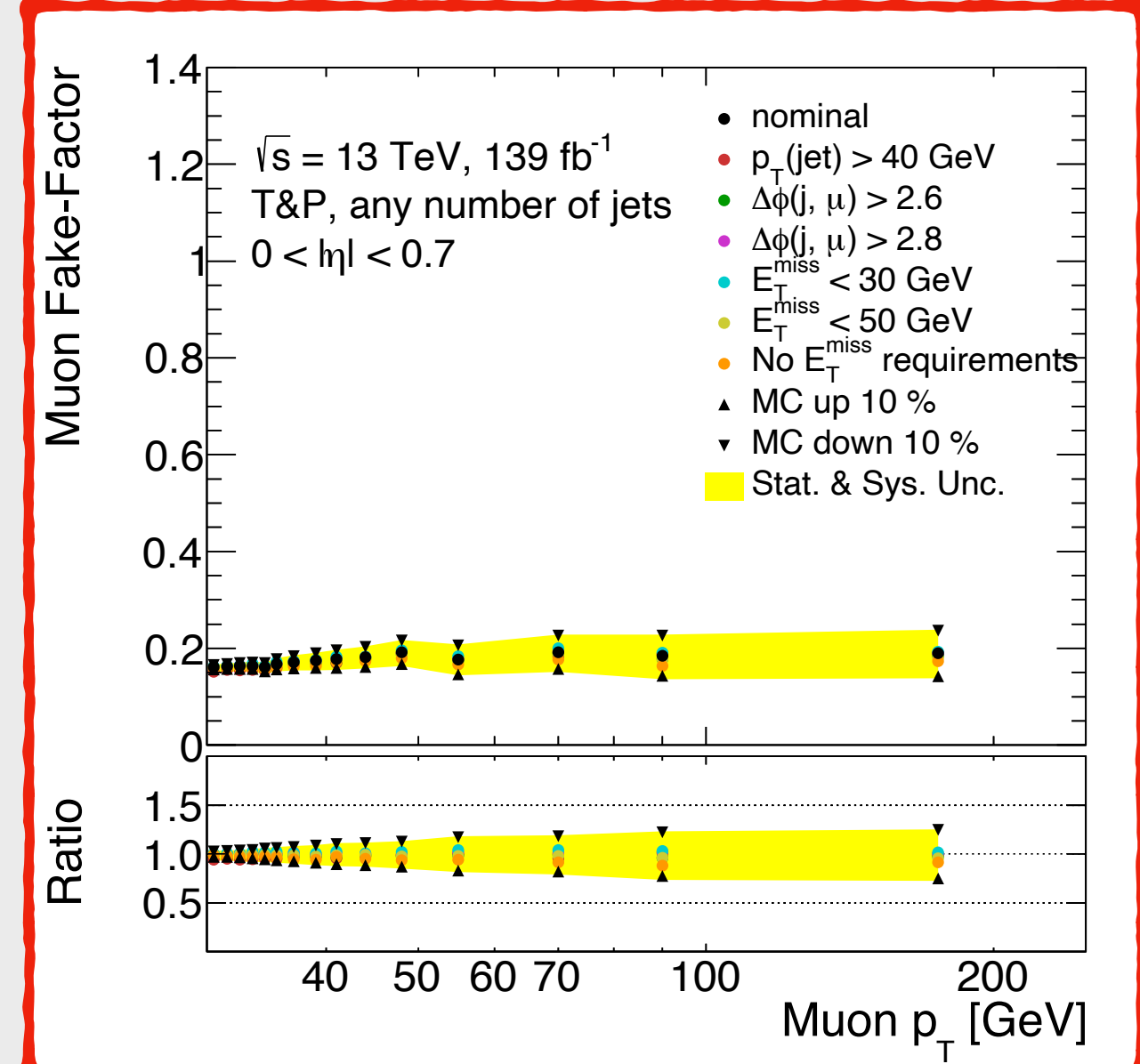
Muon Fakes

- Fake enriched region:
 - exactly 1 baseline muon,
 - any number of jets,
 - b - jet veto,
 - $E_T^{miss} < 40 \text{ GeV}, p_T(j) > 35 \text{ GeV}, \Delta\phi(\mu, j) > 2.7.$

- Five η bins.

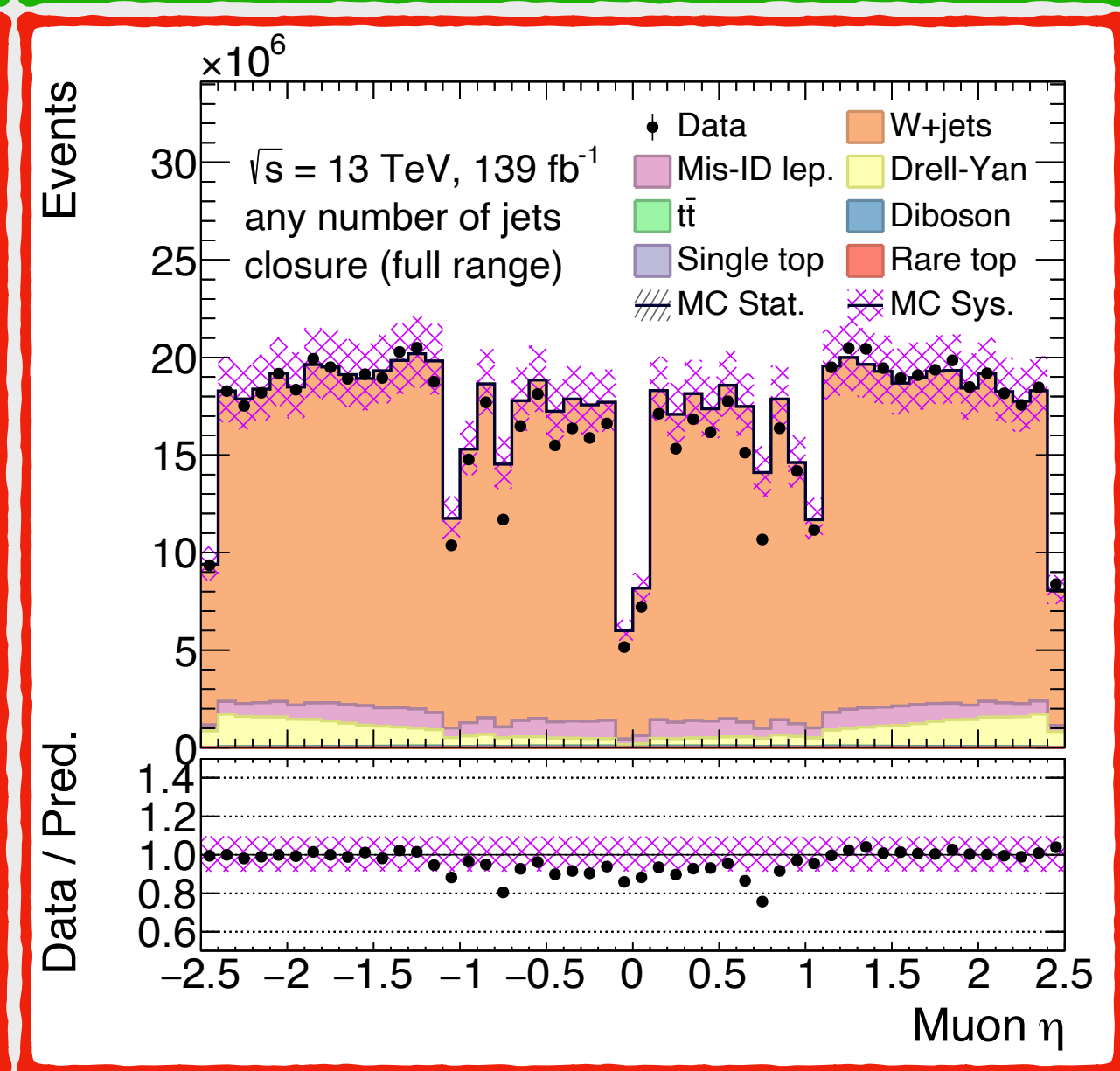
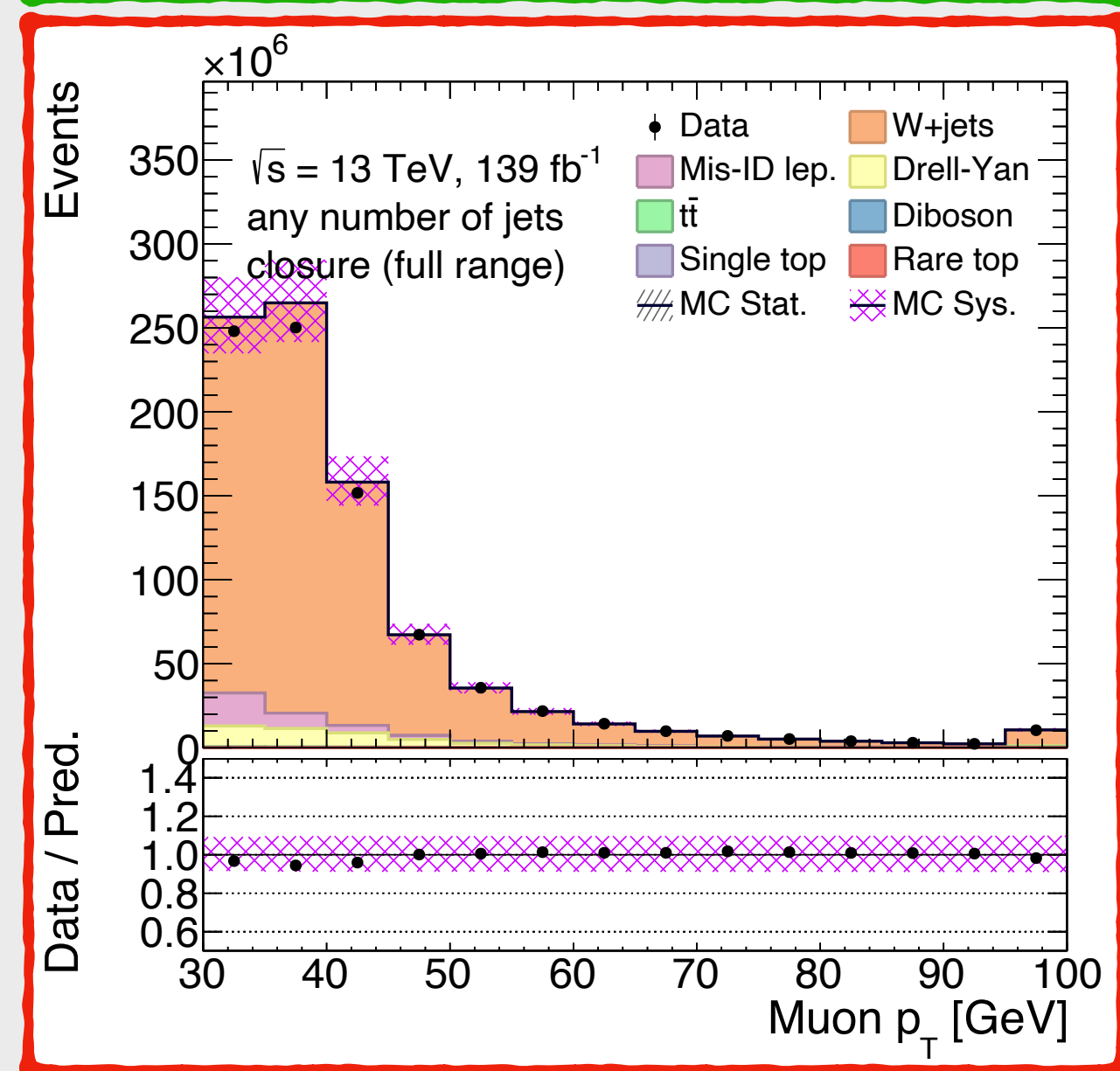
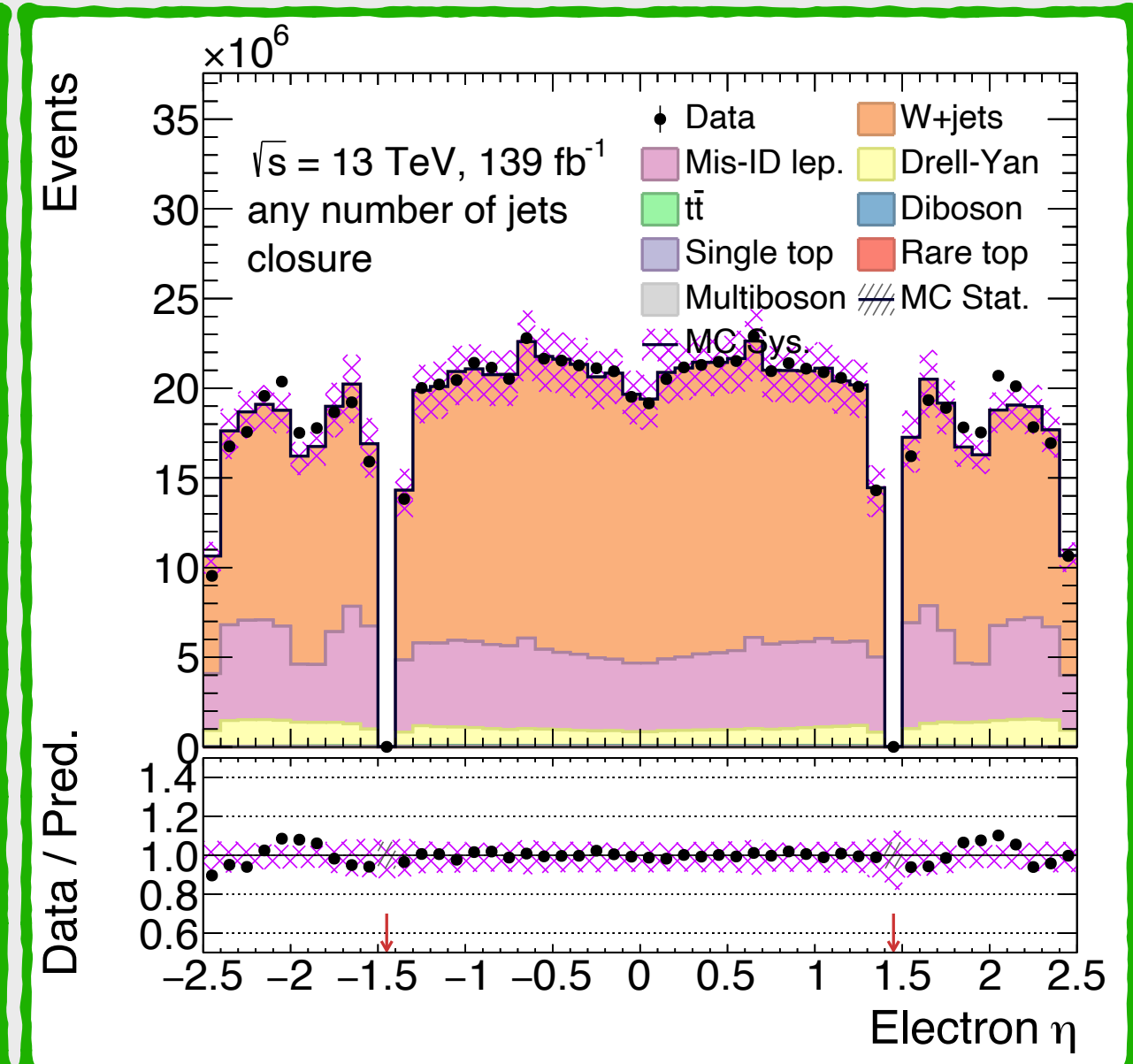
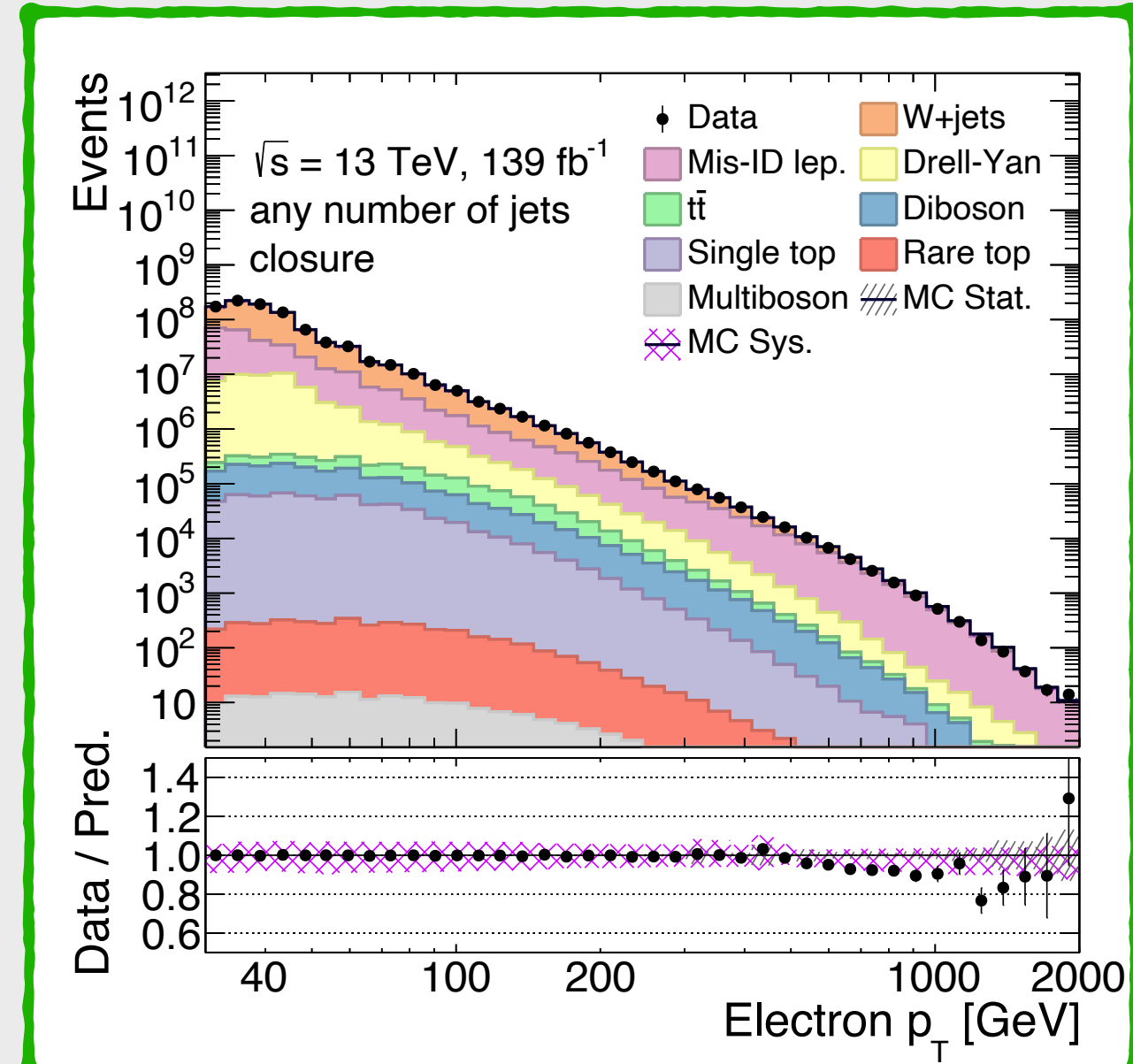
- Triggers used:

Trigger	Average prescale	Period
HLT_mu24	49.36	2015-2018
HLT_mu50	-	2015-2018



Electron and Muon Fake Factor Closure

- We see good agreement in the whole p_T range.
- Muon fakes mainly present in low p_T region ($\lesssim 100 \text{ GeV}$).
- Systematic uncertainty only from fakes and W +jets normalisation in muon case.

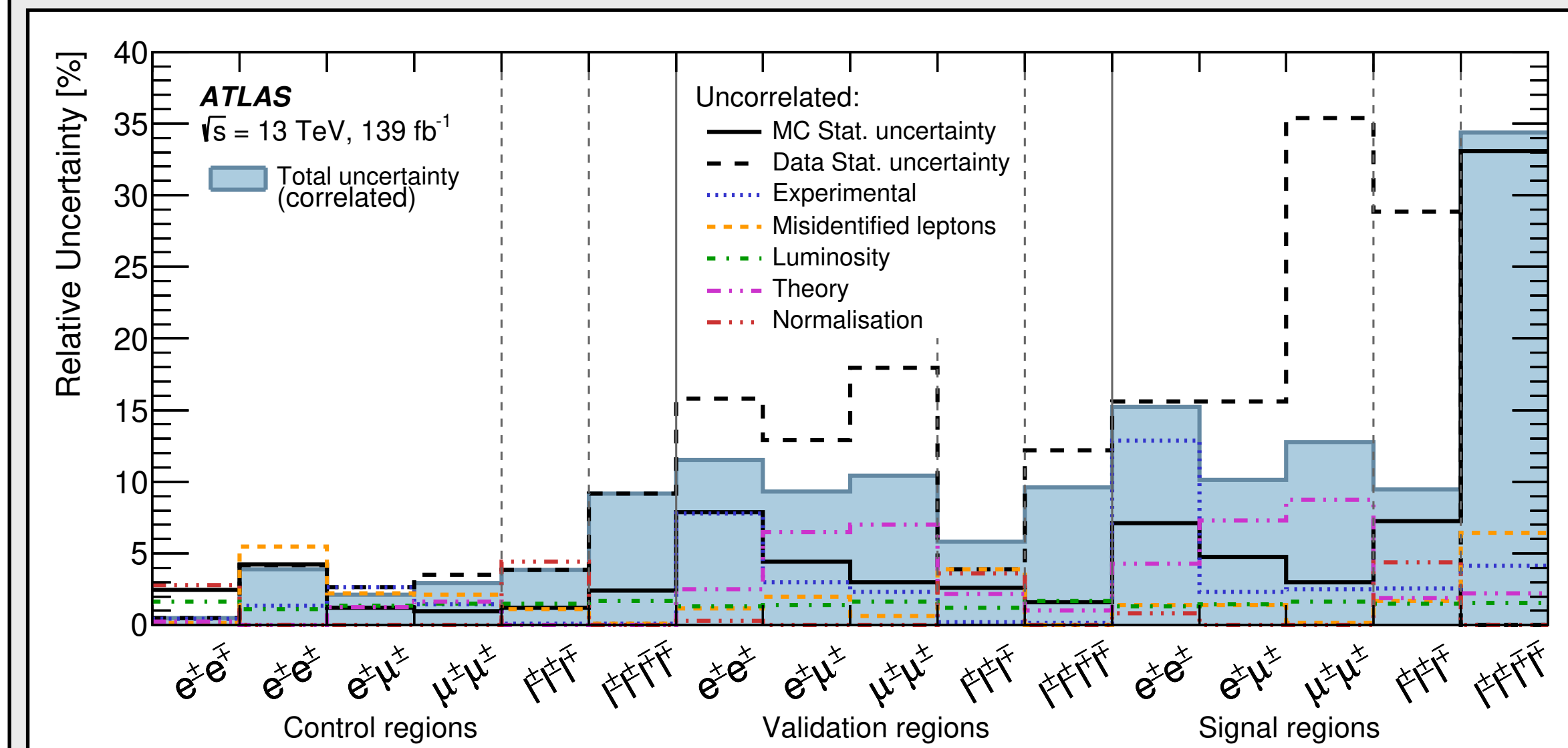
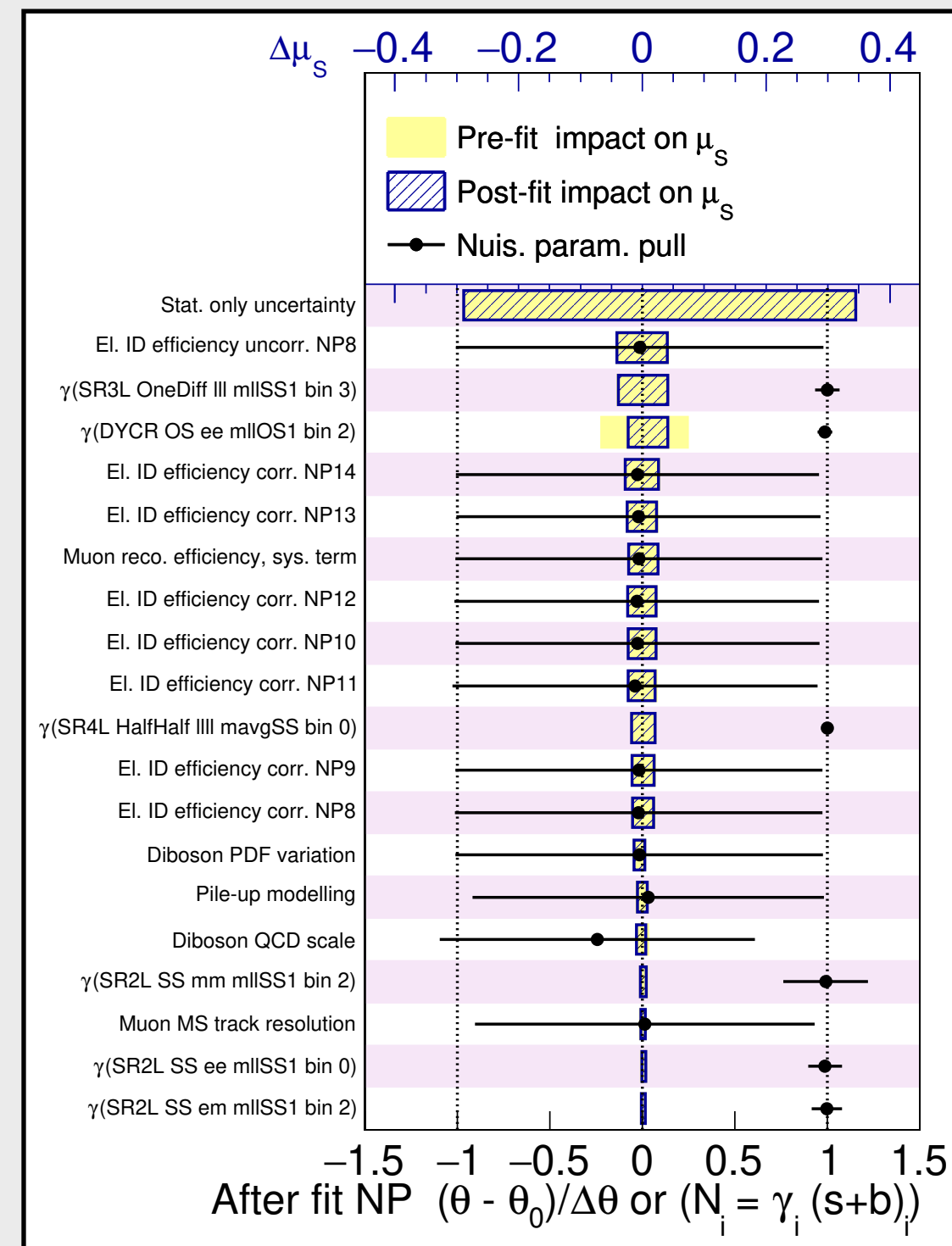
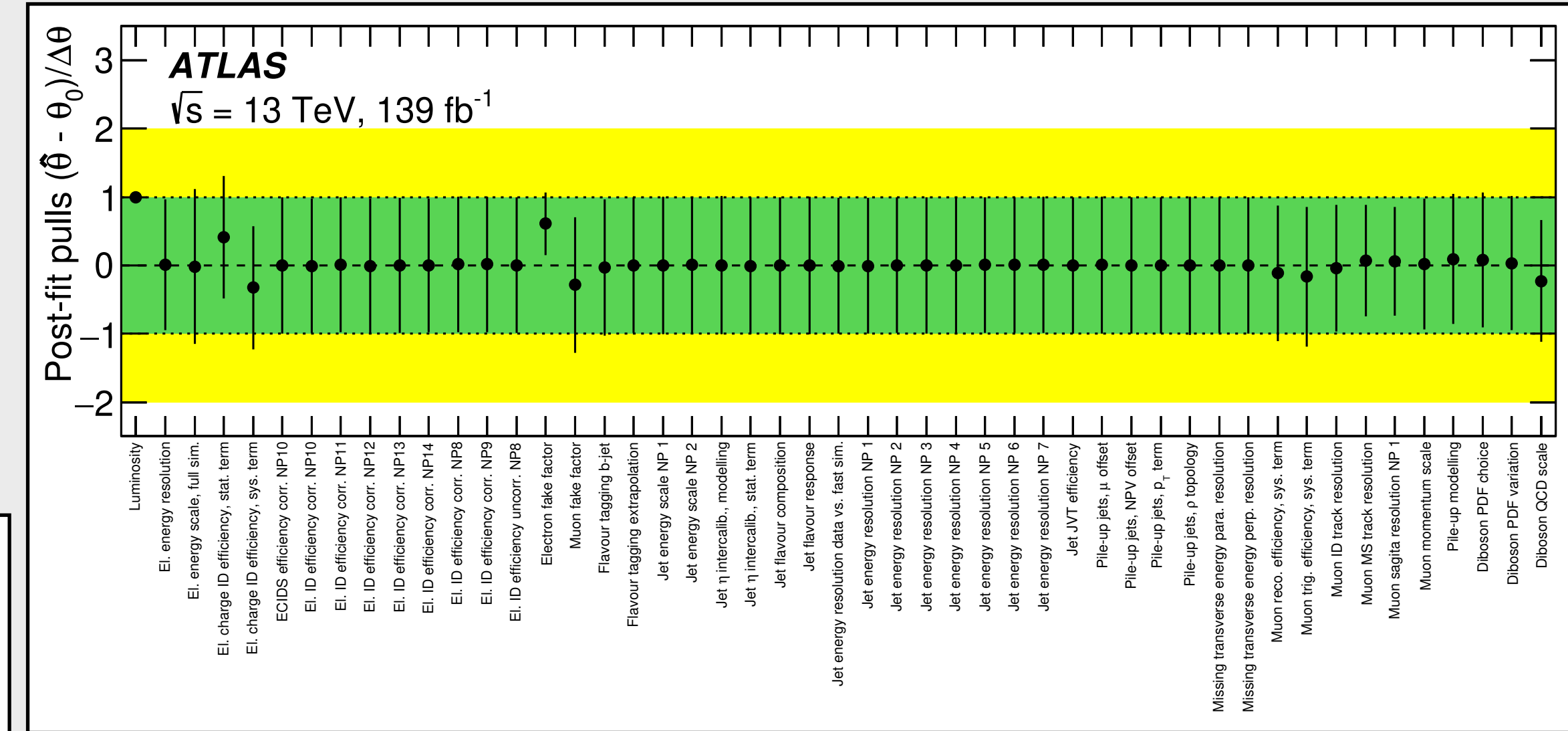




Results of the cut-based analysis

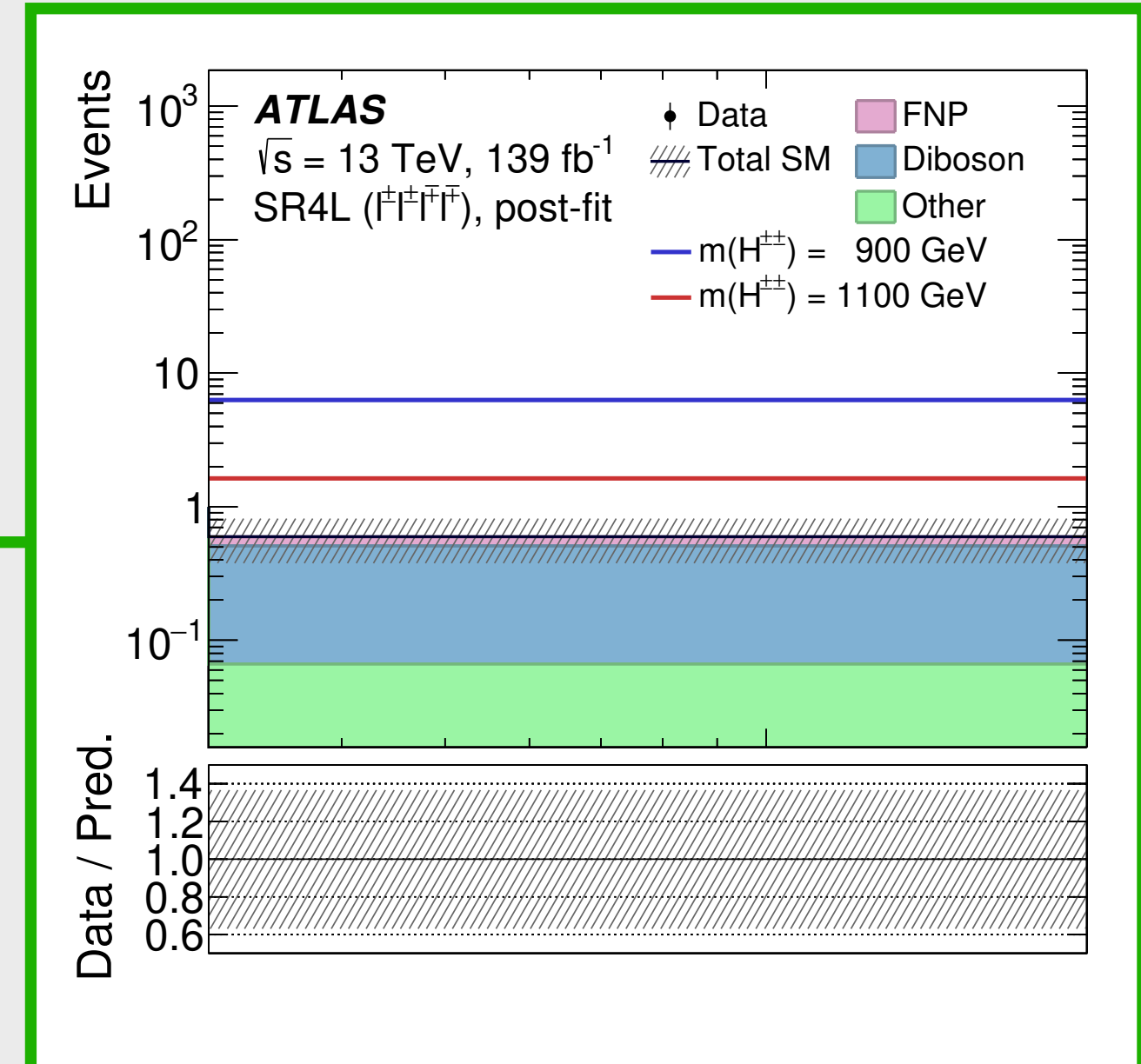
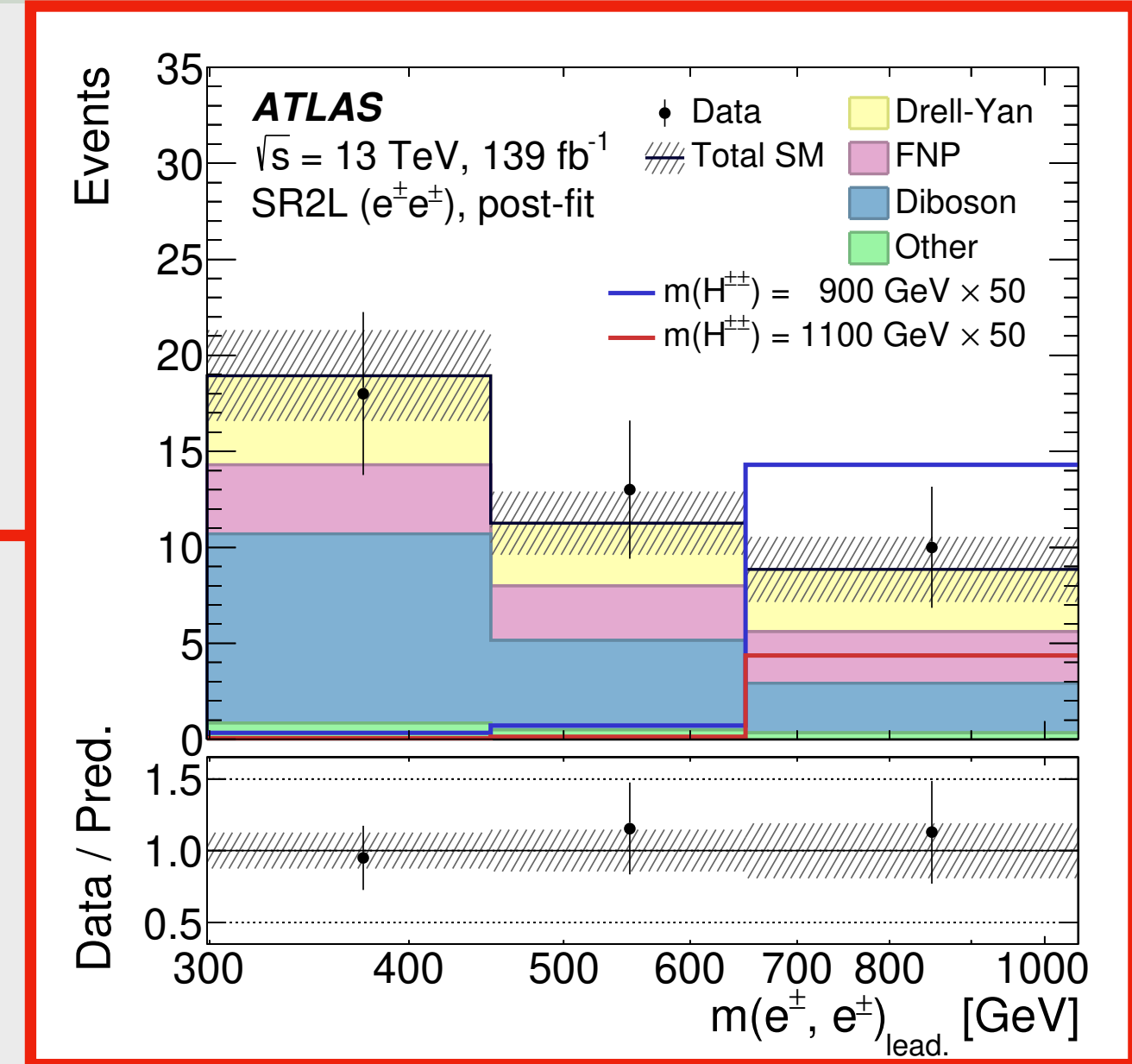
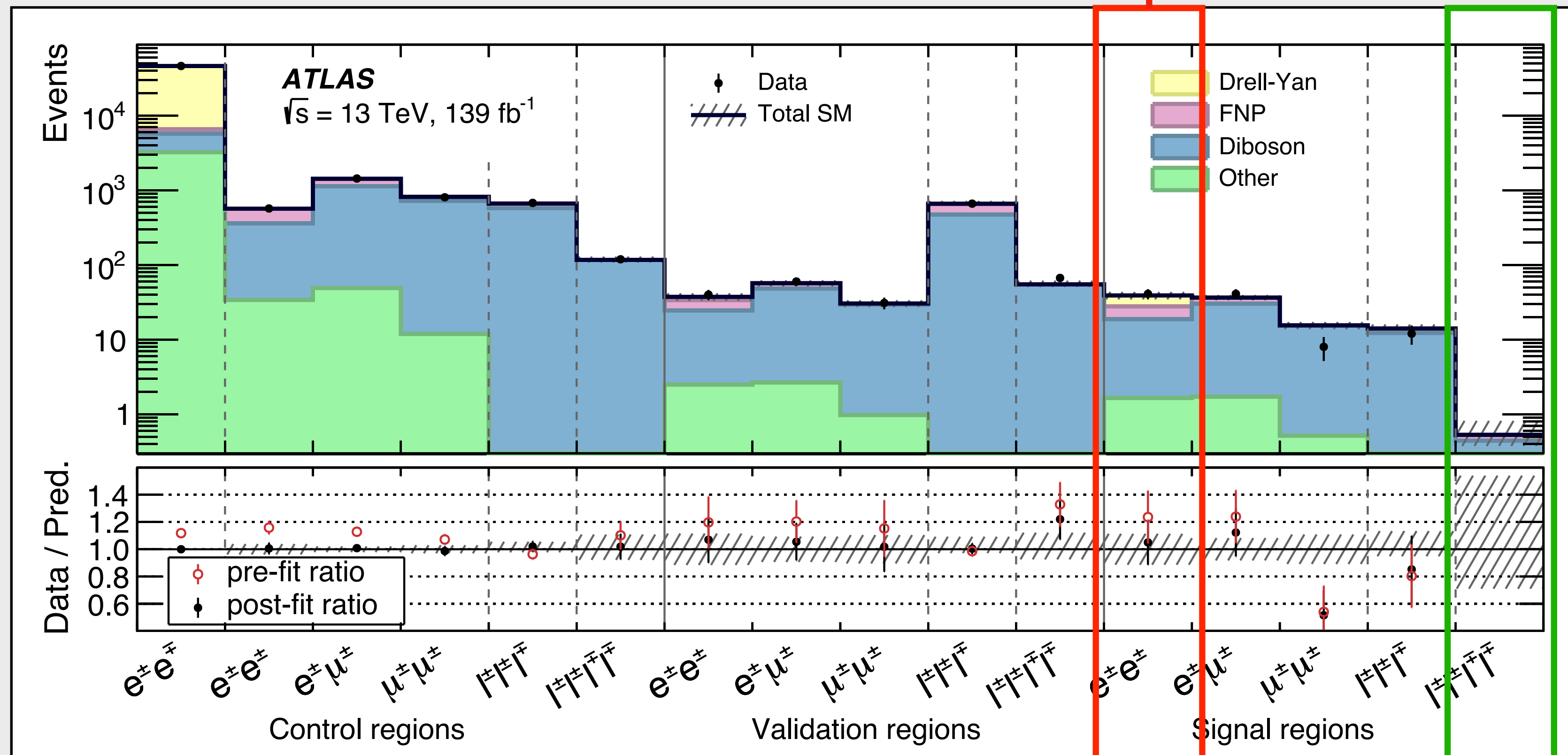
Systematic uncertainties and statistical analysis

- Both theoretical and experimental uncertainties are considered.
- Leading sources of uncertainty:
 - **statistical uncertainty,**
 - electron identification efficiency,
 - fake background estimation.
- Number of nuisance parameters (NPs) used: **254.**
- After using standard HistFitter pruning only **49** NPs remain.



Yields and distributions

- Overall good post-fit agreement between expected and observed events.
- Unfortunately, no data event passed the SR4L requirements.





Results of the multivariate analysis

Systematic uncertainties and statistical analysis

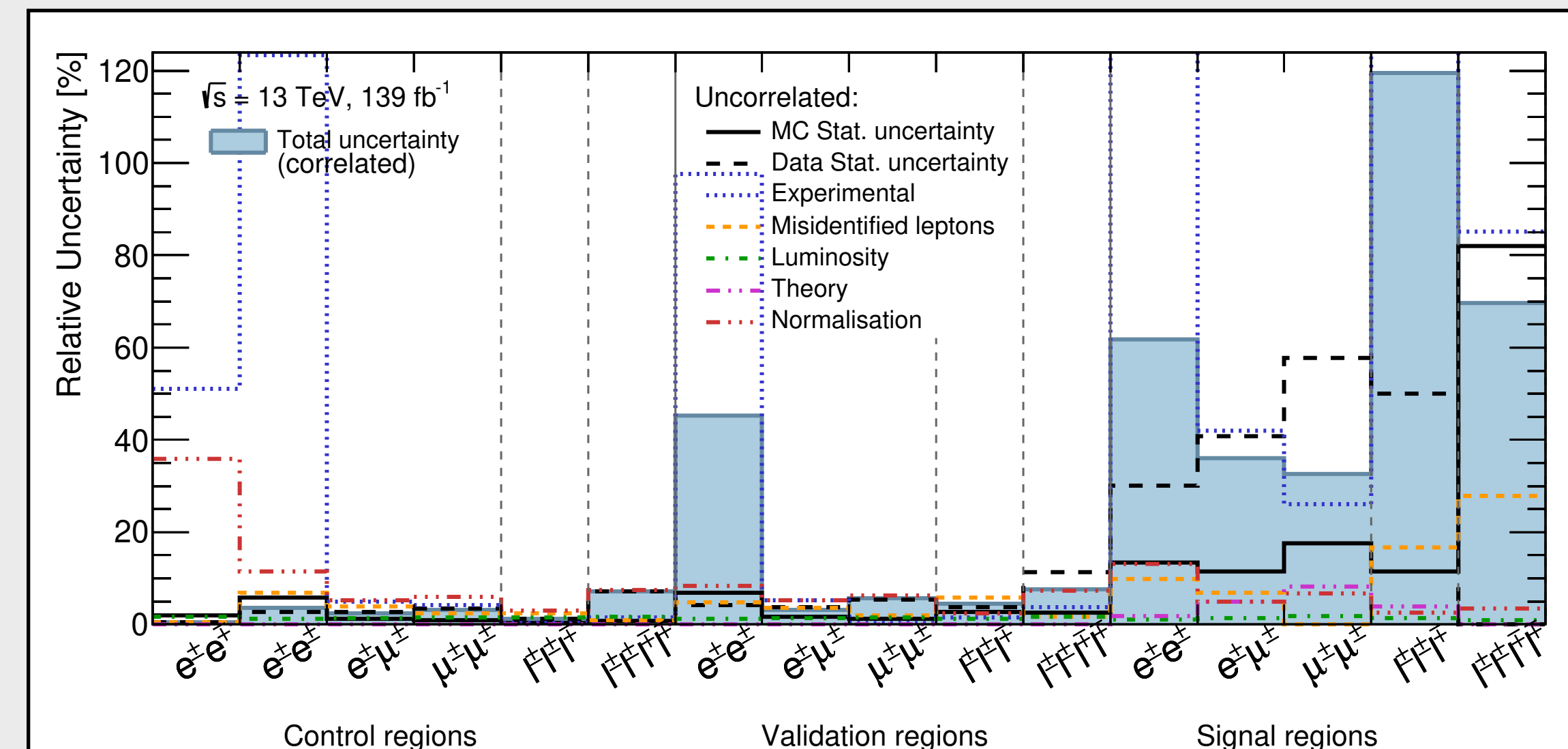
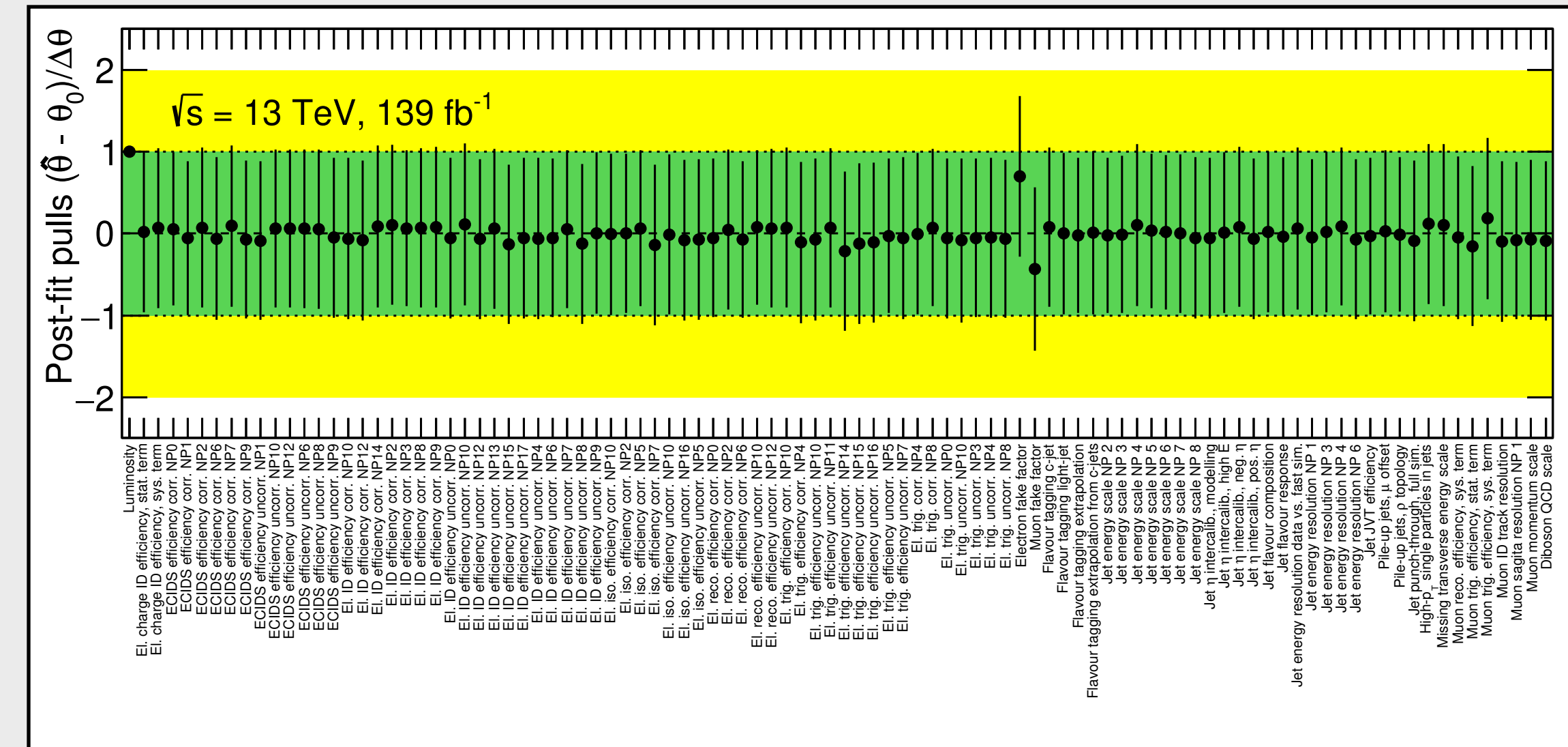
- Systematic uncertainties are **much harder** to keep **under control**:

- pull plots are more scattered,
- more NPs remain after pruning,
- uncertainties exceed 100% in some regions.

- Leading sources of uncertainty:

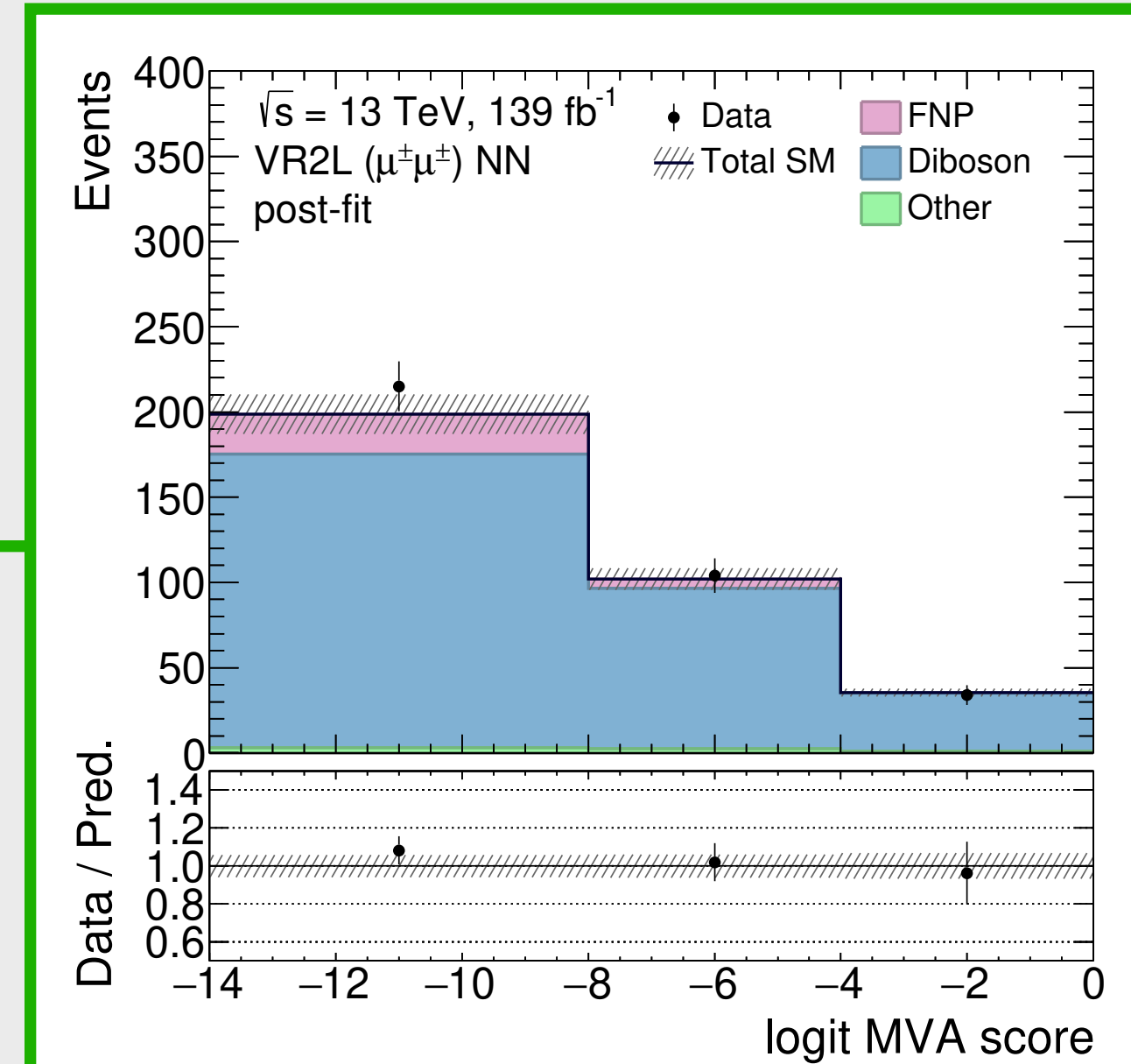
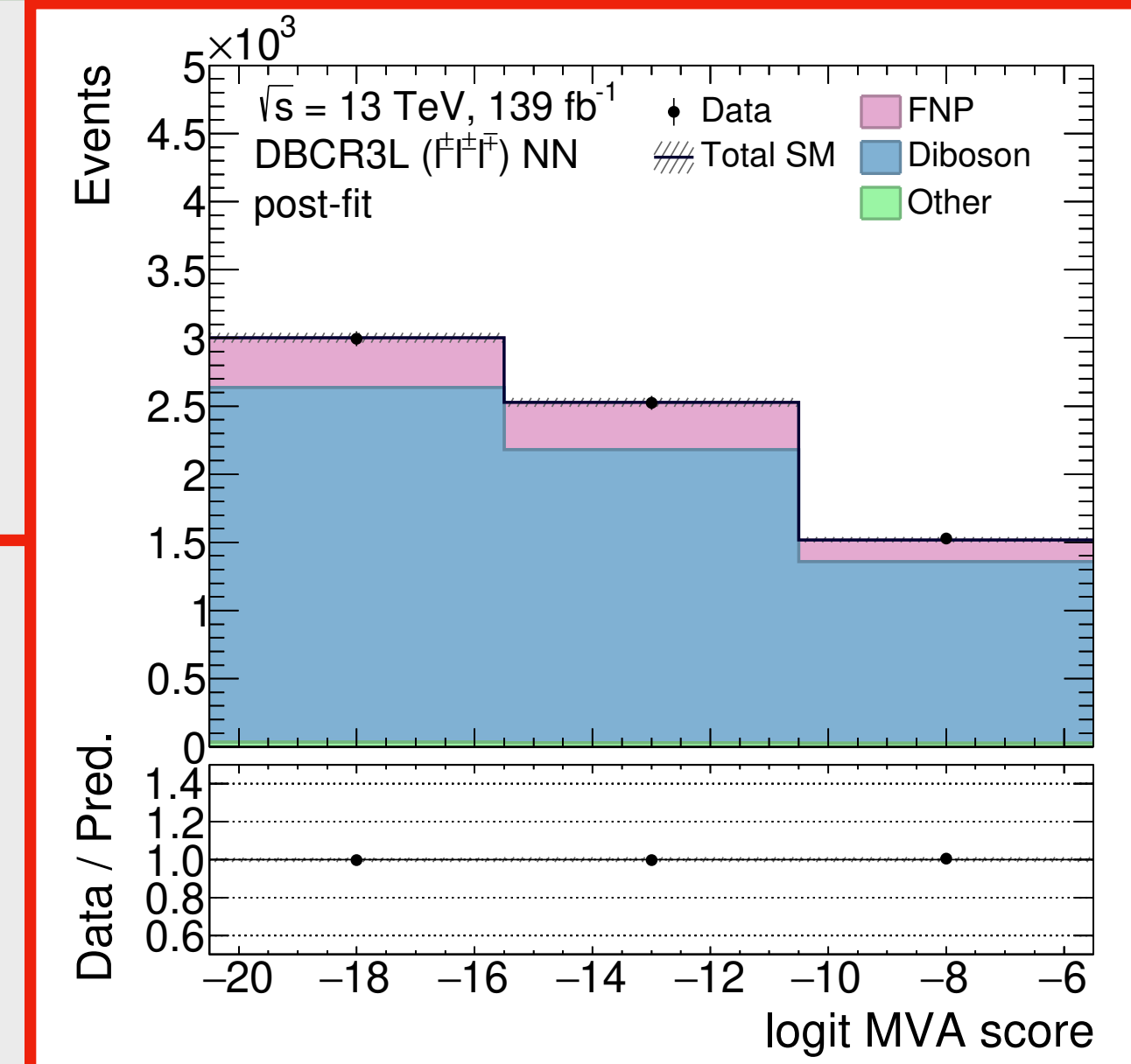
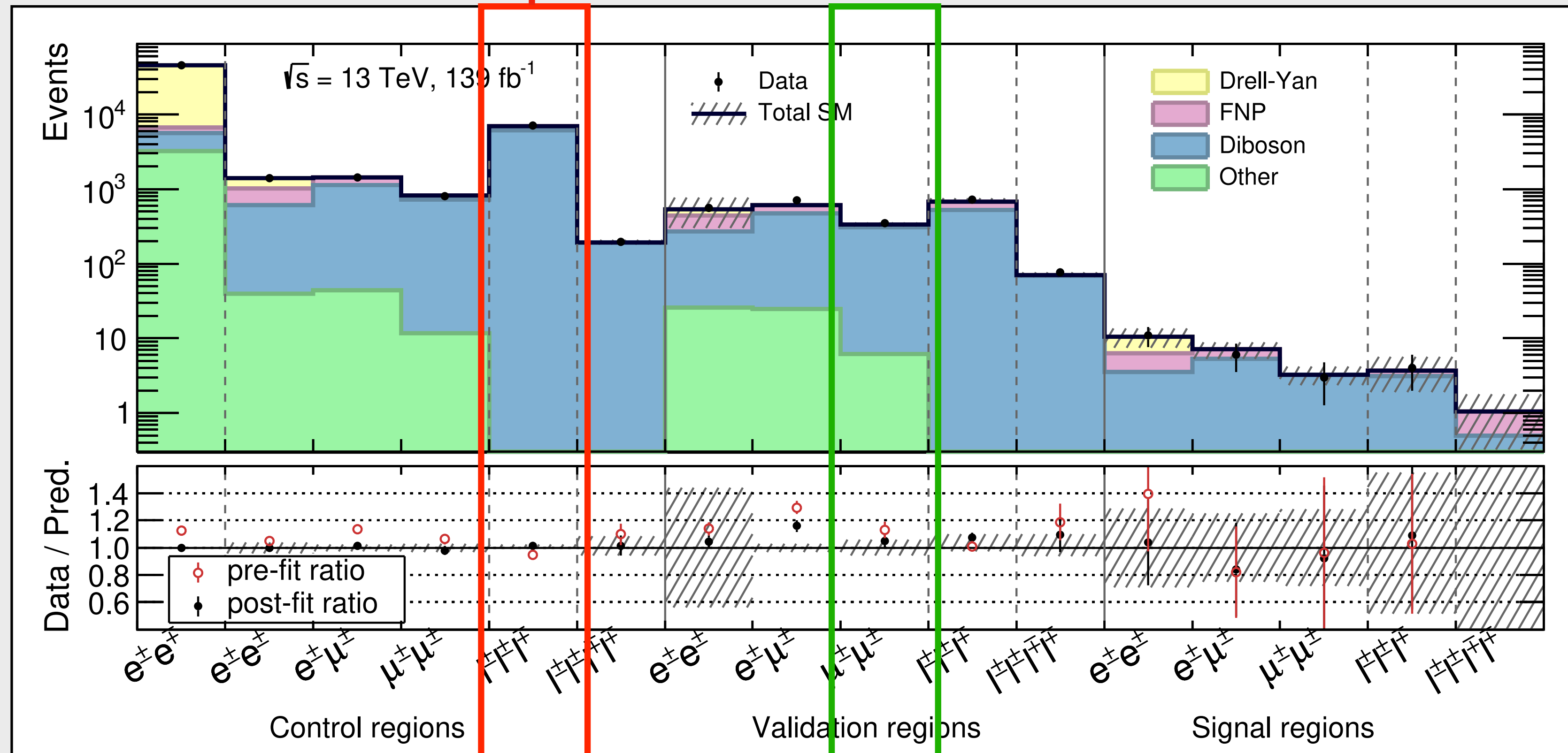
- statistical uncertainty**,
- muon fake estimation,
- muon efficiencies.

- Studies indicate that this is due to the sensitive logit MVA score. Slight systematic variation can produce very different output.

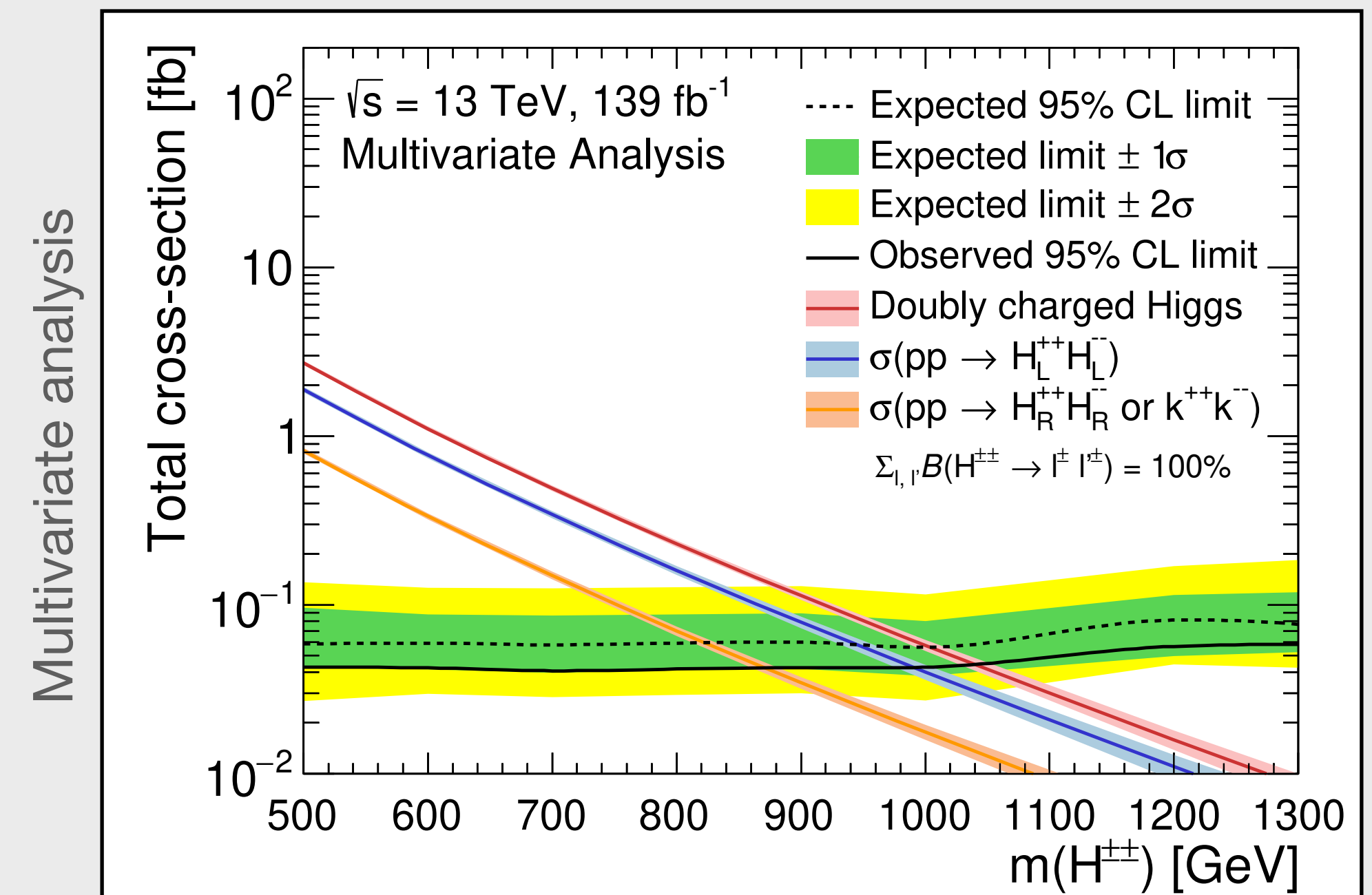
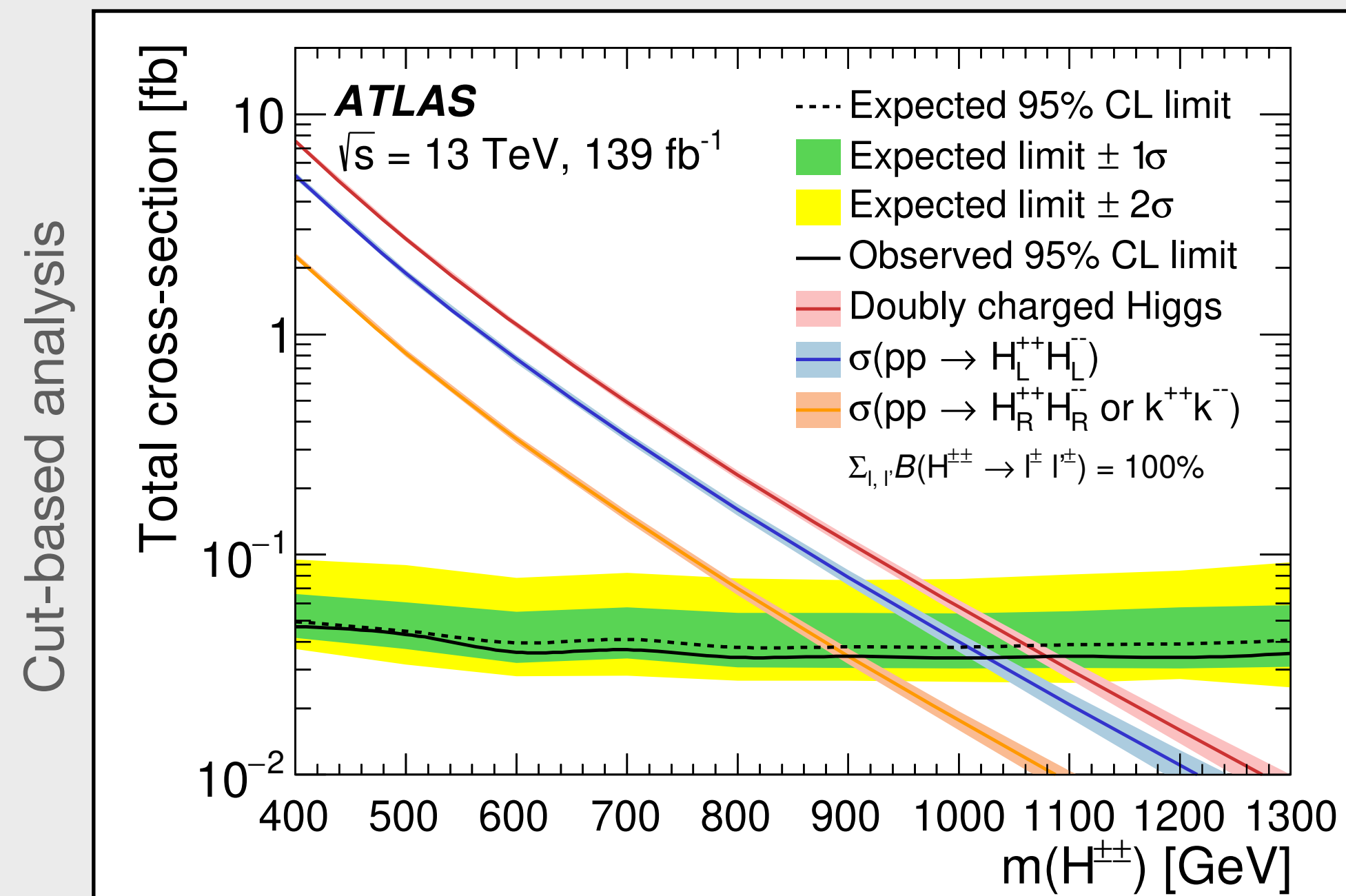


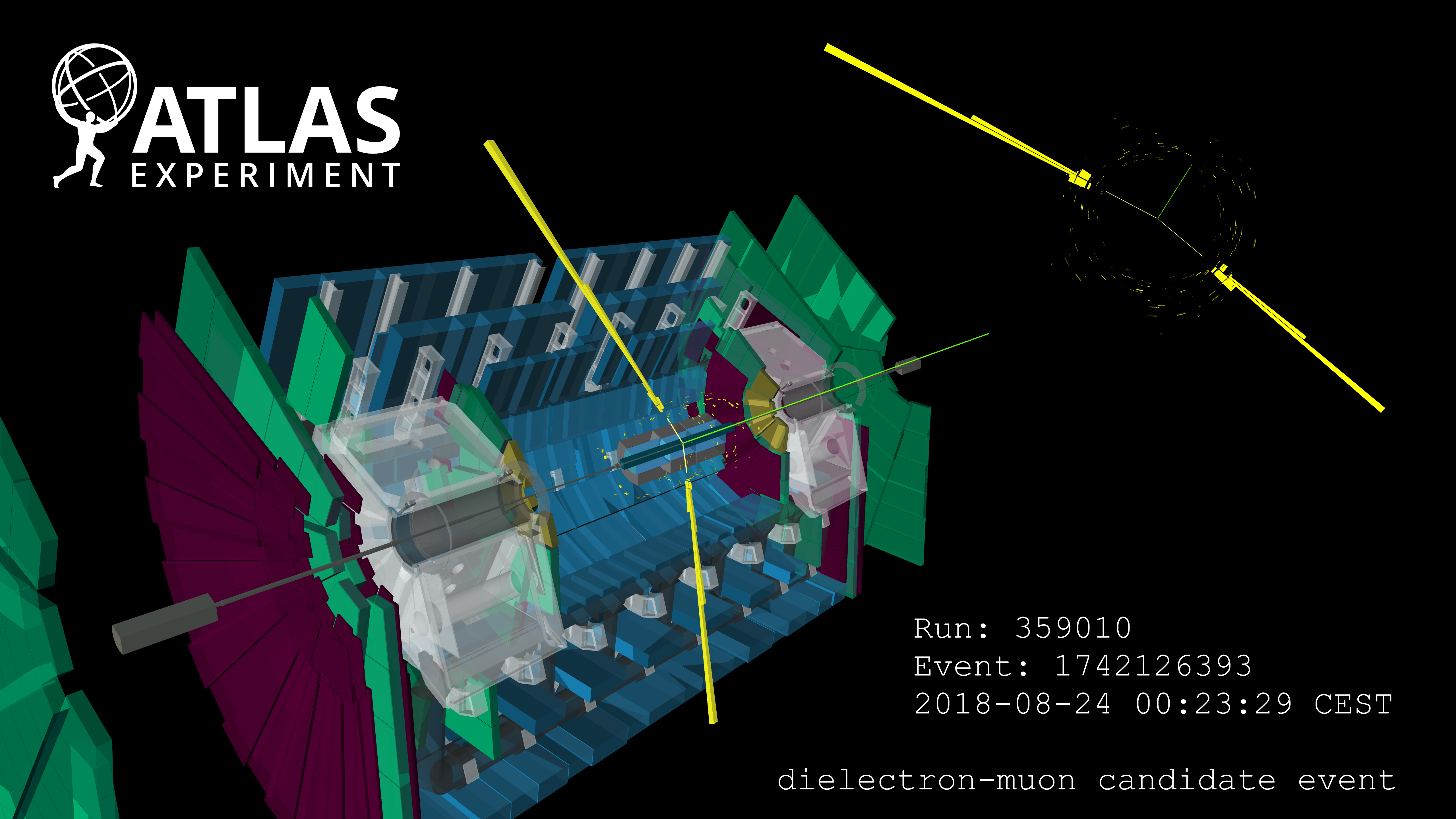
Yields and distributions

- Overall good post-fit agreement between expected and observed events with **higher uncertainties**.
- No data event passed the SR4L requirements in the MVA analysis, too.



- The expected exclusion limit is 1065^{+30}_{-50} GeV for LRSM and 880^{+30}_{-40} GeV for the Zee-Babu model.
- The observed lower limit on the $H^{\pm\pm}$ mass reaches **1080 GeV** and **900 GeV** when combining all three channels for LRSM and the Zee-Babu model, respectively.
- The four-lepton channel limit is the strongest and drives the combined result.
- MVA did **not** provide **much improvement** with respect to the cut-based analysis.





Run: 359010

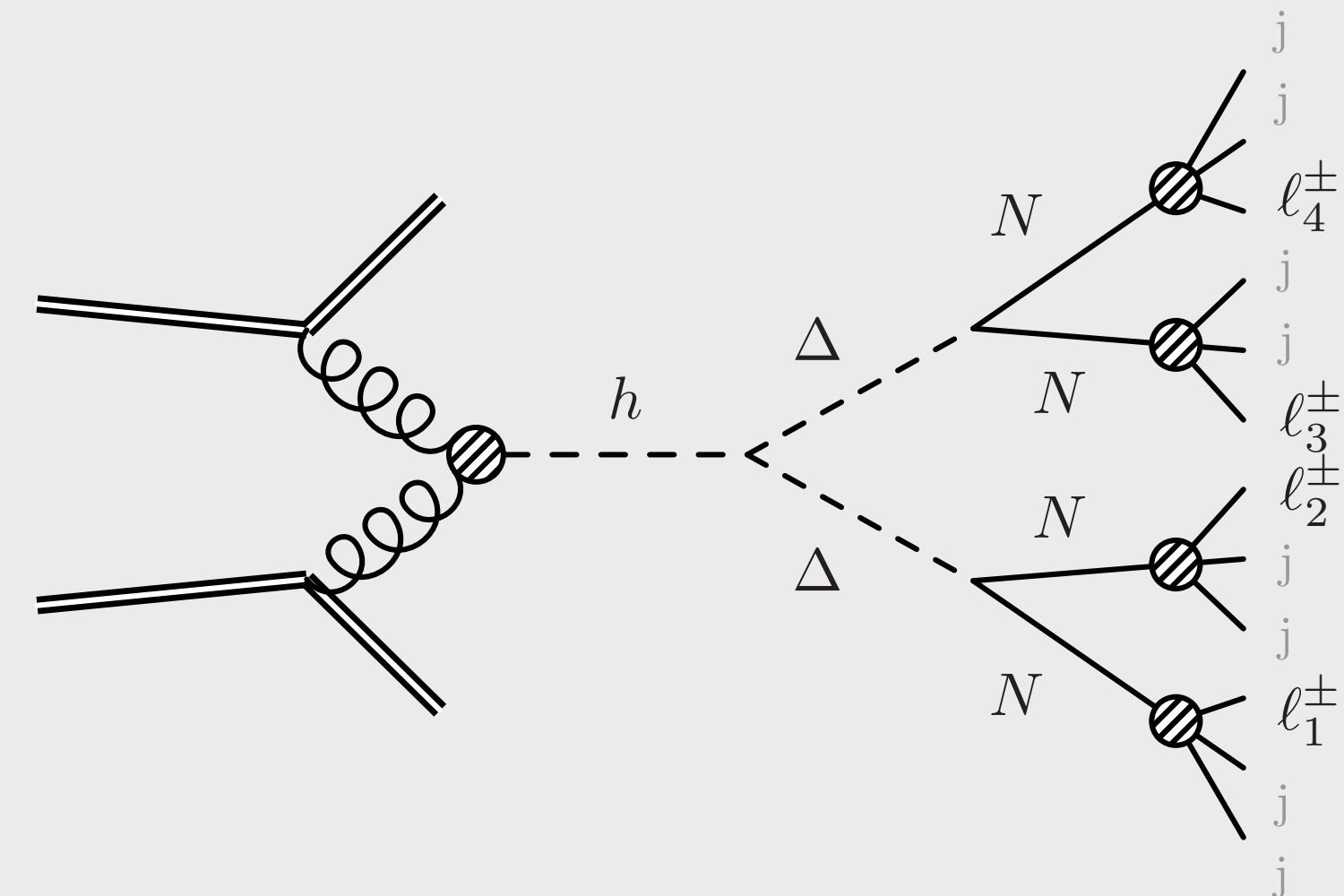
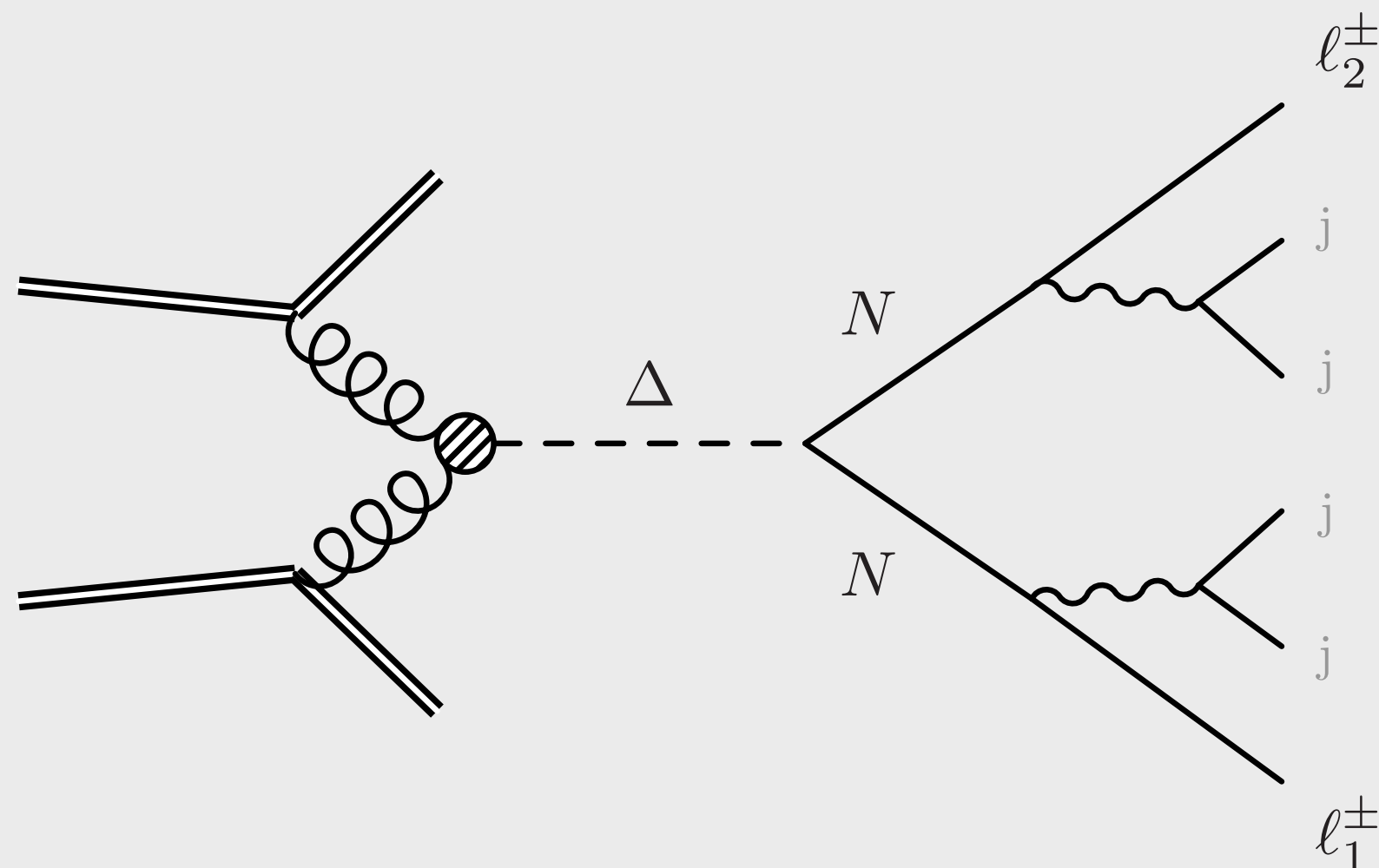
Event: 1742126393

2018-08-24 00:23:29 CEST

dielectron-muon candidate event

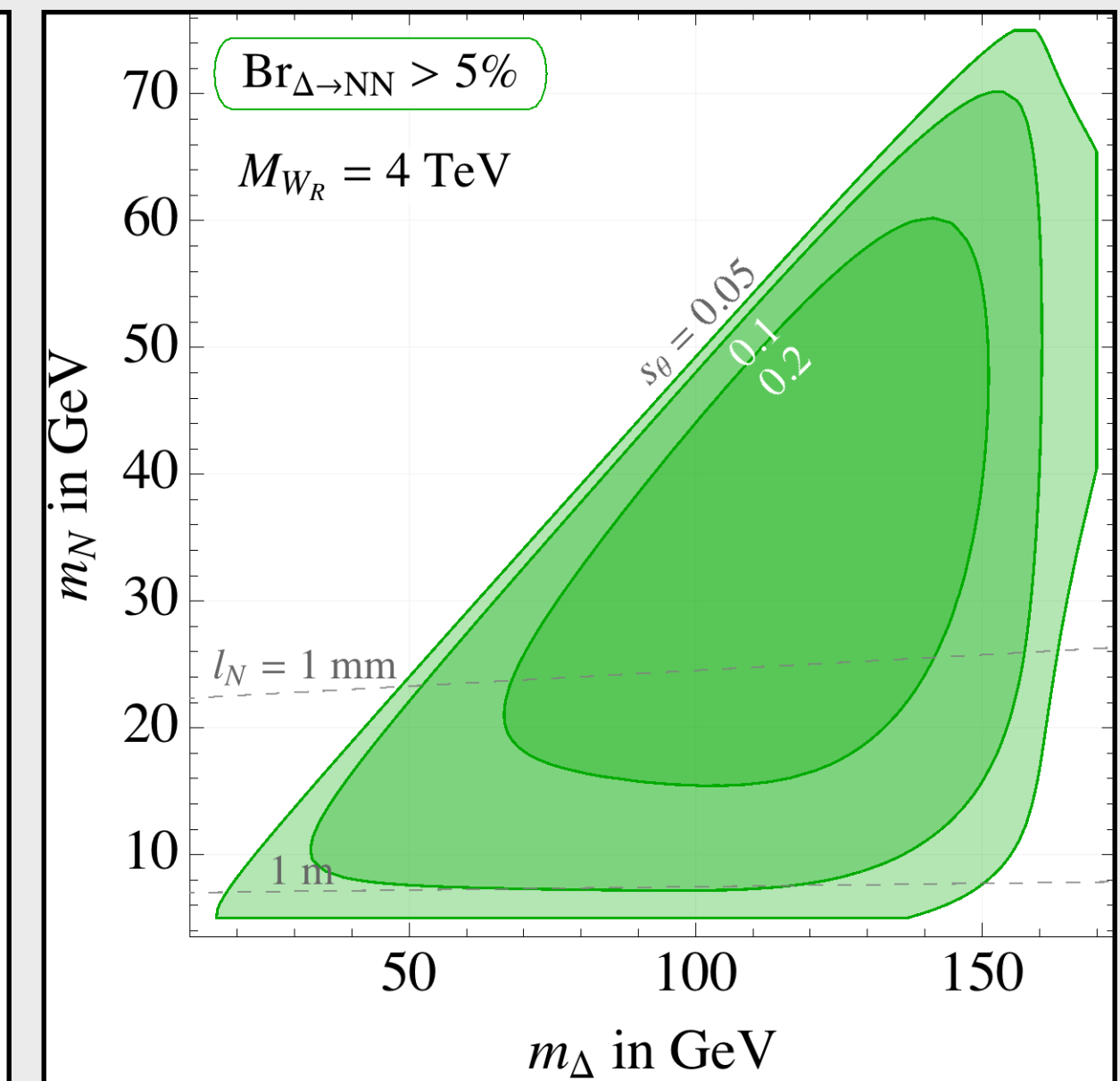
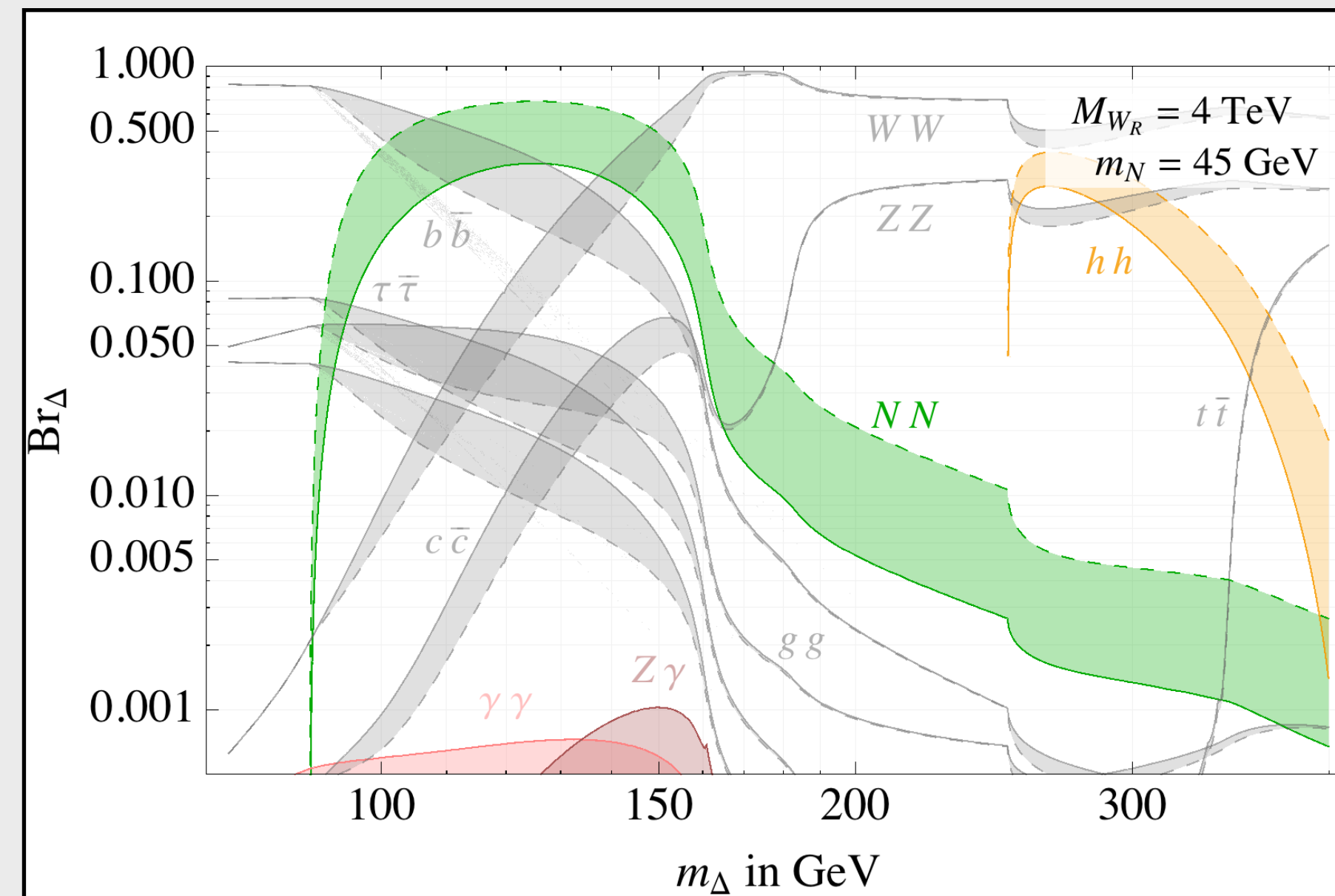
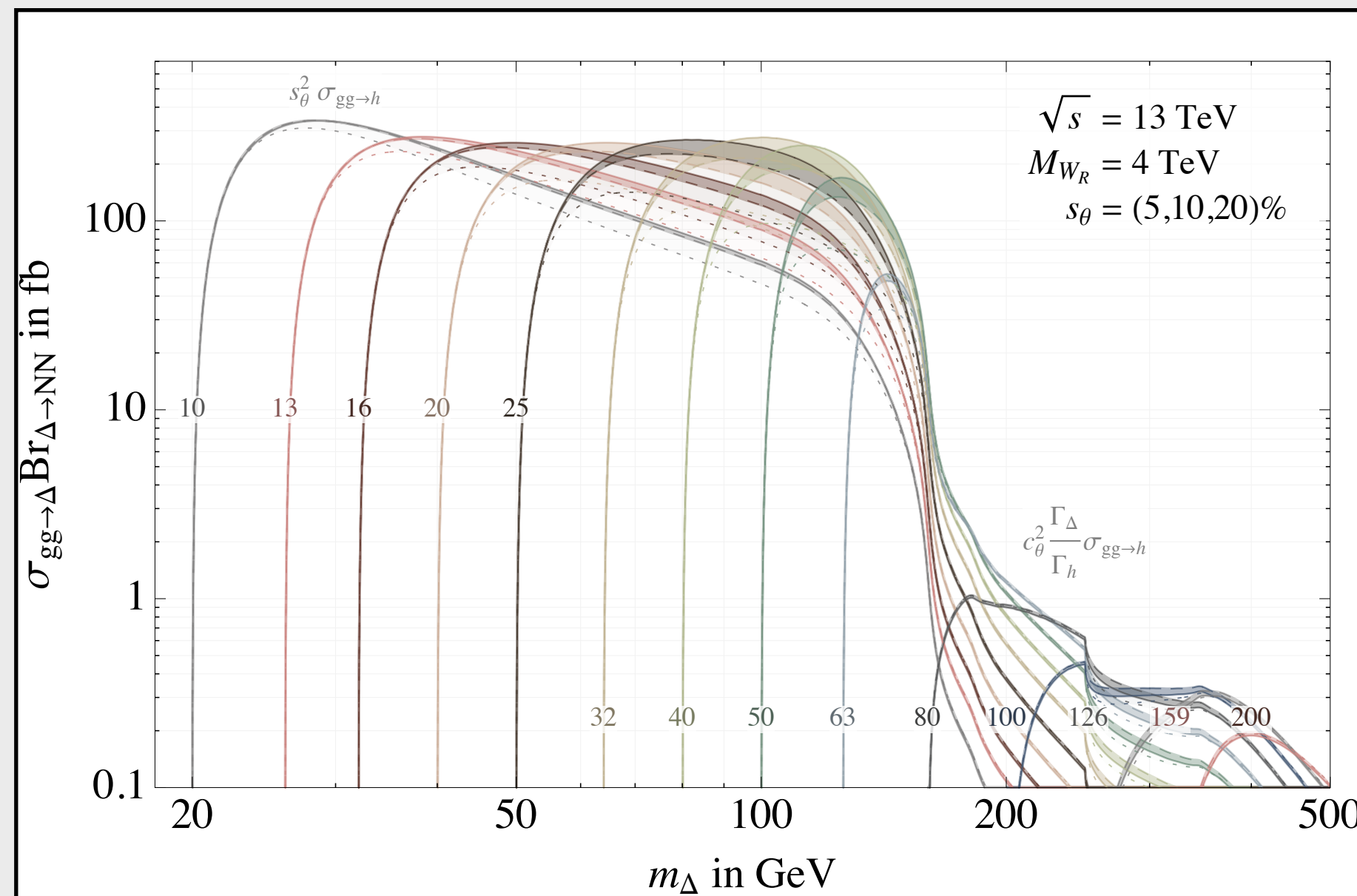
LNV Higgs Decays

- Lepton Number Violating Higgs decay is a signal within the Left-Right Symmetric Model that can create a (same-sign) two-lepton final state.
- The Higgs can act as a gateway to the **origin of heavy Majorana neutrino mass**. This process is complementary to the existing nuclear and collider searches for lepton number violation and can probe the scale of parity restoration even above other direct searches.
- We relaxed the “ Δ constraint” and are generating inclusive $gg \rightarrow NN$ process (goes through $\Delta/A/H/Z/Z_R/h$). Gain $O(20\%)$ in production cross-section.

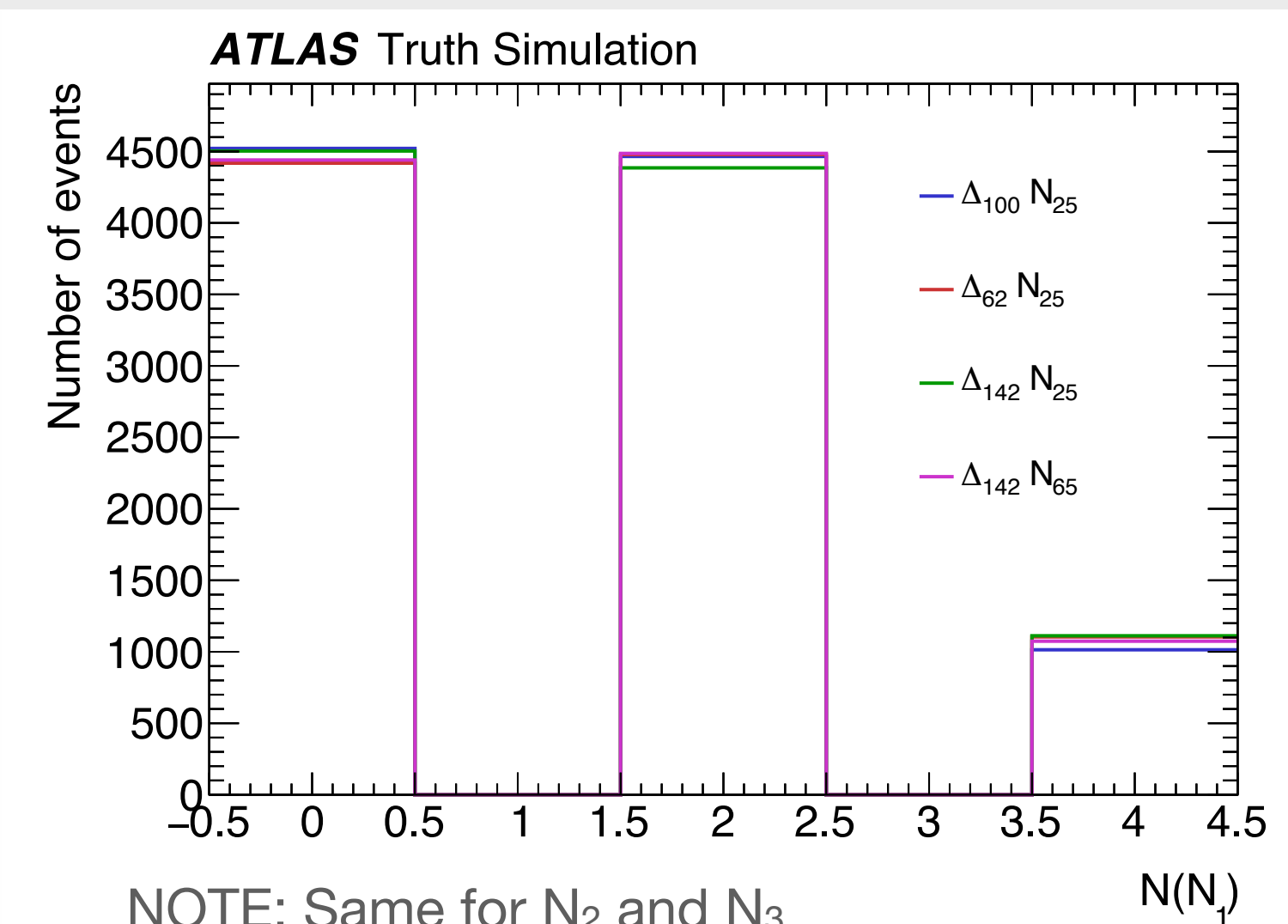
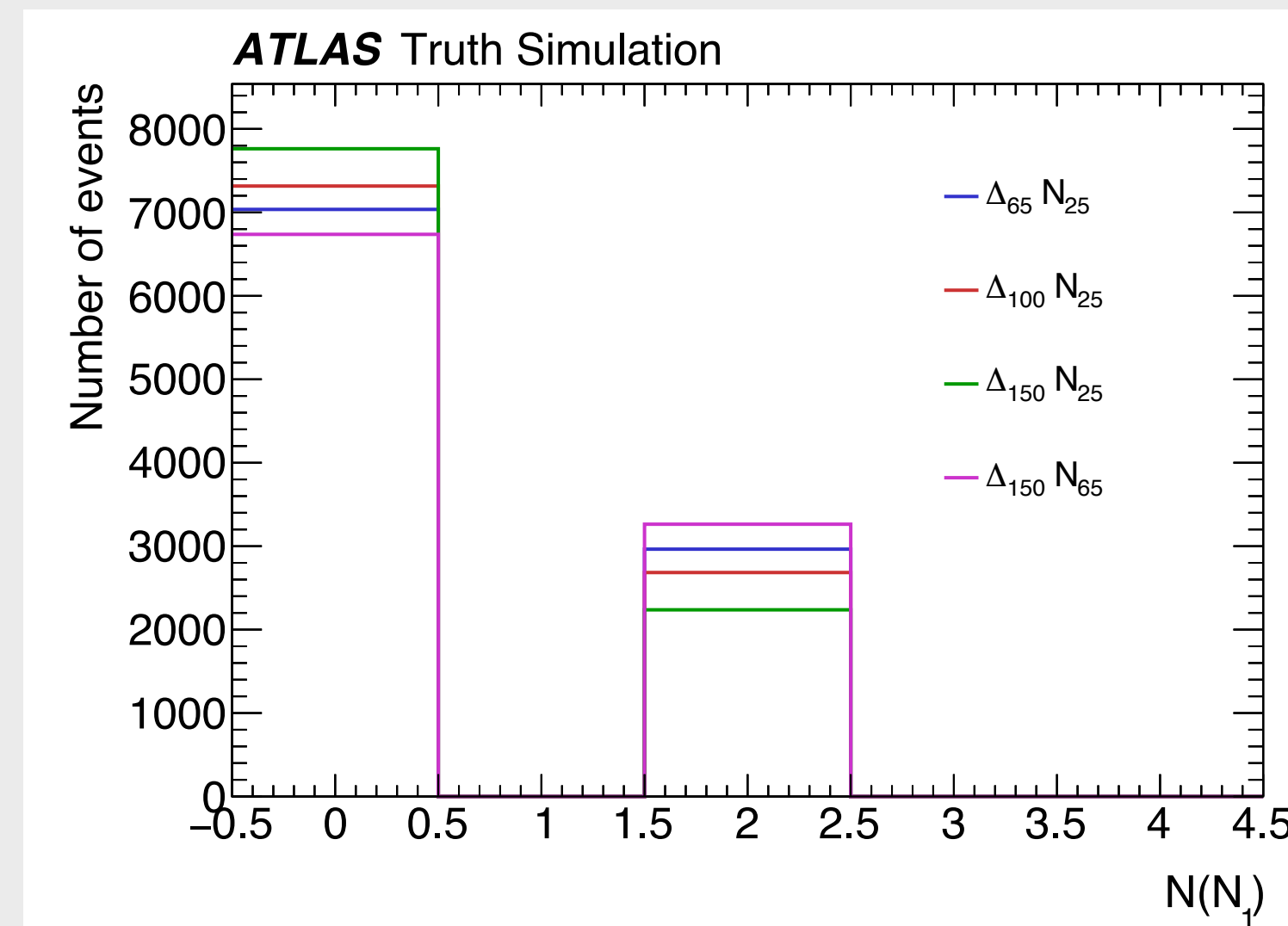
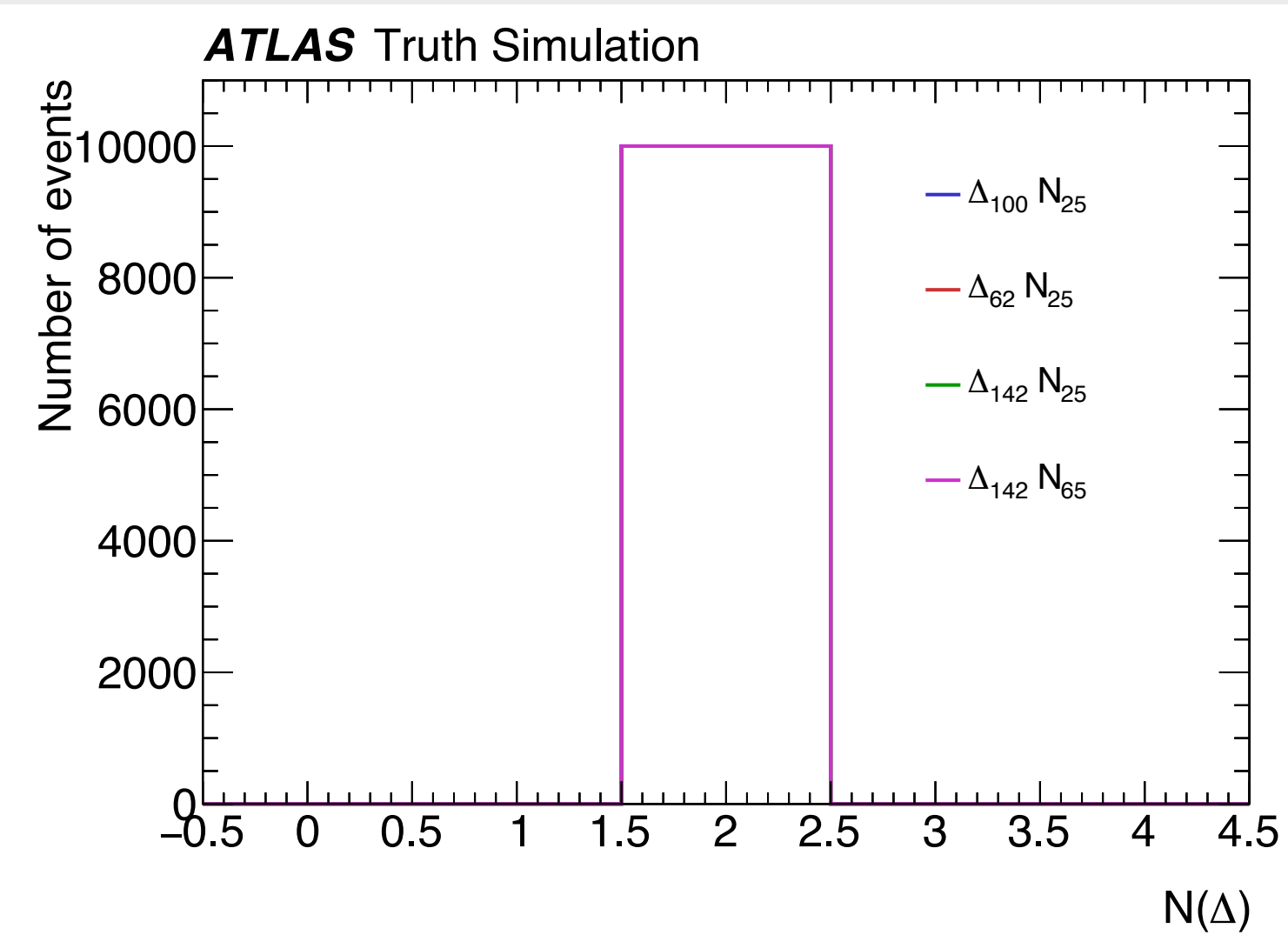
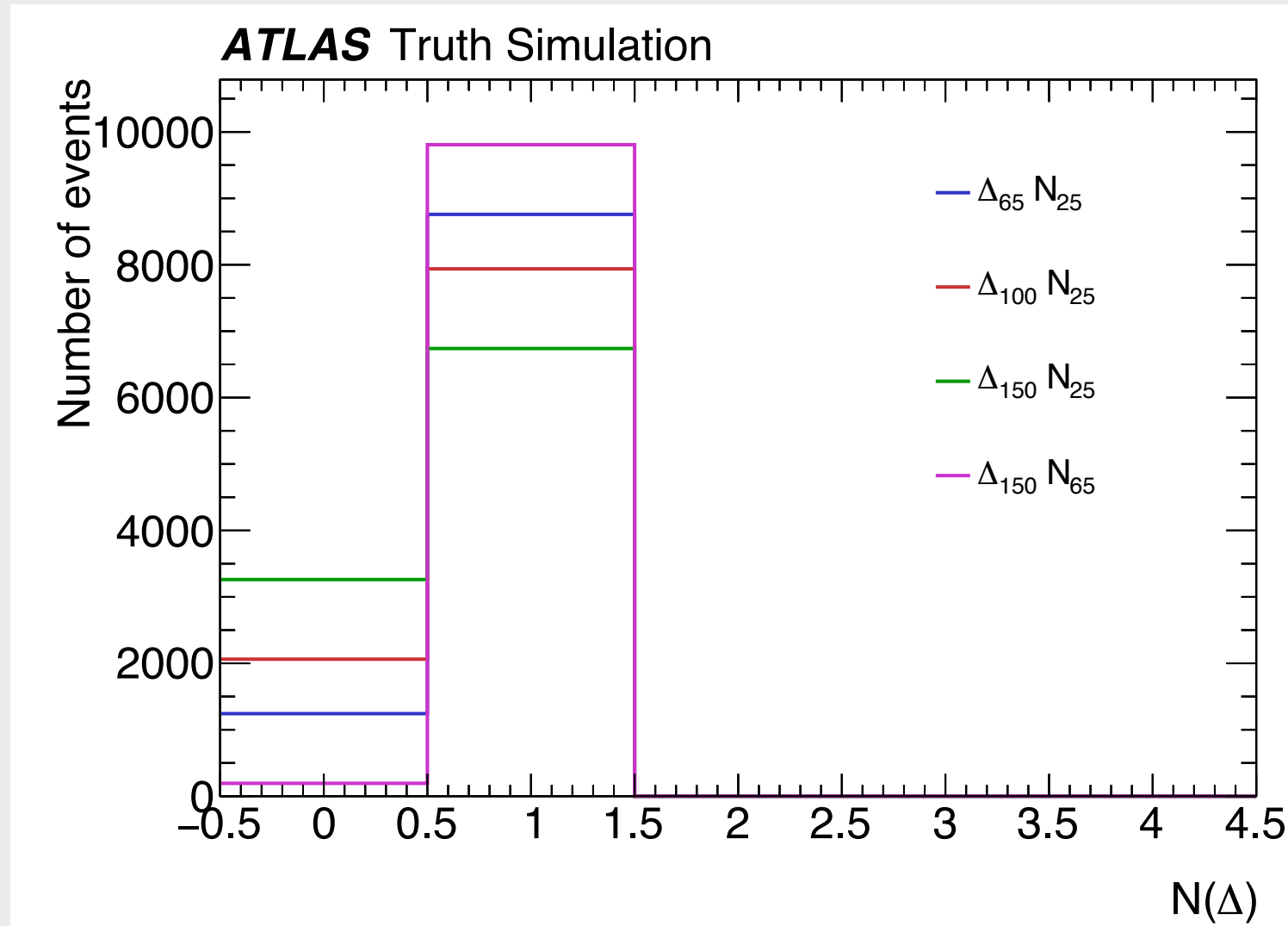


- Mass points were chosen based on [arXiv:1612.06840](https://arxiv.org/abs/1612.06840) studies. **Masses** of all three heavy neutrino flavours **are degenerate**.

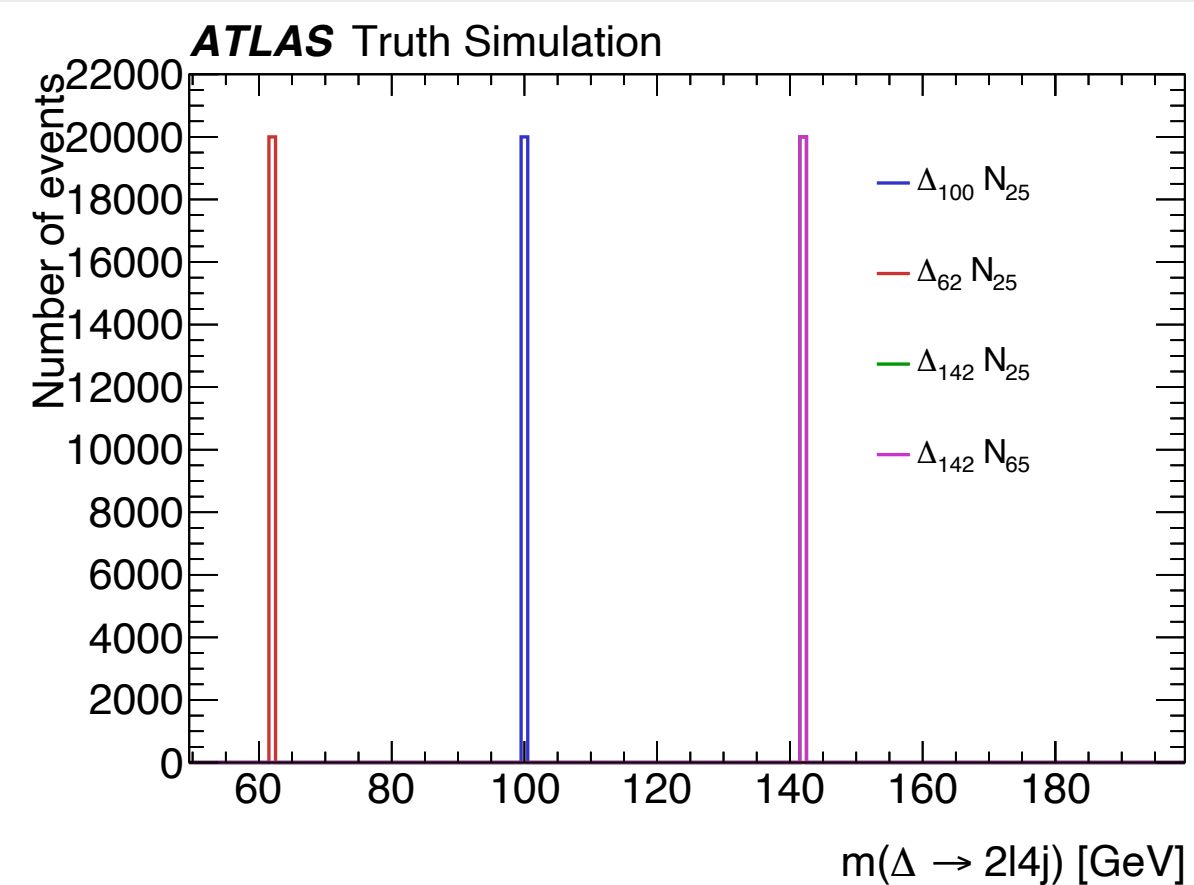
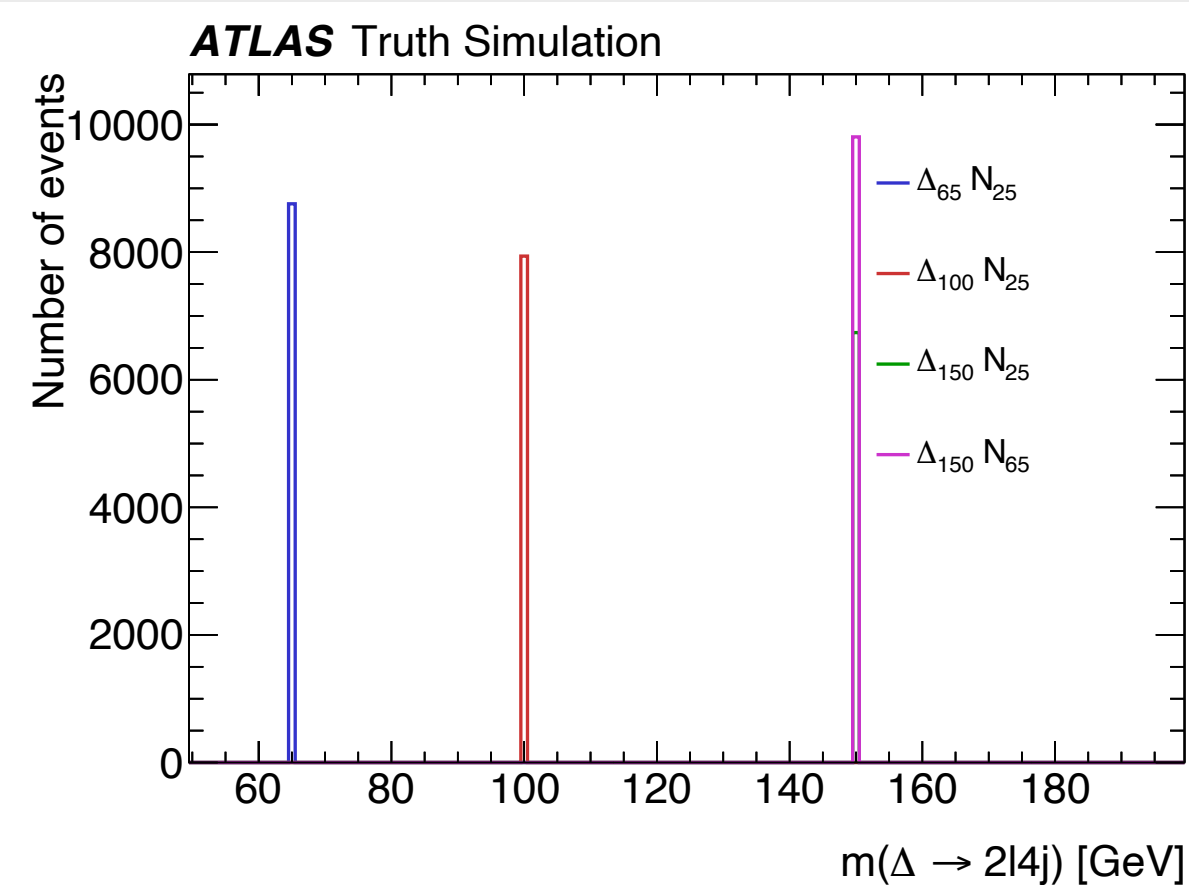
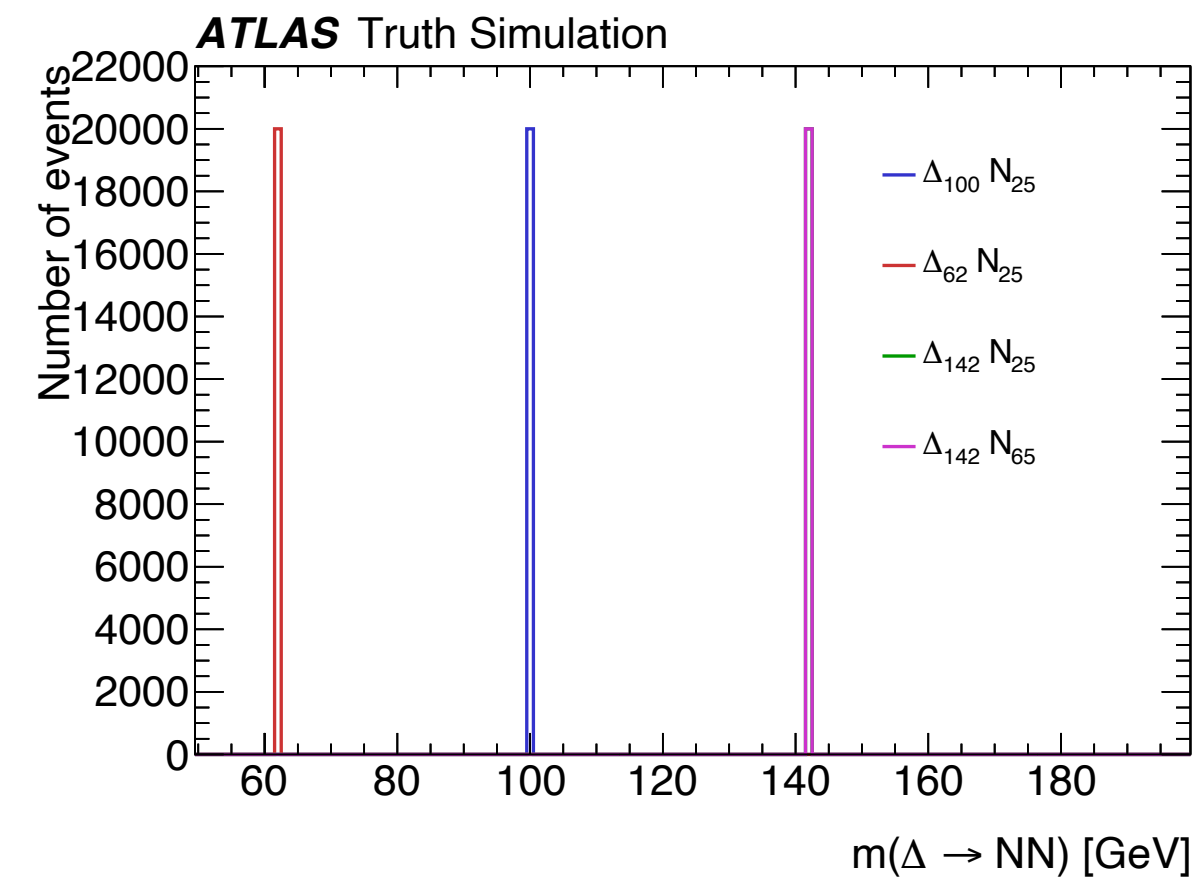
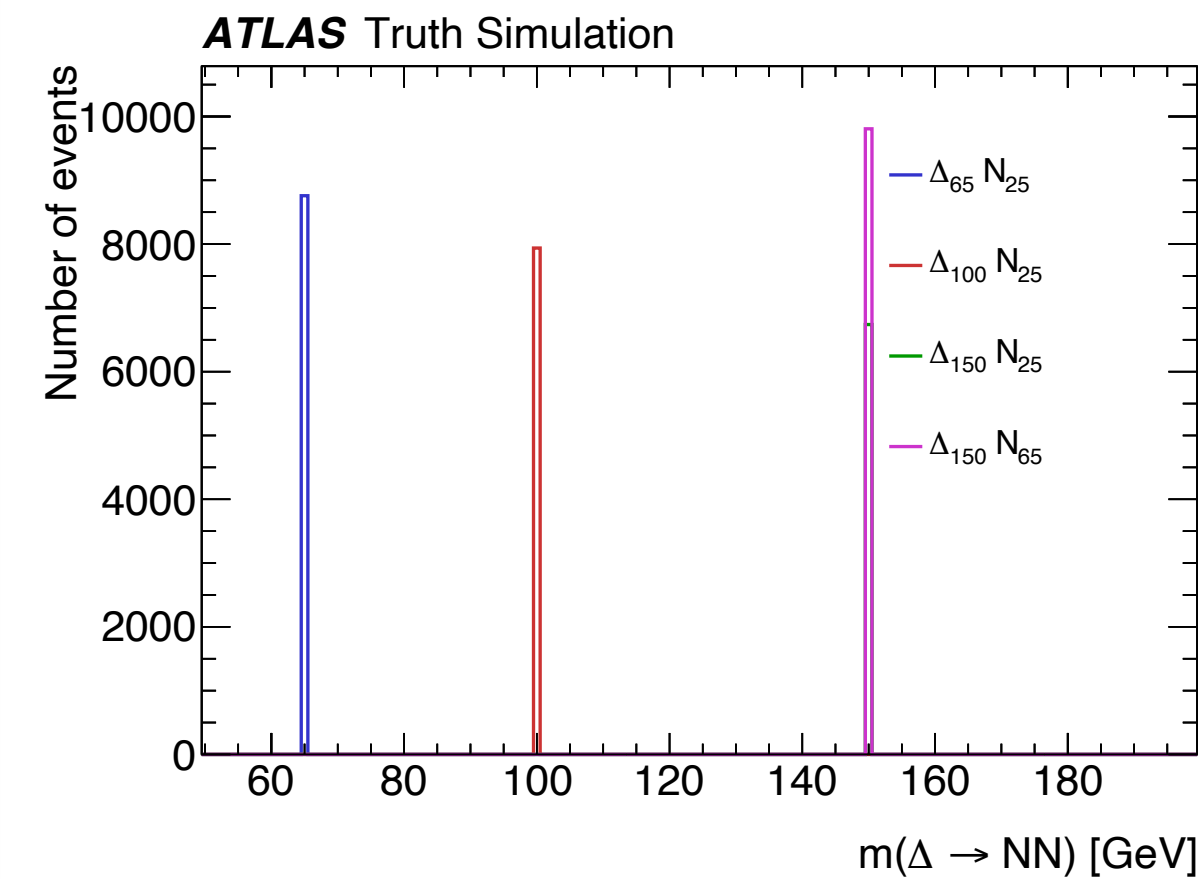
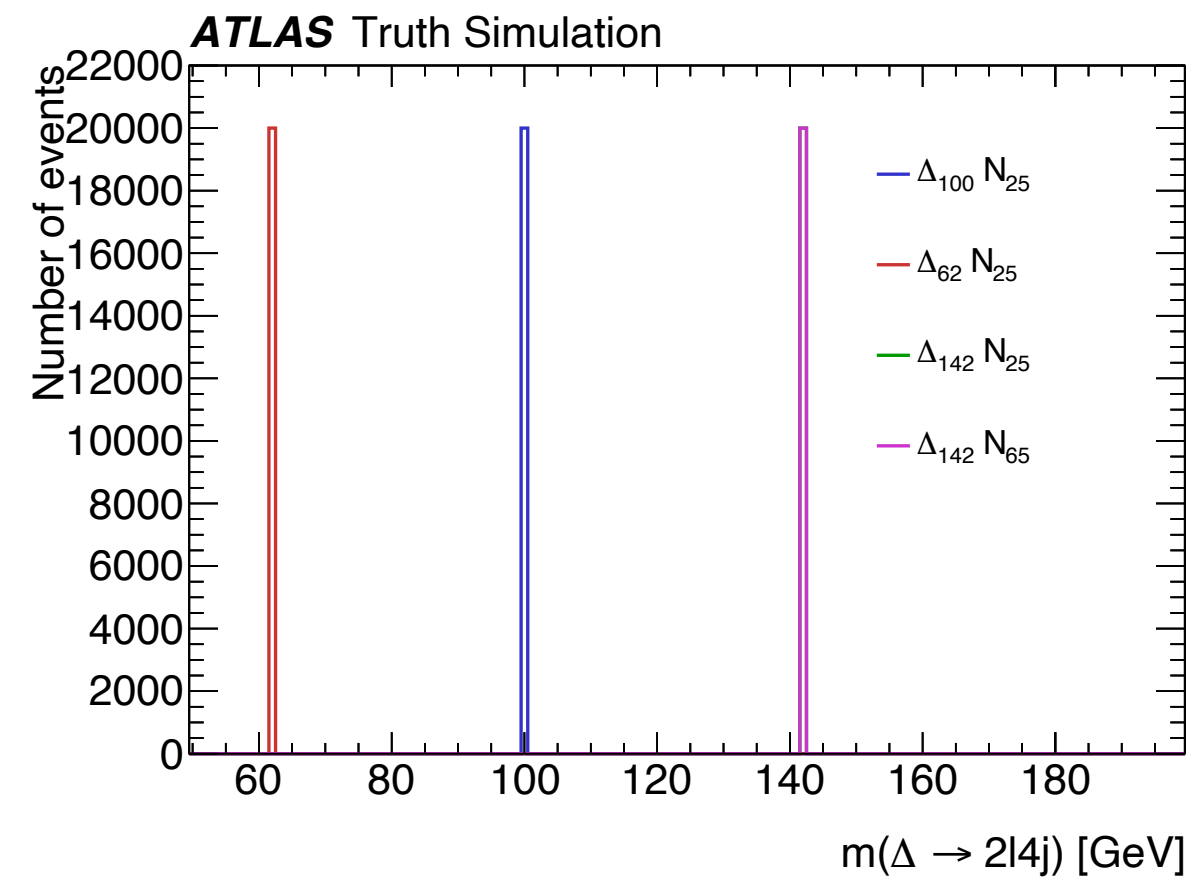
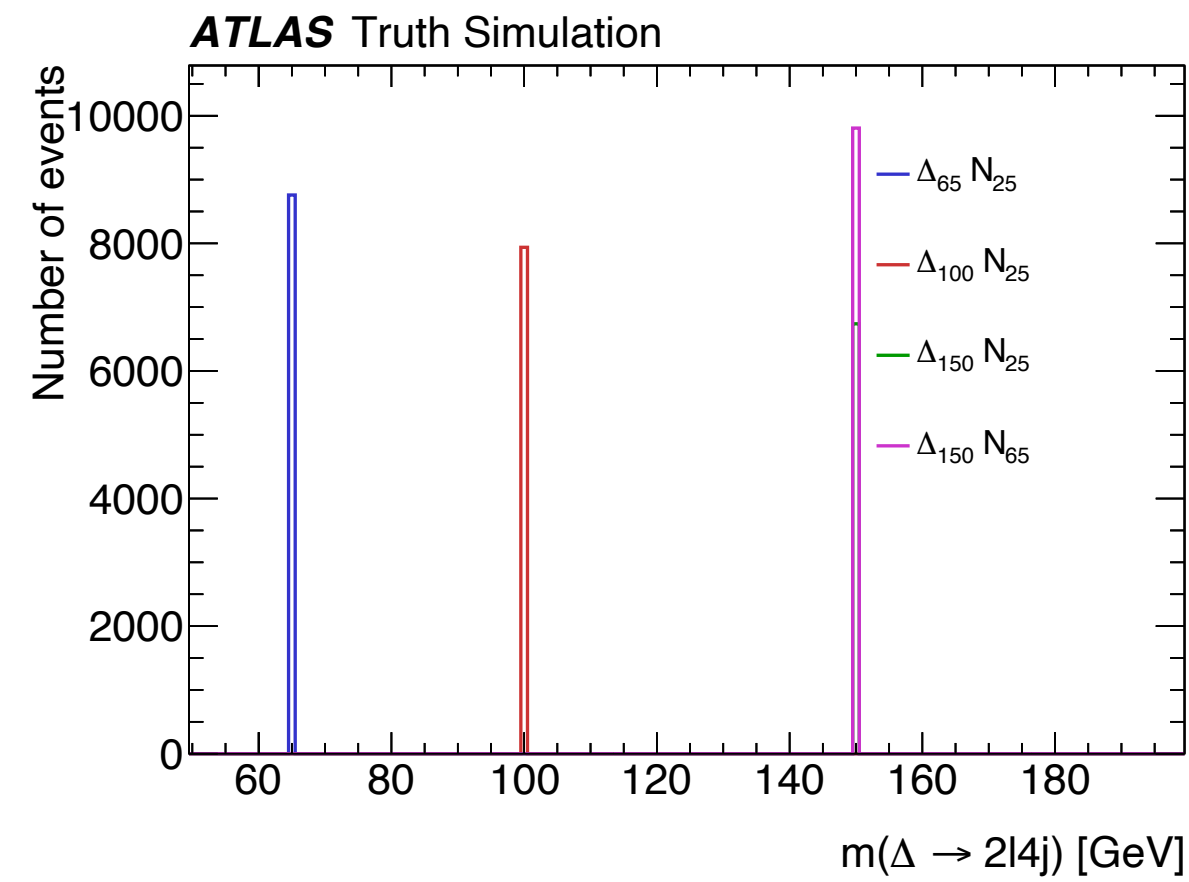
$m(N)$ [GeV]	$m(\Delta)$ [GeV]	$m(W_R)$ [TeV]	cross-section [fb] @13 TeV	cross-section [fb] @13.6 TeV
25	65	4.0	4.60	4.90
25	100	4.0	2.59	2.77
25	150	4.0	1.54	1.67
65	150	4.0	1.06	1.15



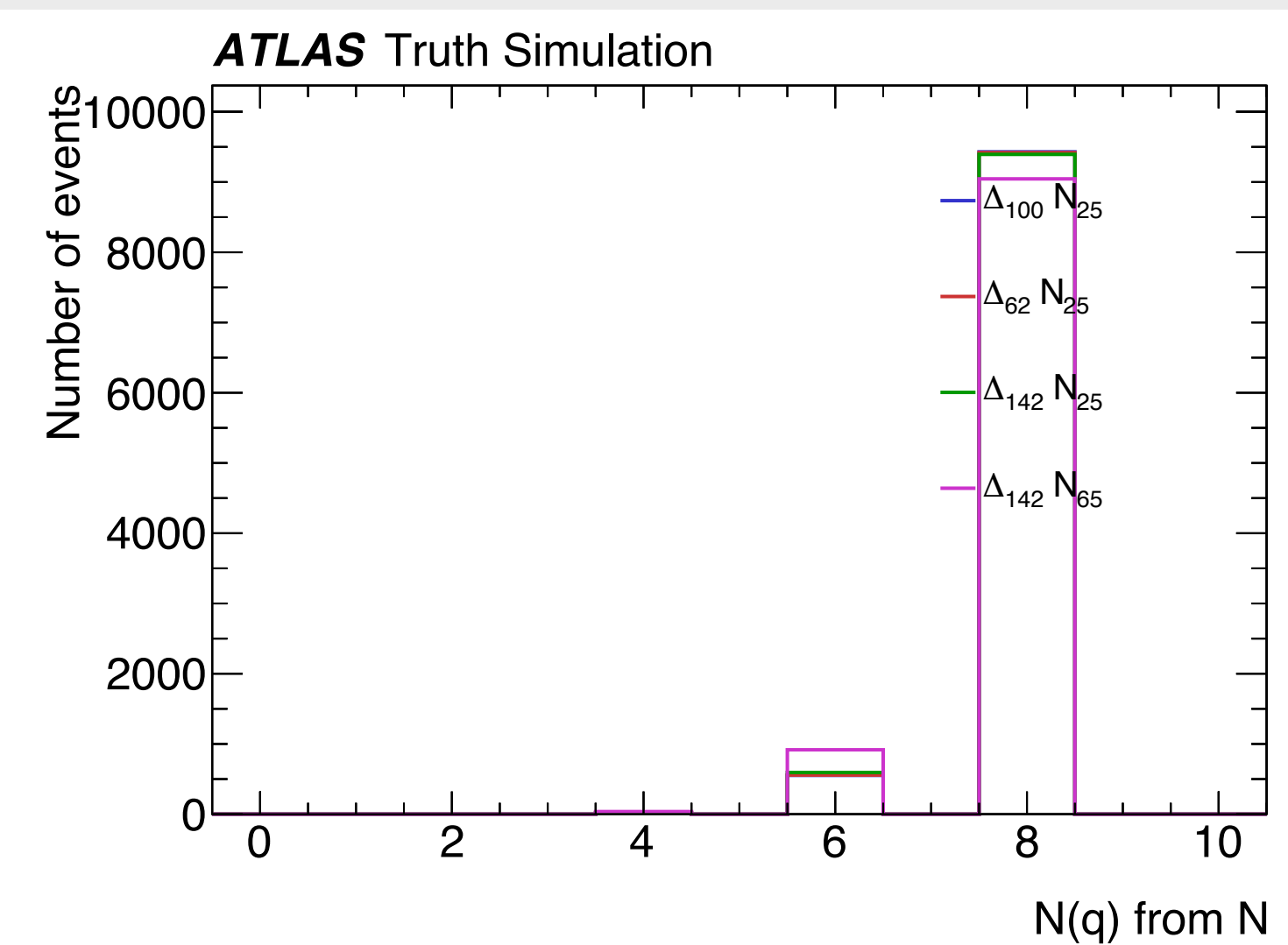
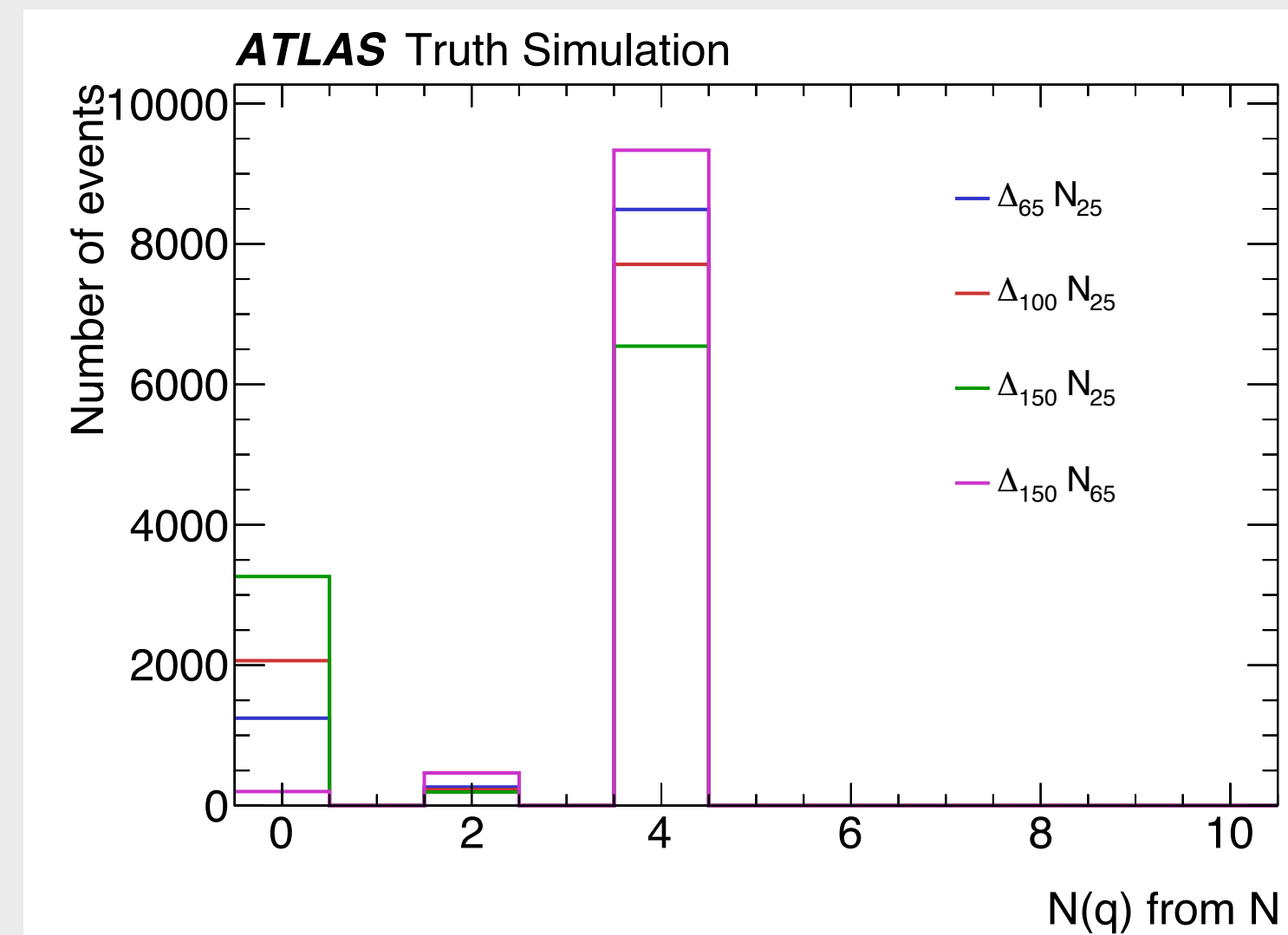
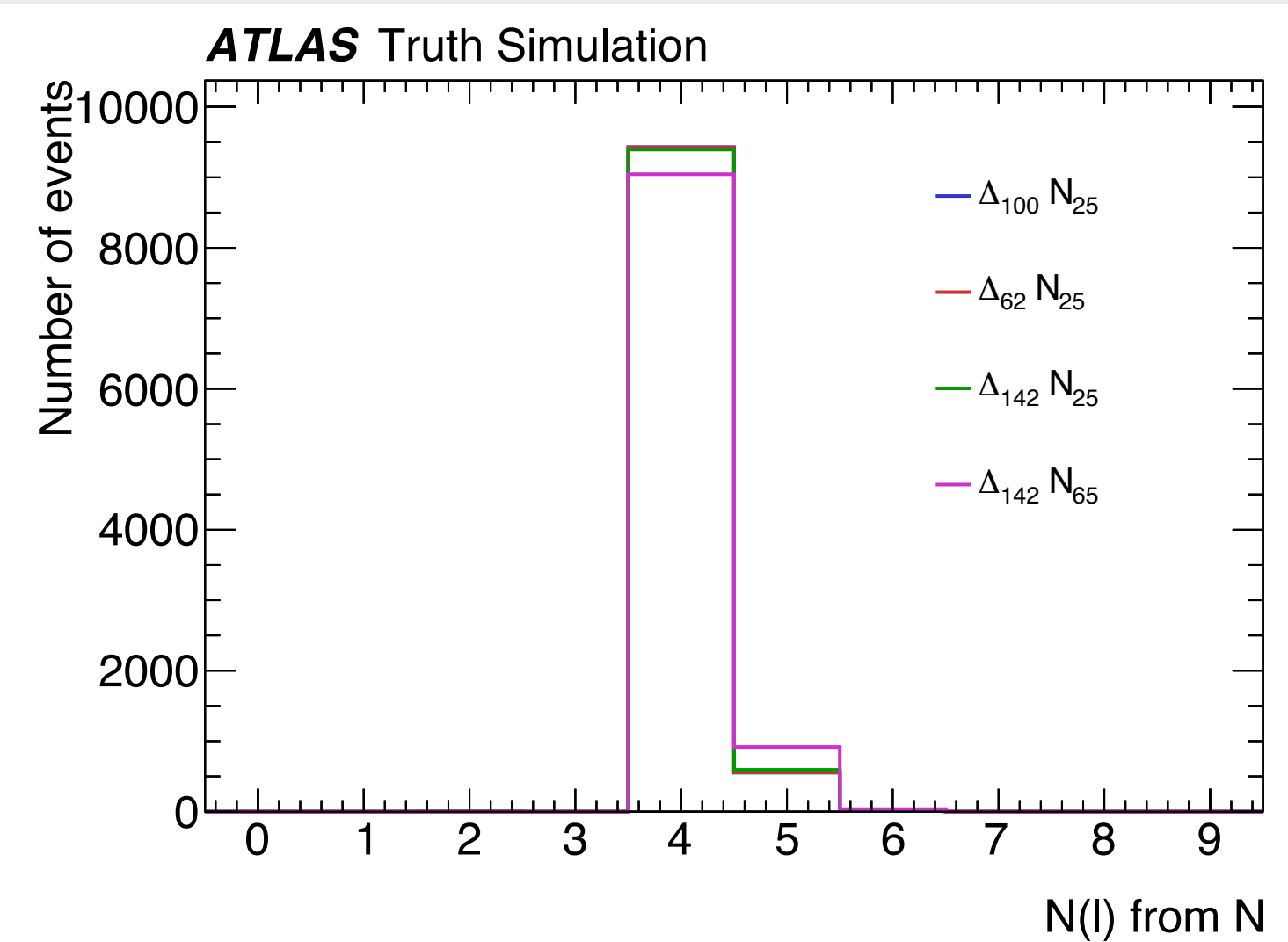
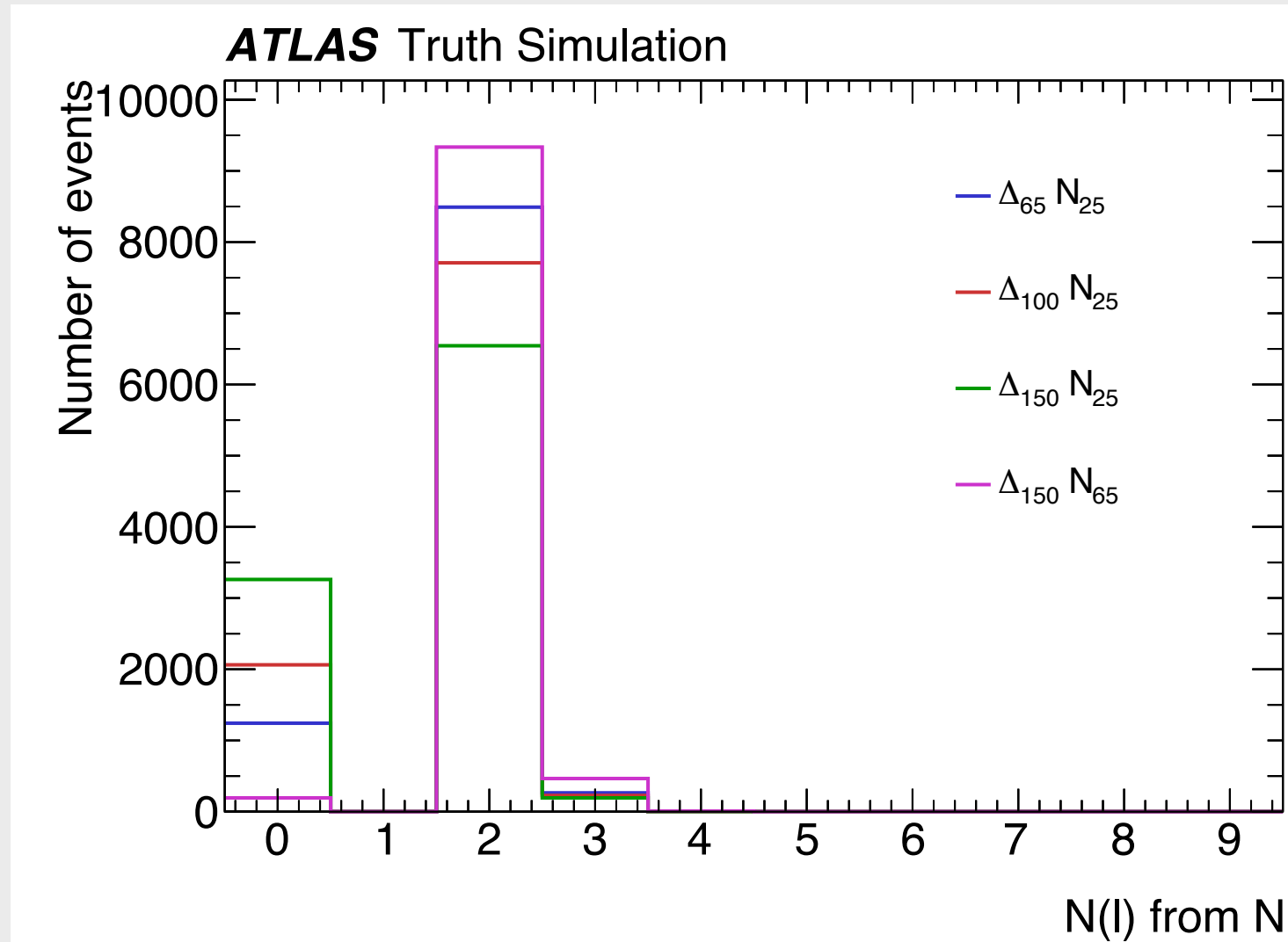
- Number of Δ s and Heavy neutrinos.



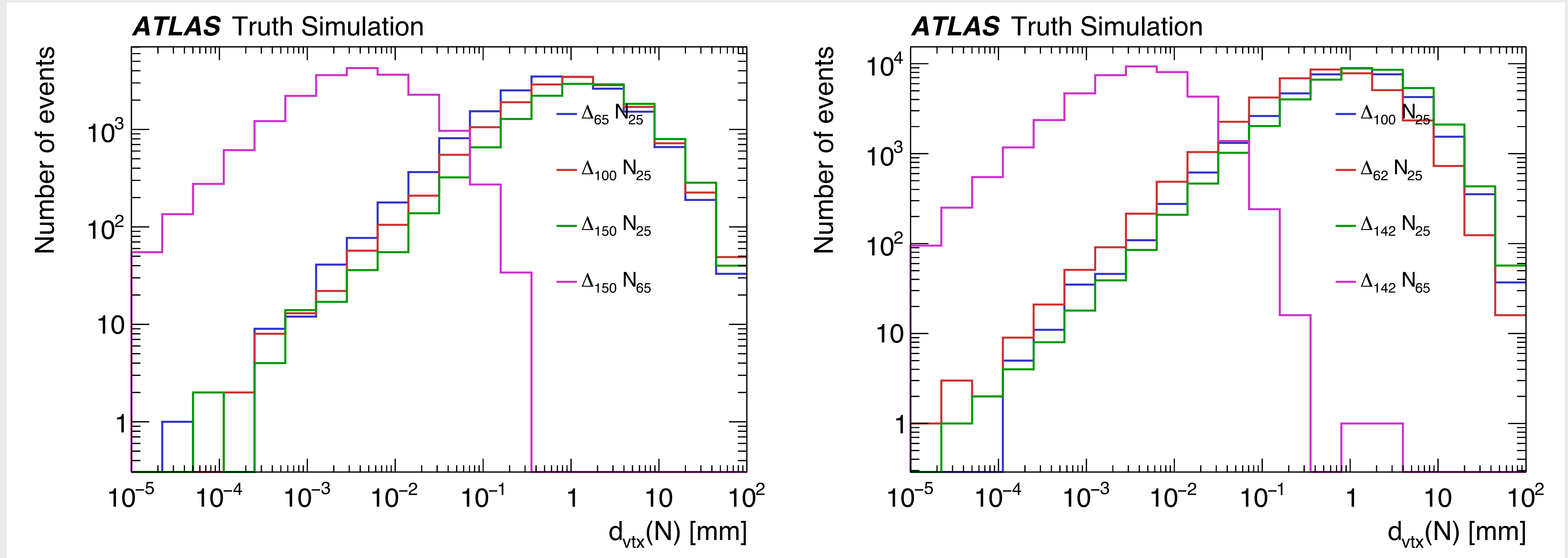
- Reconstructed invariant masses.



- Number of leptons and quarks.



- Heavy neutrino **displacement is detectable** and can play a significant role in our analysis!



$$c\tau = 0.1\text{mm} \left(\frac{40 \text{ GeV}}{m_N} \right)^5 \left(\frac{M_{W_R}}{5 \text{ TeV}} \right)^4$$

- Theoretical models relevant for the DCH analysis were presented.
- Search for a doubly charged Higgs analysis was outlined:
 - Two analysis approaches were tested,
 - Results are interpreted in different models,
 - Background modelling is crucial,
 - Preliminary studies show that MVA techniques provide modest improvements, indicating that the current cut-based analysis is well executed.
- Improvements can be made in the MVA approach, especially with the new person power joining the new round of the analysis.
- Signal sample validation of the LNV Higgs decays was presented.

Thank you...



Istituto Nazionale di Fisica Nucleare
Sezione di Bologna

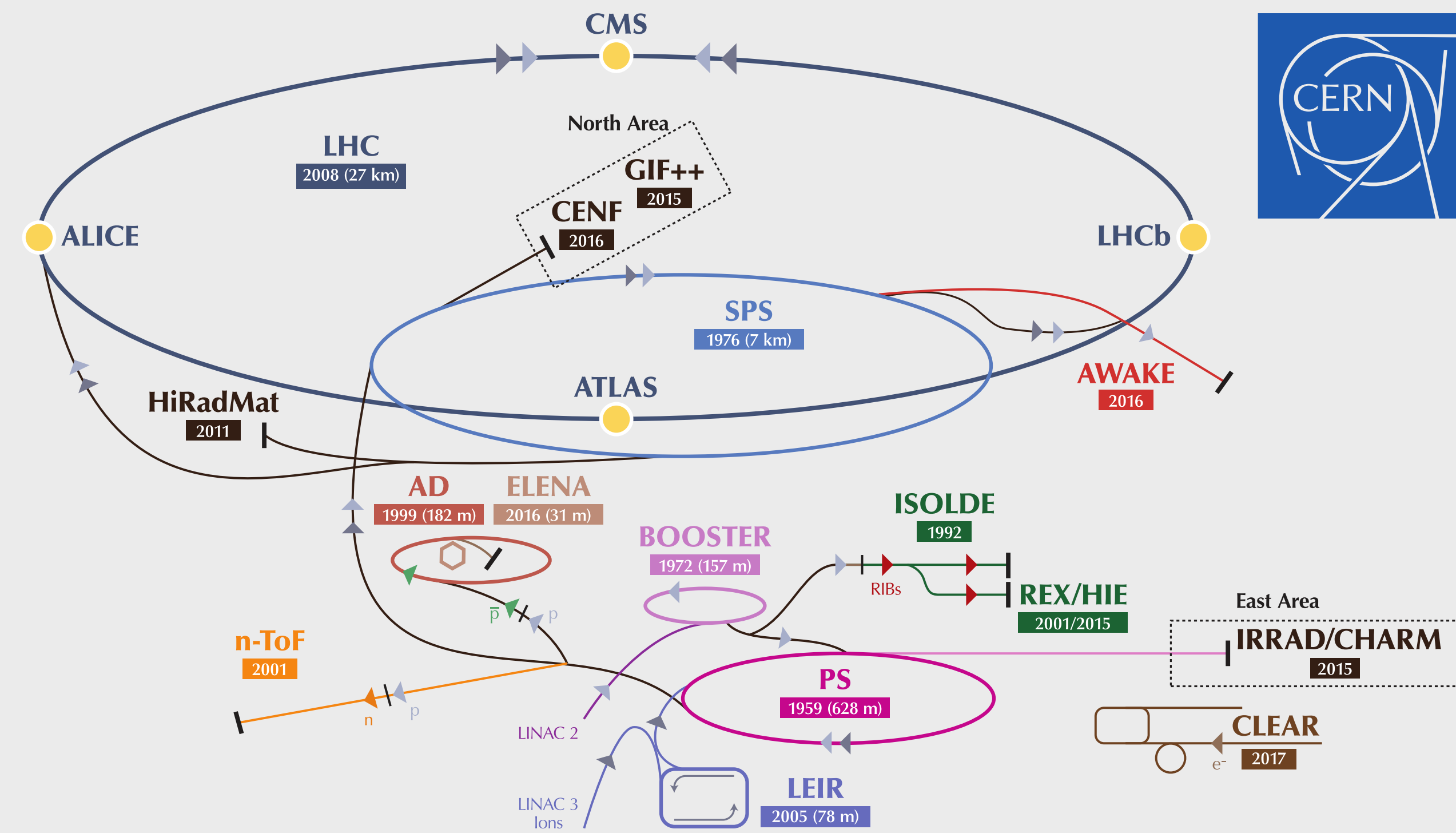




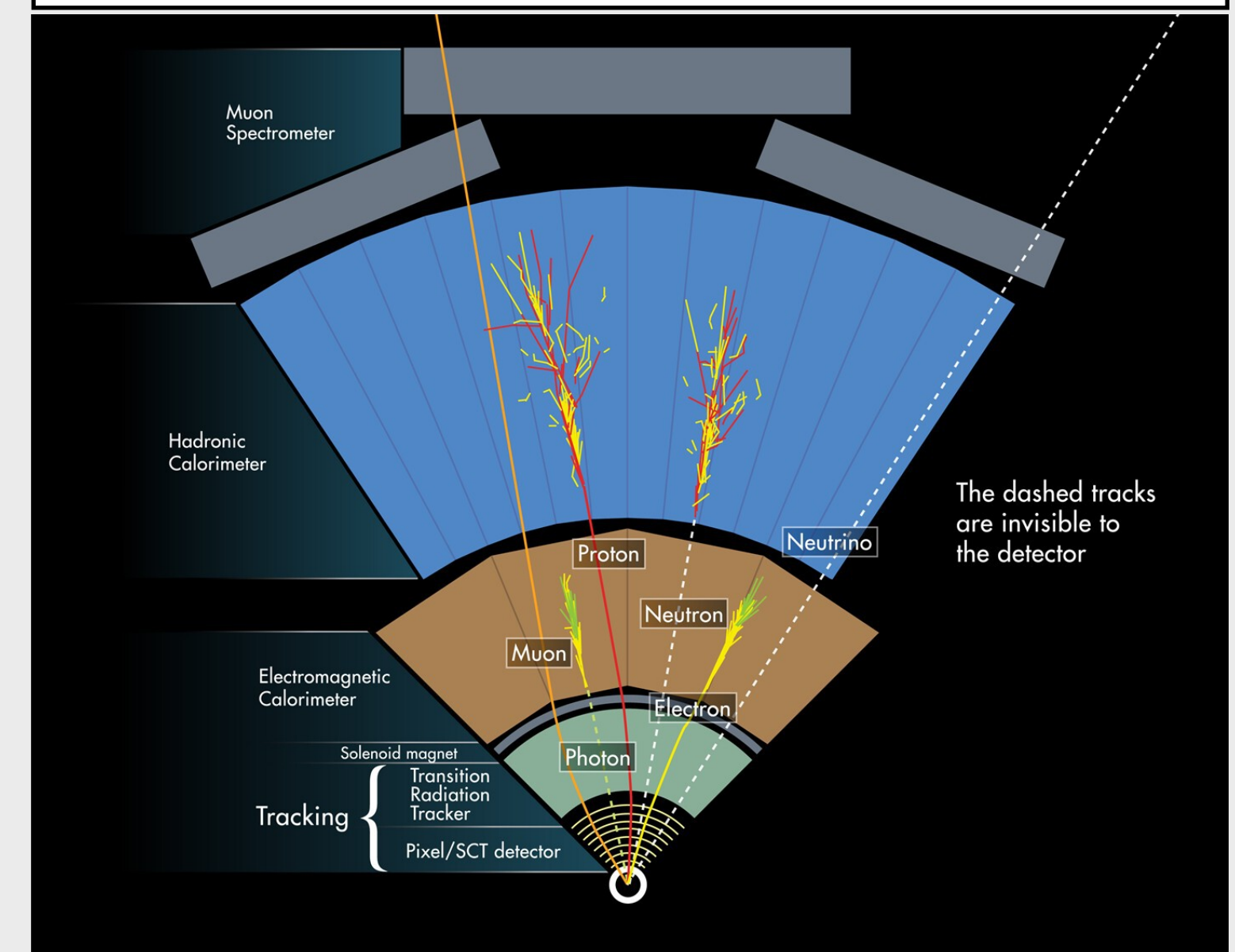
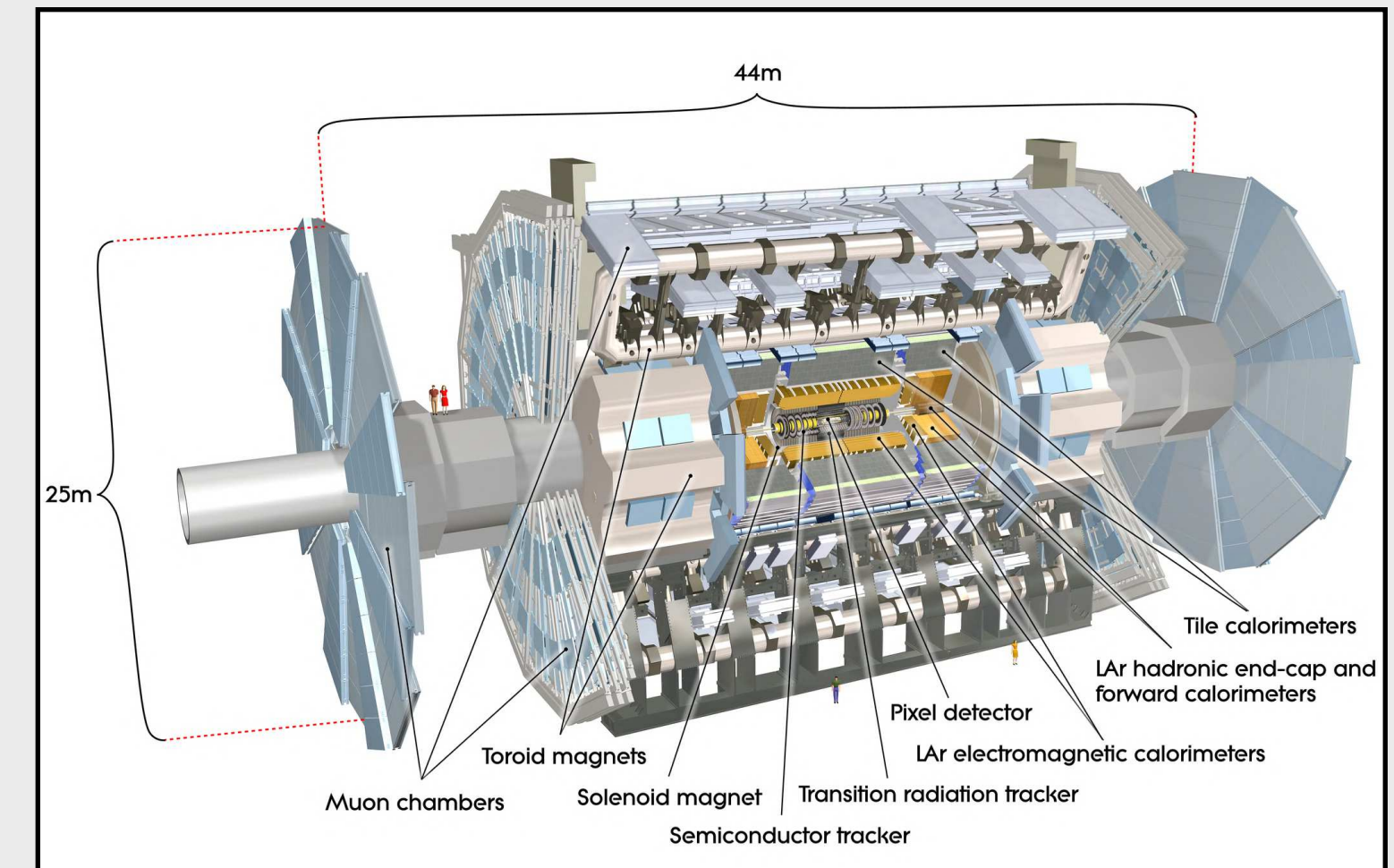
Backup slides

The Large Hadron Collider

- The Large Hadron Collider (LHC) is the most powerful particle accelerator and collider ever built.
- It is located on France - Switzerland border at CERN ~100m below the surface.
- It can accelerate both **protons** and **heavy ions**, which is done in multiple stages.
- It was designed to reach centre-of-mass energy of pp collisions $\sqrt{s} = 14$ TeV.
- Four major experiments:
 - ALICE, **ATLAS**, CMS, LHCb.
- It delivered:
 - $5.46 \text{ fb}^{-1} + 22.8 \text{ fb}^{-1}$ in Run 1,
 - 156 fb^{-1} in Run 2, and
 - 70 fb^{-1} in Run 3to ATLAS.

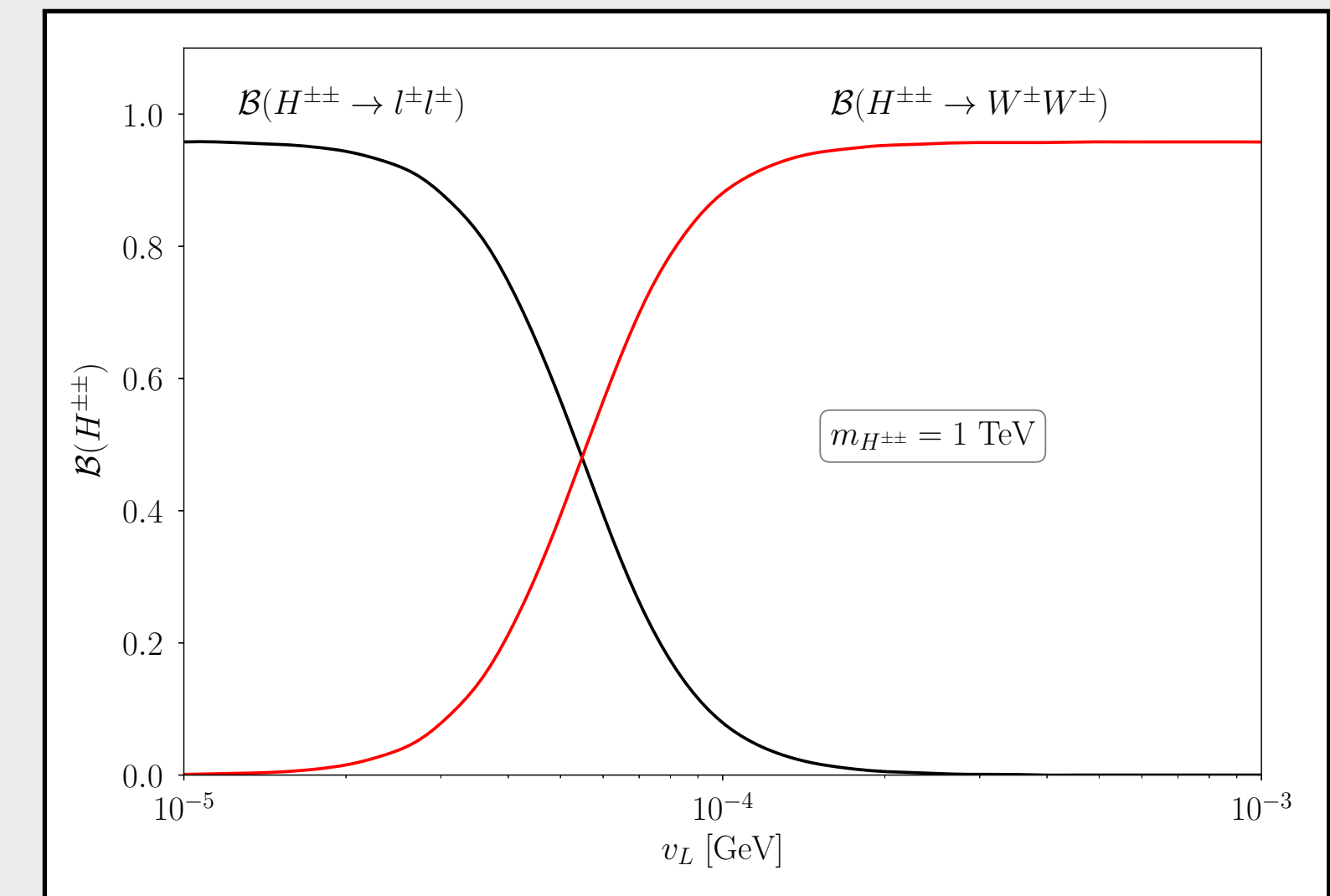
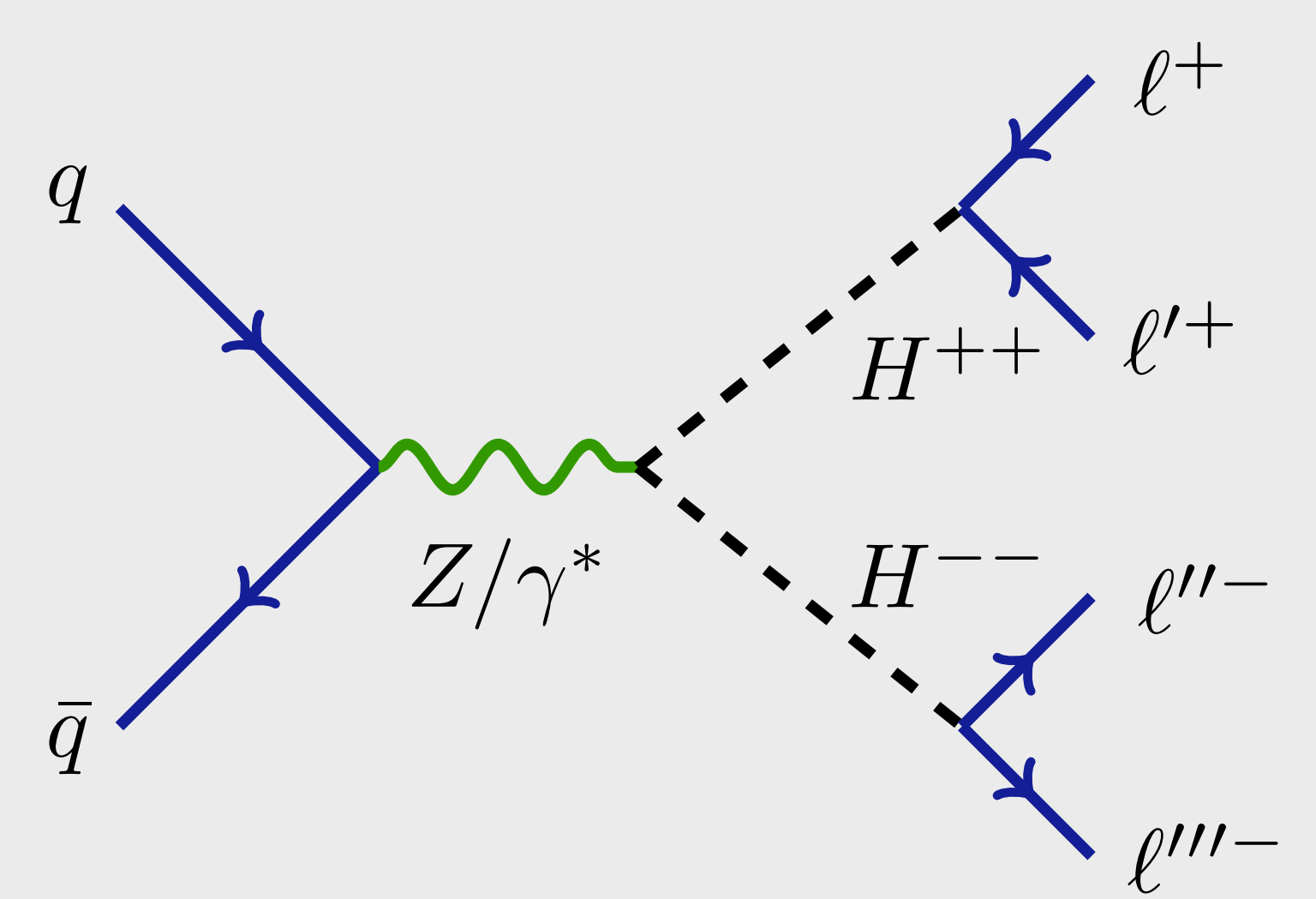


- The ATLAS detector is a multilayered general purpose detector, located at Point 1 of the LHC. It consists of several subsystems:
 - Inner Detector,
 - Electromagnetic and Hadronic calorimeters,
 - Muons spectrometers,
 - Central solenoid and toroid magnets,
 - Forward detectors (luminosity),
 - Two-level trigger.
- It was developed to detect various types of elementary particles with different reconstruction and identification techniques:
 - Electrons, muons, jets, tau leptons, missing transverse energy.



Doubly Charged Higgs Analysis

- Search for $H^{\pm\pm}$ pair production in all lepton flavour and charge combinations: $H^{\pm\pm} \rightarrow \ell^\pm \ell'^{\pm}$, $\ell, \ell' = e, \mu, \tau$.
- Results can be interpreted for:
 - Type-II Seesaw mechanism within Left-Right Symmetric Model: two chiralities $H_L^{\pm\pm}$ and $H_R^{\pm\pm}$ [[arXiv:0305288](#), [arXiv:1105.1379](#)].
 - Zee-Babu neutrino mass model: $k^{\pm\pm}$ [[arXiv:2206.14833](#)]
- Type-II seesaw mechanism: one of the **simplest** known ways to account for the **smallness of the neutrino masses**.
- **Lepton flavour violation is allowed.**
- $v_\Delta \rightarrow 0$ GeV: exclude decays to W bosons.
- Search for $m(H^{\pm\pm}) \in [300 - 1300]$ GeV and focusing on two-, three- and four-lepton final states.
- Using NLO cross-sections with corresponding uncertainties from [Phys. Rev. D 101, 075022](#).



Electrons e

Requirement	Signal electrons (tight)	Background electrons (loose)
Identification	LHTight	LHLoose
Isolation	FCLoose	fail tight selection
p_T cut	$p_T > 40$ GeV	$p_T > 40$ GeV
η cut	$ \eta < 2.47$ and veto $1.37 < \eta < 1.52$	$ \eta < 2.47$ and veto $1.37 < \eta < 1.52$
$ d_0 /\sigma_{d_0}$ cut	$ d_0 /\sigma_{d_0} < 5.0$	$ d_0 /\sigma_{d_0} < 5.0$
$ z_0 \sin(\theta) $ cut	$ z_0 \sin(\theta) < 0.5$ mm	$ z_0 \sin(\theta) < 0.5$ mm
Bad cluster veto	yes	yes

Jets j

Requirement	Analysis jets
Jet type	AntiKt4EMPFLOWJETS
JVT working point	Tight
fJVT working point	-
p_T cut	$p_T > 20$ GeV
η cut	$ \eta < 2.5$
b -tagging	DL1r with FixedCutBEff_77

Muons μ

Requirement	Signal muons (tight)	Background muons (loose)
Quality	HighPt if $p_T > 300$ GeV else Medium	HighPt if $p_T > 300$ GeV else Medium
Bad muon veto	yes	yes
Isolation	FixedCutTightTrackOnly	fail FixedCutTightTrackOnly
p_T cut	$p_T > 40$ GeV	$p_T > 40$ GeV
η cut	$ \eta < 2.5$	$ \eta < 2.5$
$ d_0 /\sigma_{d_0}$ cut	$ d_0 /\sigma_{d_0} < 3.0$	$ d_0 /\sigma_{d_0} < 3.0$
$ z_0 \sin(\theta) $ cut	$ z_0 \sin(\theta) < 0.5$ mm	$ z_0 \sin(\theta) < 0.5$ mm

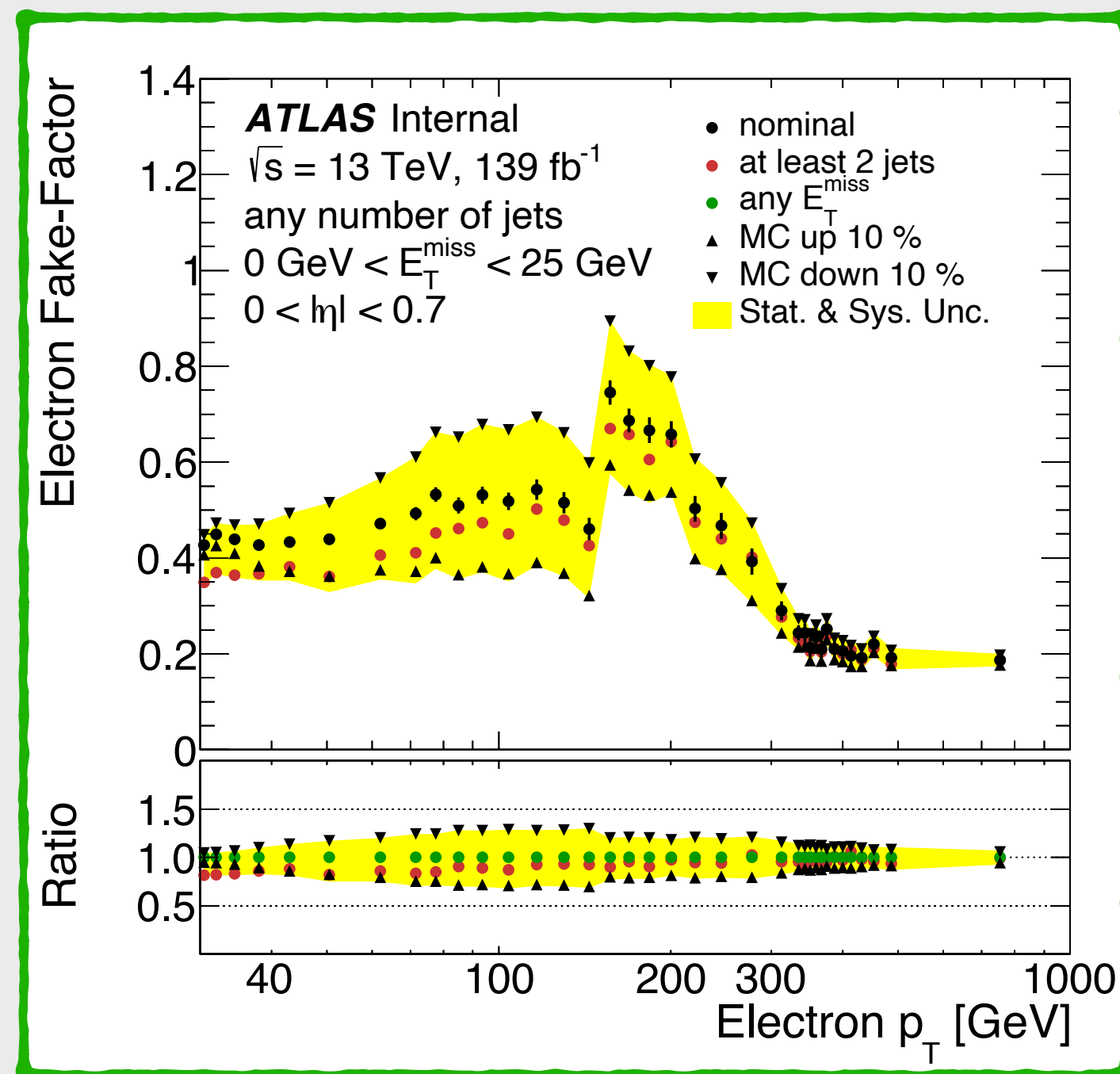
Missing Transverse Energy

Requirement	Nominal
Type	Track-based Soft Term
Working point	Tight
Forward jets	yes
Pile-up jets in significance calculation	no

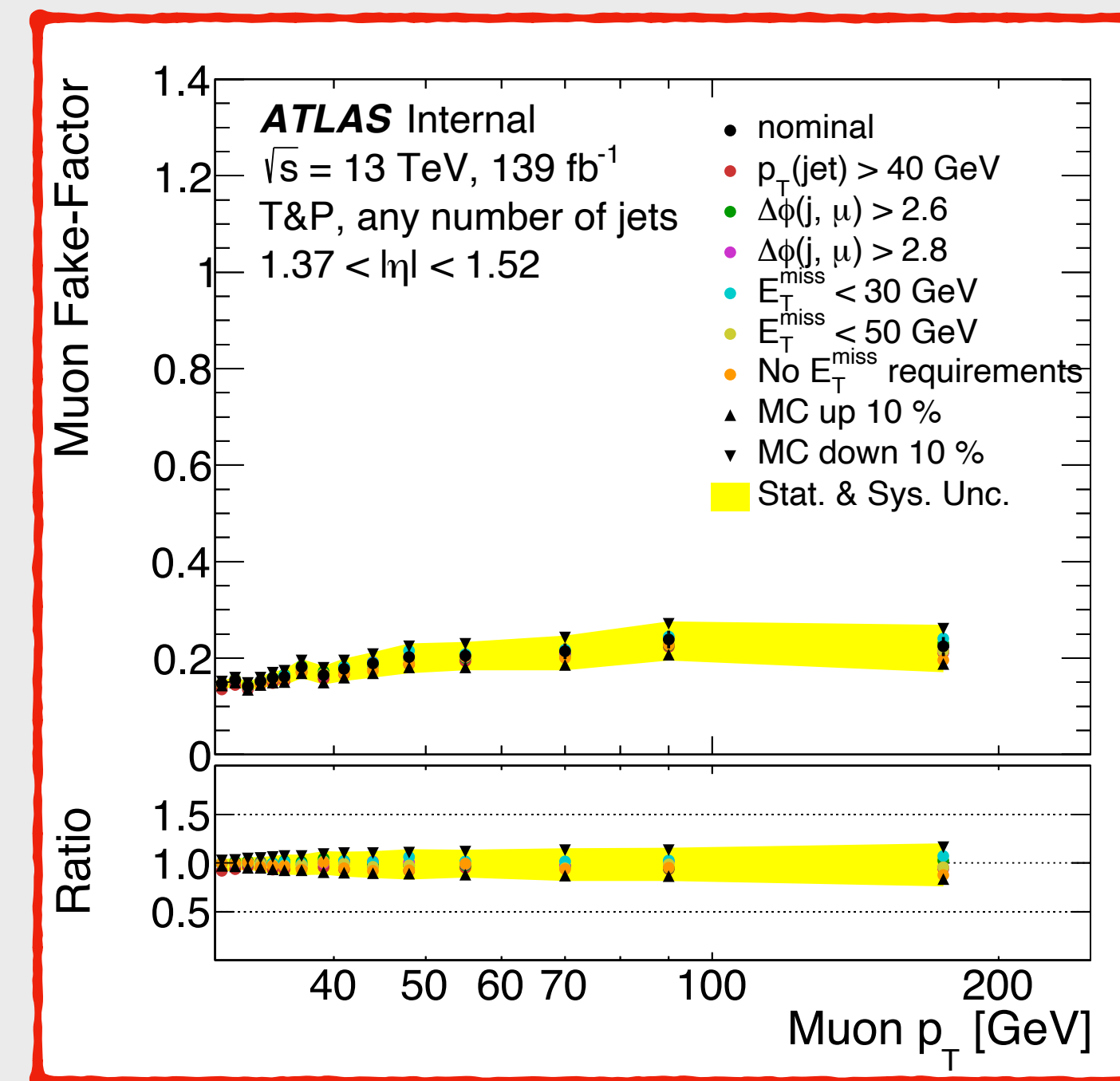
Fake enriched regions

- Fake enriched regions require **exactly 1 baseline lepton** with **any number of jets** and a **b-jet veto** (+ additional kinematic constraints in the muon fake estimation due to tag-and-probe technique).
- Estimation performed in **4 (5) η** and two E_T^{miss} intervals with variable p_T binning.
- Prescaled single lepton triggers were used.

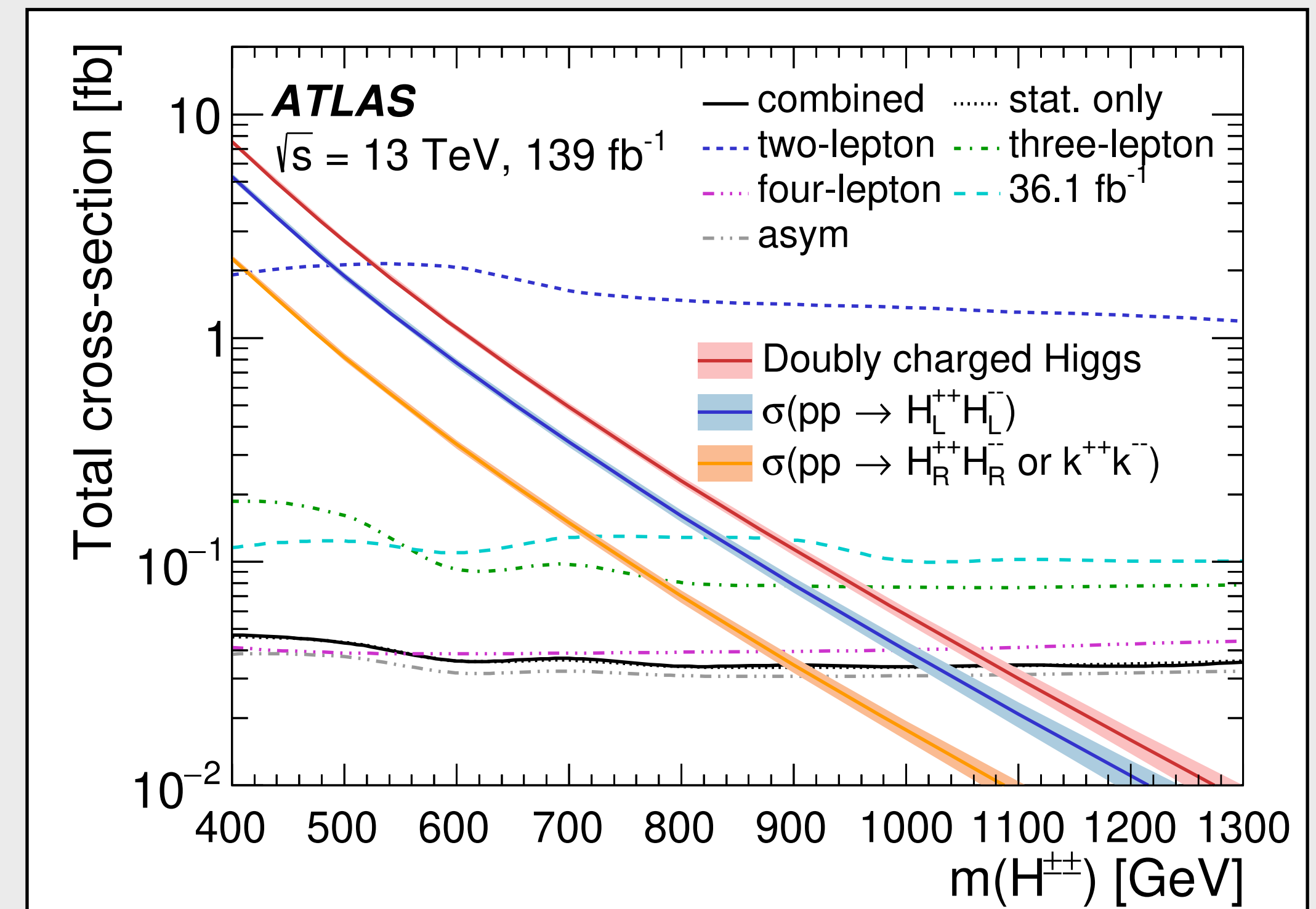
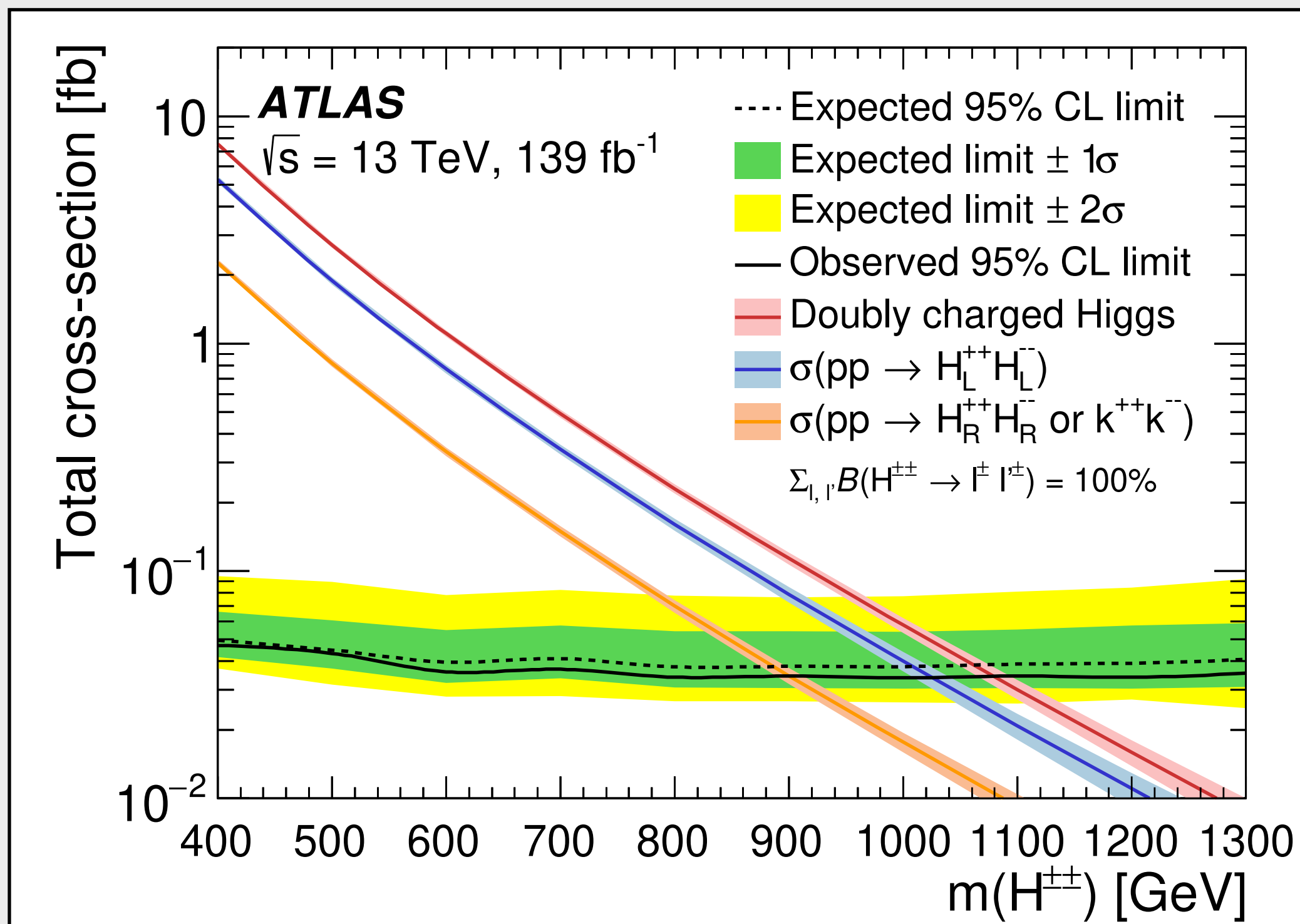
electrons



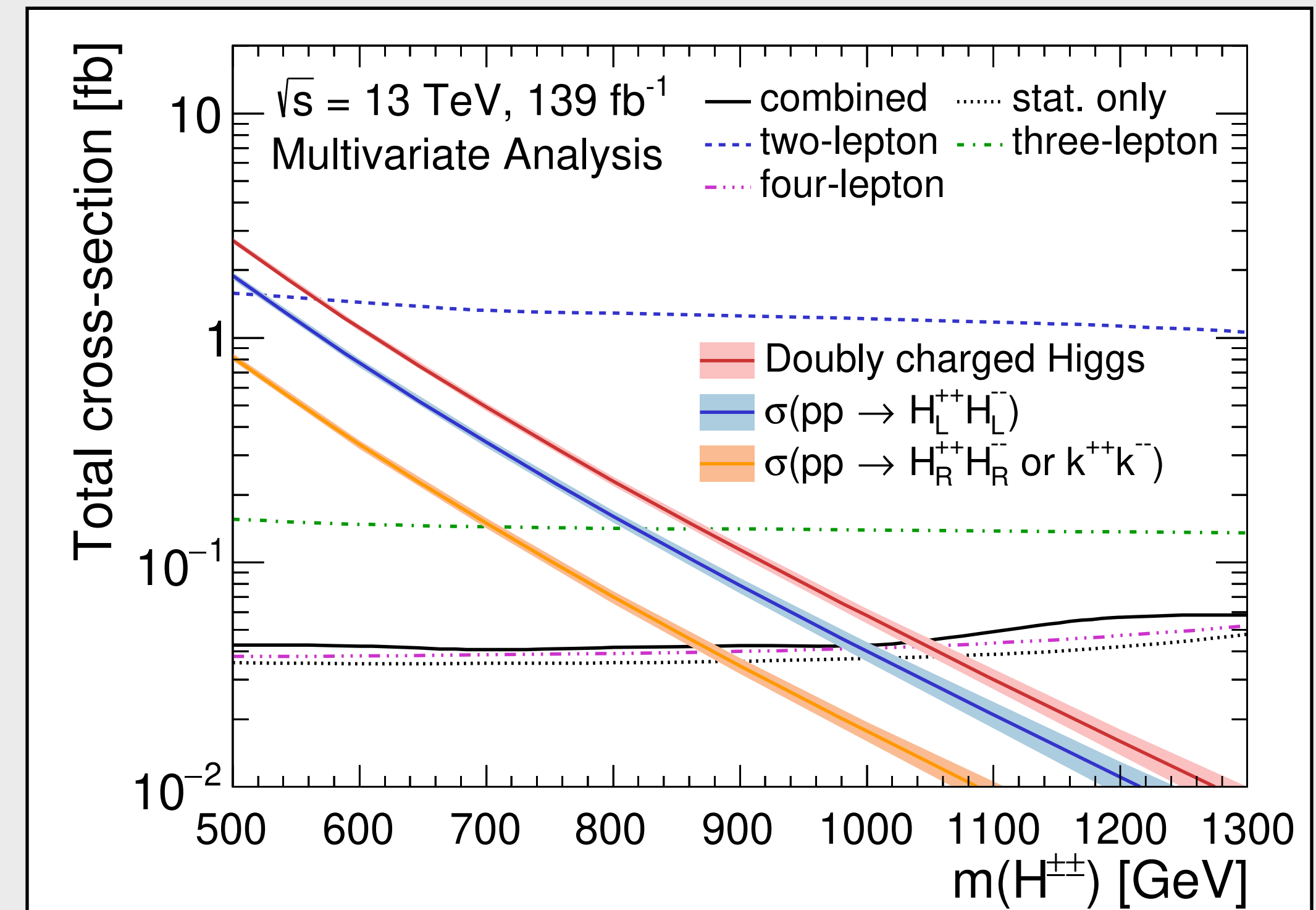
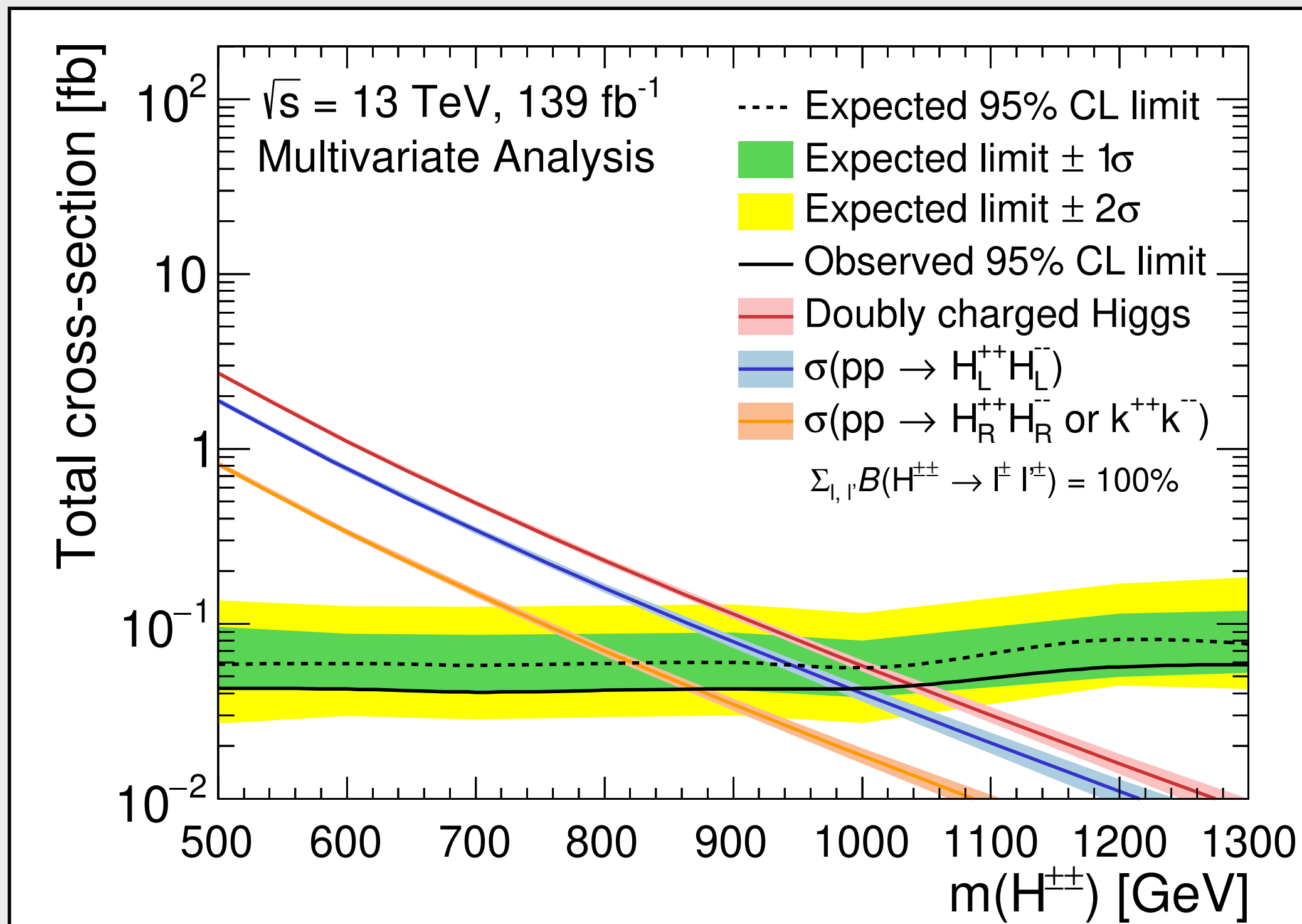
muons



- The expected exclusion limit is 1065^{+30}_{-50} GeV for LRSM and 880^{+30}_{-40} GeV for the Zee-Babu model.
- The observed lower limit on the $H^{\pm\pm}$ mass reaches **1080 GeV** and **900 GeV** when combining all three channels for LRSM and the Zee-Babu model, respectively.
- The four-lepton channel limit is the strongest and drives the combined result.

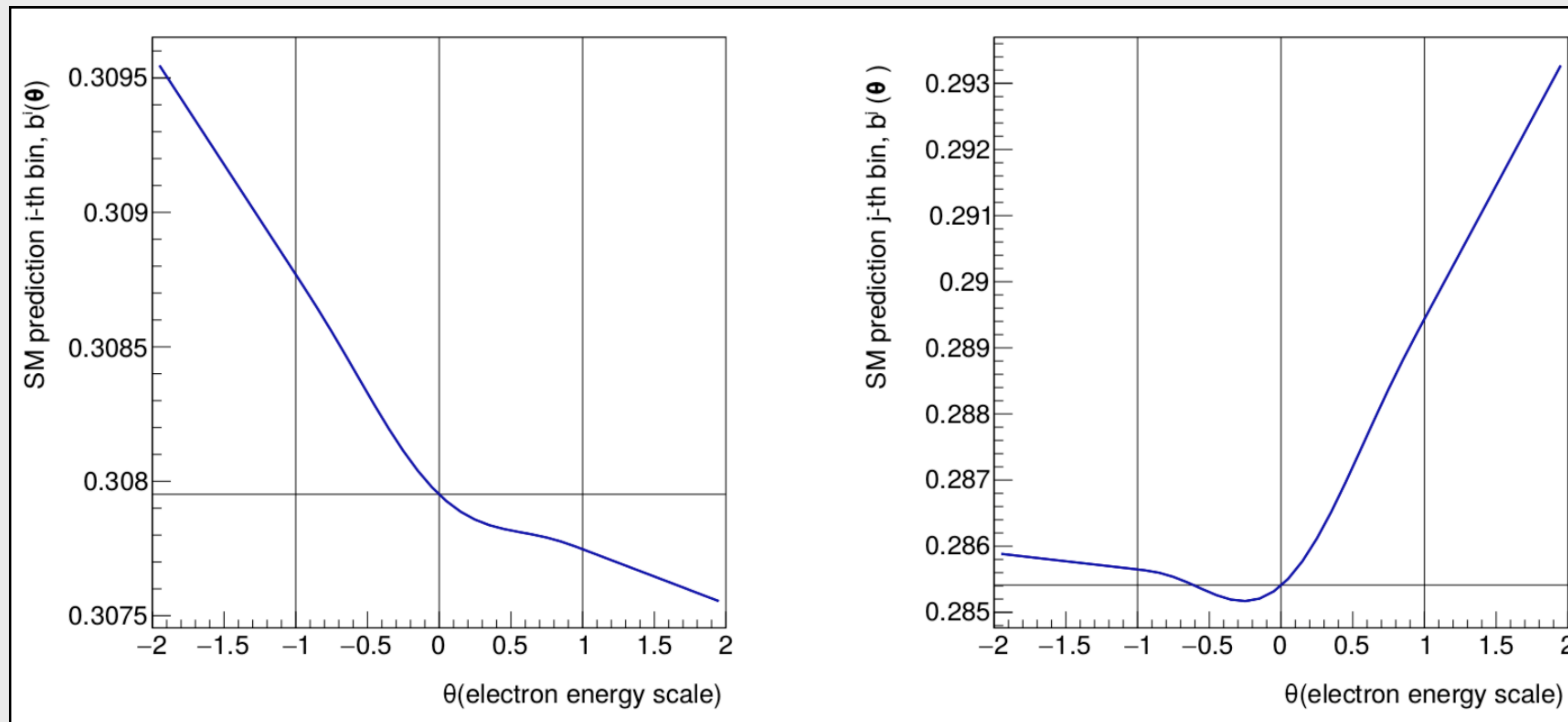


- The expected exclusion limit is 1010^{+80}_{-70} GeV for LRSM and 810^{+20}_{-30} GeV for the Zee-Babu model.
- The observed lower limit on the mass reaches **1050 GeV** and **890 GeV** when combining all three channels for LRSM and the Zee-Babu model, respectively.
- **Not much improvement** with respect to the cut-based analysis.



$$L_{\text{phys}}(\text{data} | \mu) = L_{\text{phys}}(\mu) = \prod_{i \in \text{bins}} \text{Pois}(N_i | \mu S_i + B_i) = \prod_{i \in \text{bins}} \frac{(\mu S_i + B_i)^{N_i}}{N_i!} e^{-(\mu S_i + B_i)}$$

- The predicted SM yield is a function of the nuisance parameter (systematics).



$$\lambda(\mu) = \frac{L(\mu, \hat{\theta}(\mu))}{L(\hat{\mu}, \hat{\theta})}$$

Fixed μ , floating θ

Floating μ , floating θ

$$p_B = P(q_\mu < q_\mu^{ref} | B) = \int_{-\infty}^{q_\mu^{ref}} f(q_\mu | B) dq_\mu$$

$$p_{S+B} = P(q_\mu > q_\mu^{ref} | S + B) = \int_{q_\mu^{ref}}^{\infty} f(q_\mu | S + B) dq_\mu$$

$$CL_s = \frac{p_{S+B}}{1 - p_B}$$

