Mining Heavy Charged Relic

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- For example, in supersymmetry with R-parity, the lightest supersymmetric particle could be a charged slepton.

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- Attracted much experimental interest in 1980s; most notably, search for heavy hydrogen in sea water achieved concentration limit of 10⁻²⁹! [Smith et. al., Nucl. Phys. B 206, 333 (1982)]
- **However**, need to assume their distribution is undisturbed by the evolution of the galaxy, solar system, planetary formation, and water current.

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- We propose to **extract the DM** still embedded in old rocks by:
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- 2. performing a modernized version of the 1980s experiment to reach single-particle sensitivity in macroscopic amount of liquid.
- We also ambitiously aim to extend that experiment's mass sensitivity from 1 TeV to 10¹⁰ TeV.

Contrast with Conventional Paleo-Detectors

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- CHAMPs interact with matter electromagnetically, lose energy by ionization, and stop in rocks.
- In both cases, win by the age of the rock.

Distribution in Geological Rocks

Stopping power of matter due to ionization:

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Atmosphere equivalent to 10 m of rock, so $m_X \leq 10^4$ TeV thermalized in air and do not penetrate earth.

Distribution in the Earth (Flat Earth Approx) Number density of a rock at depth z

$$n_X = \Phi_X \Delta t \frac{z}{2L_{\text{stop}}^2} \Theta(L_{\text{stop}} - z)$$



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$$= \int_{X} \int_{X} \int_{X} \frac{\rho_{\text{DM}}}{m_{X}} v_{\text{DM}}$$

$$f_{X} = f_{X} \left(\frac{\rho_{\text{DM}}}{m_{X}}\right) v_{\text{DM}}$$

$$f_{Y} = \text{fraction of DM}$$







Number density of a rock at depth z





More dark matter at larger z due to the effective area of the surface relative to the flux:

$$\cos\beta = \frac{Z}{L_{\rm stop}}$$

Distribution in the Earth (Round Earth)

$$n_X = \frac{\Phi_X \Delta t}{4R_{\oplus}} \left(\frac{2R_{\oplus}z}{L_{\text{stop}}^2} + 1 \right) \Theta(L_{\text{stop}} - z)$$



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"plateau regime": when $L_{\rm stop} \gtrsim \sqrt{2R_{\oplus}z}$, flat Earth does not apply, and $L_{\rm stop}$ drops out of the equation, $n_X \propto m_X^{-1}$.











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Benchmark rock volume: $V_{rock} = 1 \text{ m}^3$.

Assume we can find a single heavy particle in 1 m³, detection scheme explained later.

10 Myr Rock, 10 m Depth



 \blacksquare X's weight breaks chemical bonds

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Deeper rock not sensitive to lighter particles since L_{stop} is too short.

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Maximally conservative approach: compare 2 scenarios,

- (1) the sample stayed at the spot we found it for 1 Gyr;
- (2) the sample spent equal amount of time at all depth < 100 km.
- We conservatively take the smaller N_X between the two.

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Detection Scheme

Detection Overview

We envision a 3-step search to reach single particle sensitivity:

- 1. Convert sample to liquid (melting, chemical processing).
- 2. Enrich the liquid.
- 3. Mass spectrometry.

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Scale height:
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Using bias random walk (Drude model) with ballistic motion:

$$t_{\rm sink} = 10 \min\left(\frac{10^4 \text{ TeV}}{m_X}\right)^{3/2} \left(\frac{L_{\rm sink}}{1 \text{ m}}\right)$$



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 F_c

Limited by smallest cut on smallest test tube (100 μ L): 1 mm \rightarrow 1 μ L.

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One possibility:

[Gao et. al. 2502.16437] claims a series of SQUID detectors can detect single charged particle with v between 3×10^{-11} and 3×10^{-6} .

Summary and Outlook

Parameter Space



Liquify old rocks, enrichment, mass spec.

Assumed single particle sensitivity.

Complementary to astrophysics [Fedderke et. al., PRD 101, 115021].

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- Rocks (last 10 Myr):
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Complementary Search: Water Pool

- Monitor a large water pool for a year, which collects the X particles thermalized in the air, targeting lighter parameter space.
- 50 simultaneous centrifuges over a year can process 10^5m^3 of water.

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Thank You

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