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Indirect Constraints on Higgs-mediated FCNCs

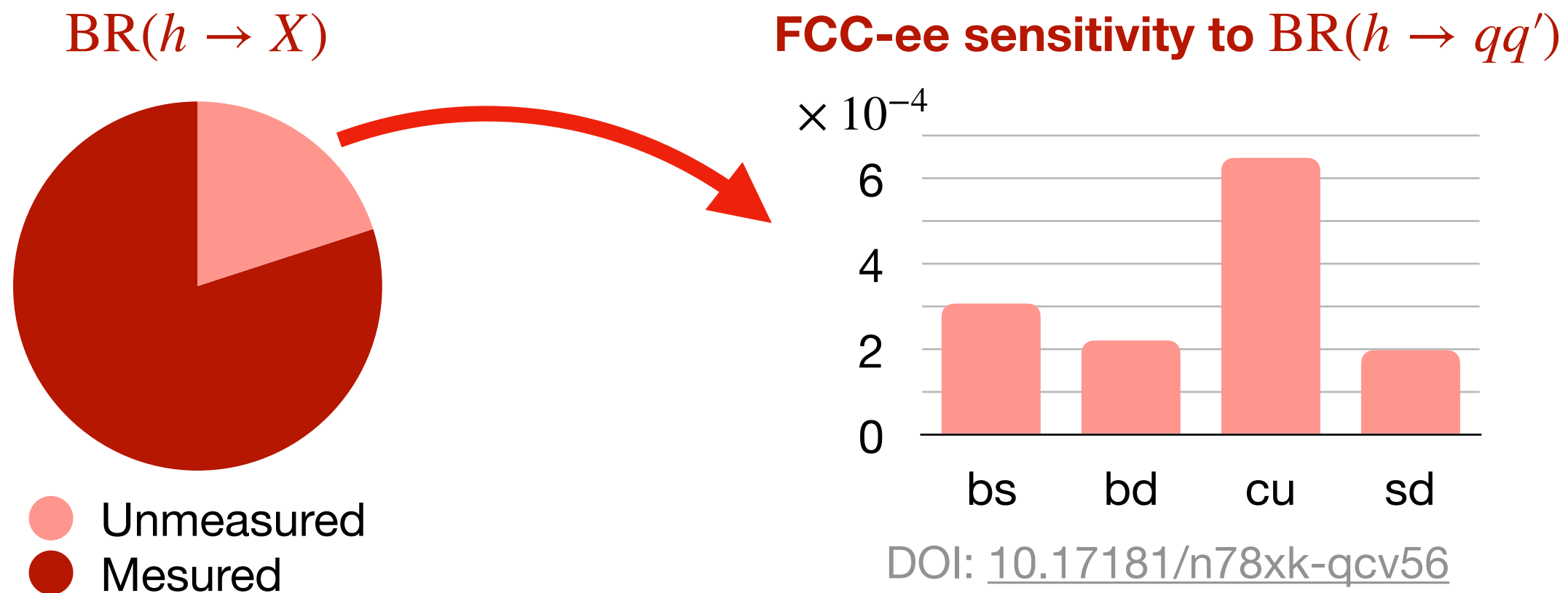
Collaborators: Prof. Dr. Jernej Fesel Kamenik, Dr. Jonathan Kriewald

BRDA 2025, October 30, 2025

Introduction and Motivation

10+ Years After the Higgs Discovery

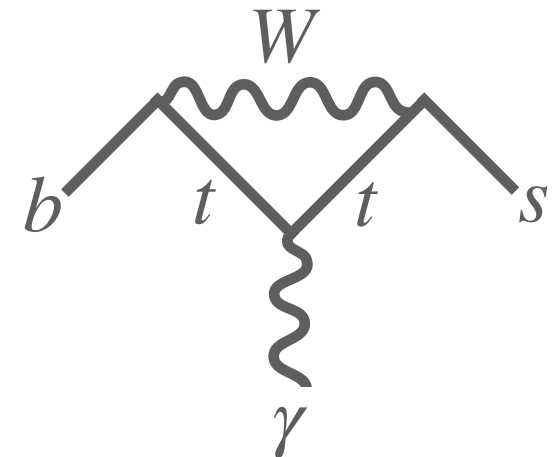
- Higgs (2012) completed SM
- Only a few couplings (t , W , Z , τ) precisely measured
- Light-quark and off-diagonal currently **unmeasured**
- Higgs factories (e.g. FCC-ee) → **sensitivity to** on off-diagonal



Flavor-Changing Neutral Currents

1. FCNCs are forbidden at tree level in SM

- Suppressed due to **GIM mechanism**
- Sensitive to New Physics (NP)
- Via loops $\rightarrow$ virtual heavy particles



2. Intensity vs. energy

- New particles modify effective couplings, even if too heavy to produce
- **Complementary to direct searches** at the energy frontier
- Rare processes $\rightarrow$ indirect search for new particles

Rare FCNC processes act as precision microscopes for heavy new physics. If we observe a Higgs-mediated FCNC, it's BSM.

Setup and Matching

Effective Lagrangian and Matching

- Effective Lagrangian:

$$\mathcal{L}_{eff} \supset y_{sb}(\bar{s}_L b_R)h + y_{bs}(\bar{b}_L s_R)h + y_{db}(\bar{d}_L b_R)h + y_{bd}(\bar{b}_L d_R)h \\ + y_{cu}(\bar{c}_L u_R)h + y_{uc}(\bar{u}_L c_R)h + y_{sd}(\bar{s}_L d_R)h + y_{ds}(\bar{d}_L s_R)h + \text{h.c.}$$

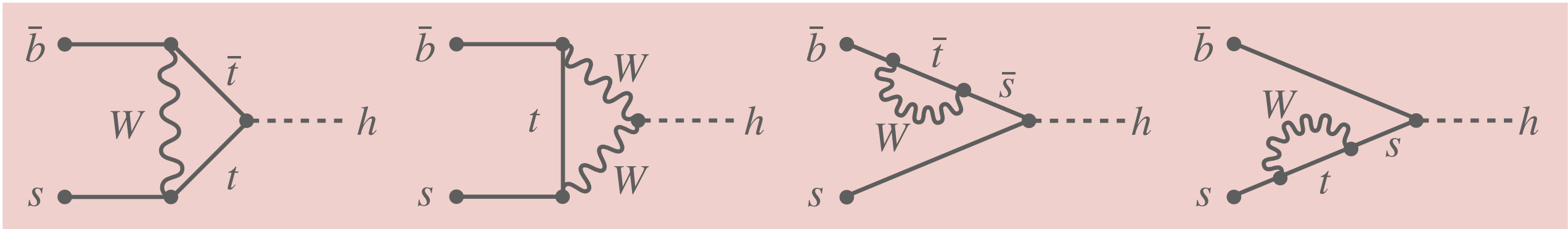
- Relevant operators and matching (example for B_s):

$$\begin{aligned} \mathcal{O}_S^{bs\ell\ell} &= \frac{\alpha_e}{(4\pi)} \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* m_b (\bar{s}_L b_R) (\bar{\ell}\ell) \longrightarrow C_S^{bs\ell\ell} = -\frac{(4\pi)}{\alpha_e} \frac{\sqrt{2}}{4G_F} \frac{1}{V_{tb} V_{ts}^* m_b} \frac{m_\ell}{m_h^2 m_W} \frac{g}{2} y_{sb} \\ \mathcal{O}_S^{'bs\ell\ell} &= \frac{\alpha_e}{(4\pi)} \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* m_b (\bar{s}_R b_L) (\bar{\ell}\ell) \longrightarrow C_S^{'bs\ell\ell} = -\frac{(4\pi)}{\alpha_e} \frac{\sqrt{2}}{4G_F} \frac{1}{V_{tb} V_{ts}^* m_b} \frac{m_\ell}{m_h^2 m_W} \frac{g}{2} y_{bs}^* \\ \mathcal{O}_{S,LL}^{bsbs} &= (\bar{s}_R b_L) (\bar{s}_R b_L) \longrightarrow C_{S,LL}^{bsbs} = \frac{1}{2m_h^2} y_{bs}^{*2} \\ \mathcal{O}_{S,RR}^{bsbs} &= (\bar{s}_L b_R) (\bar{s}_L b_R) \longrightarrow C_{S,RR}^{bsbs} = \frac{1}{2m_h^2} y_{sb}^2 \\ \mathcal{O}_{S,LR}^{bsbs} &= (\bar{s}_R b_L) (\bar{s}_L b_R) \longrightarrow C_{S,LR}^{bsbs} = \frac{1}{m_h^2} y_{bs}^* y_{sb} \end{aligned}$$

Standard Model Prediction

Standard Model Prediction

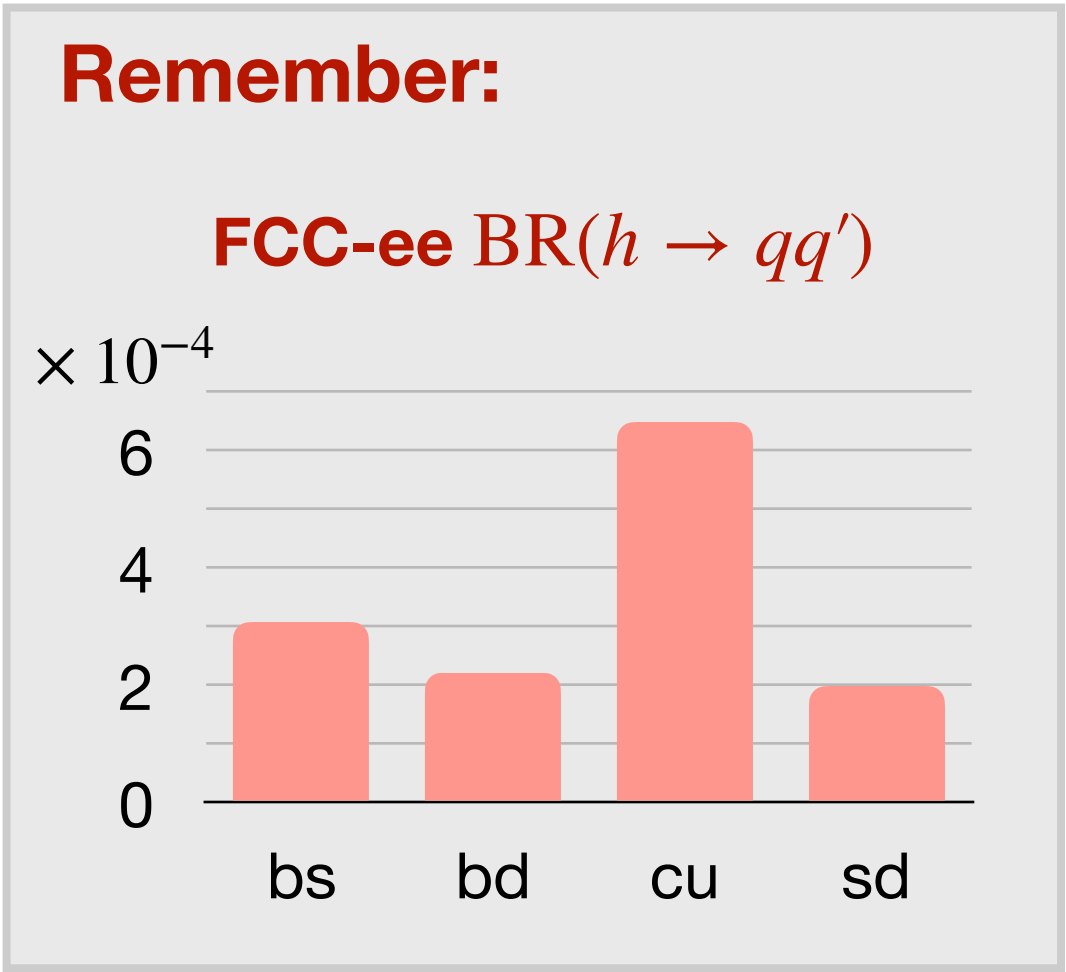
- Unitary gauge (example for $\bar{b}s$):



- Results:

y_{bs}	$(-1.45 - 0.03 i) \times 10^{-7}$
y_{sb}	$(-8.77 + 0.16 i) \times 10^{-6}$
y_{bd}	$(1.43 - 0.57 i) \times 10^{-9}$
y_{db}	$(1.71 + 0.68 i) \times 10^{-6}$
y_{uc}	$(4.79 - 3.96 i) \times 10^{-12}$
y_{cu}	$(1.16 + 0.96 i) \times 10^{-14}$
y_{sd}	$(-6.08 + 2.45 i) \times 10^{-11}$
y_{ds}	$(-1.60 - 0.65 i) \times 10^{-9}$

$BR(h \rightarrow qq')$
4.64×10^{-8}
2.04×10^{-9}
2.34×10^{-20}
1.81×10^{-15}



Current Indirect Bounds

Computation of the Bounds

- Running of Wilson coefficients done with **wilson**
- Low-energy flavor observables computed with **flavio**
- Considered observables:

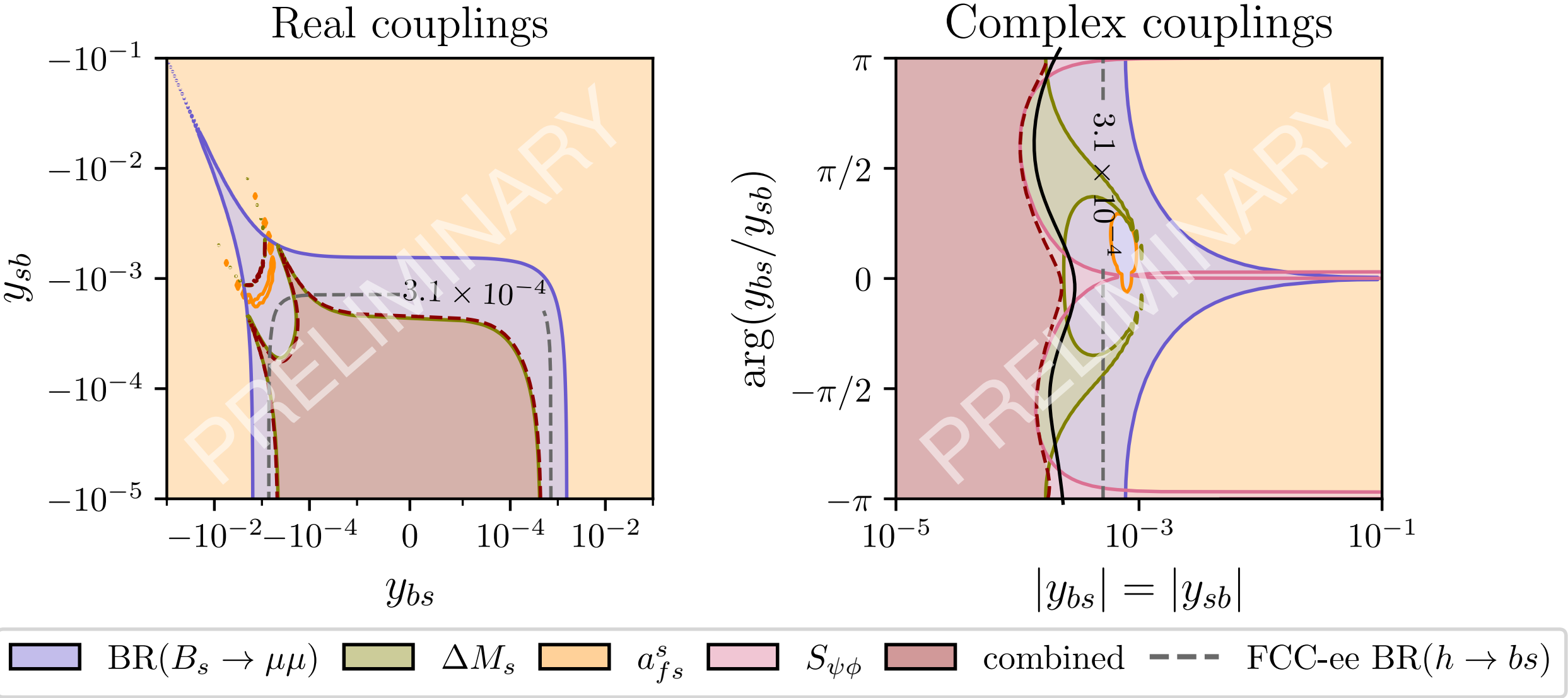
Decay

$$\text{BR}(B_s(B_d, D^0, K^0) \rightarrow \mu^+ \mu^-)$$

Mixing

- $B_{s(d)} - \bar{B}_{s(d)}$: $\Delta M_{s(d)}, \quad a_{fs}^{s(d)}, \quad S_{\psi\phi(K)}$
- $D^0 - \bar{D}^0$: $x_D, \quad \left| \frac{q}{p} \right|_D, \quad \phi_D$
- $K^0 - \bar{K}^0$: $|\varepsilon_K|, \quad \Delta M_K$

Bounds on y_{bs} and y_{sb}

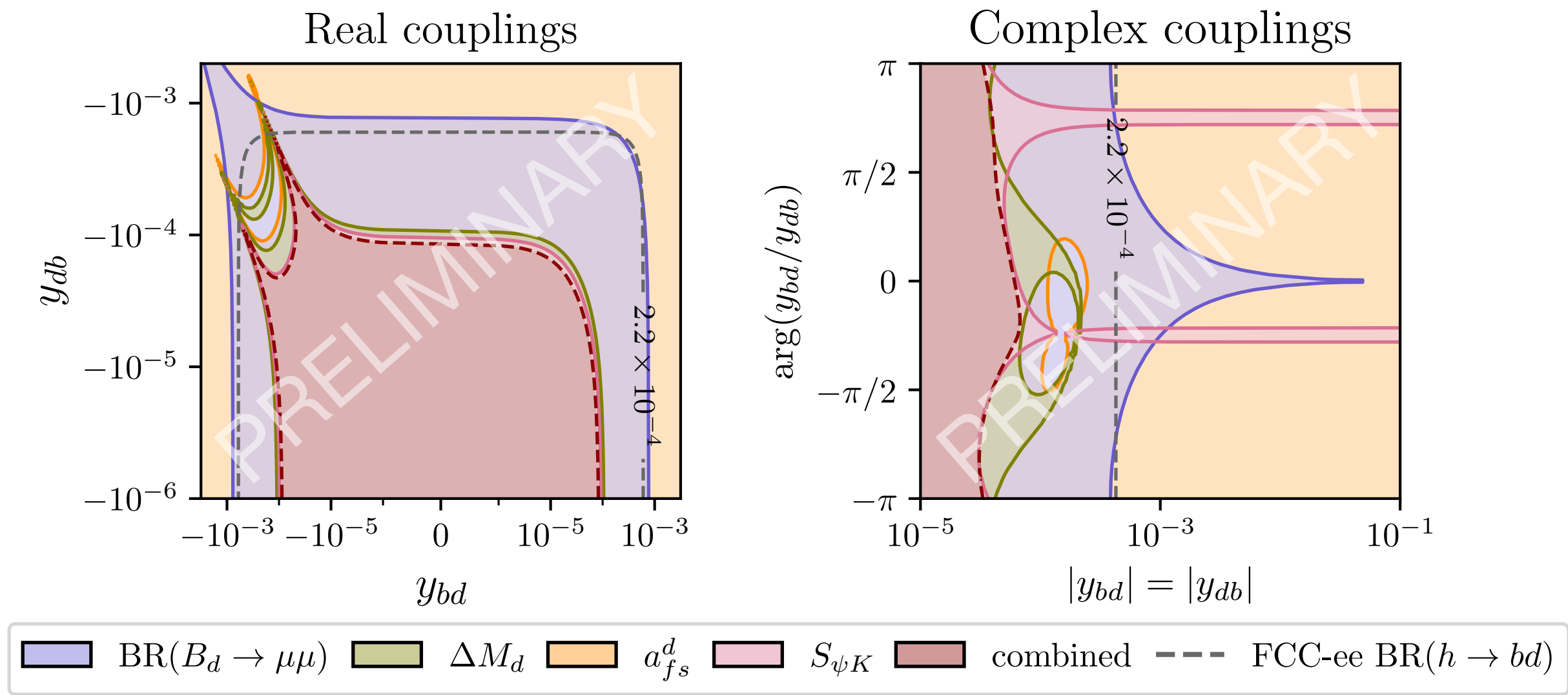


Upper limit **BR(h → bs)**

1.0×10^{-4}	1.3×10^{-3}	3.2×10^{-3}

	6.5×10^{-5}
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Bounds on y_{bd} and y_{db}



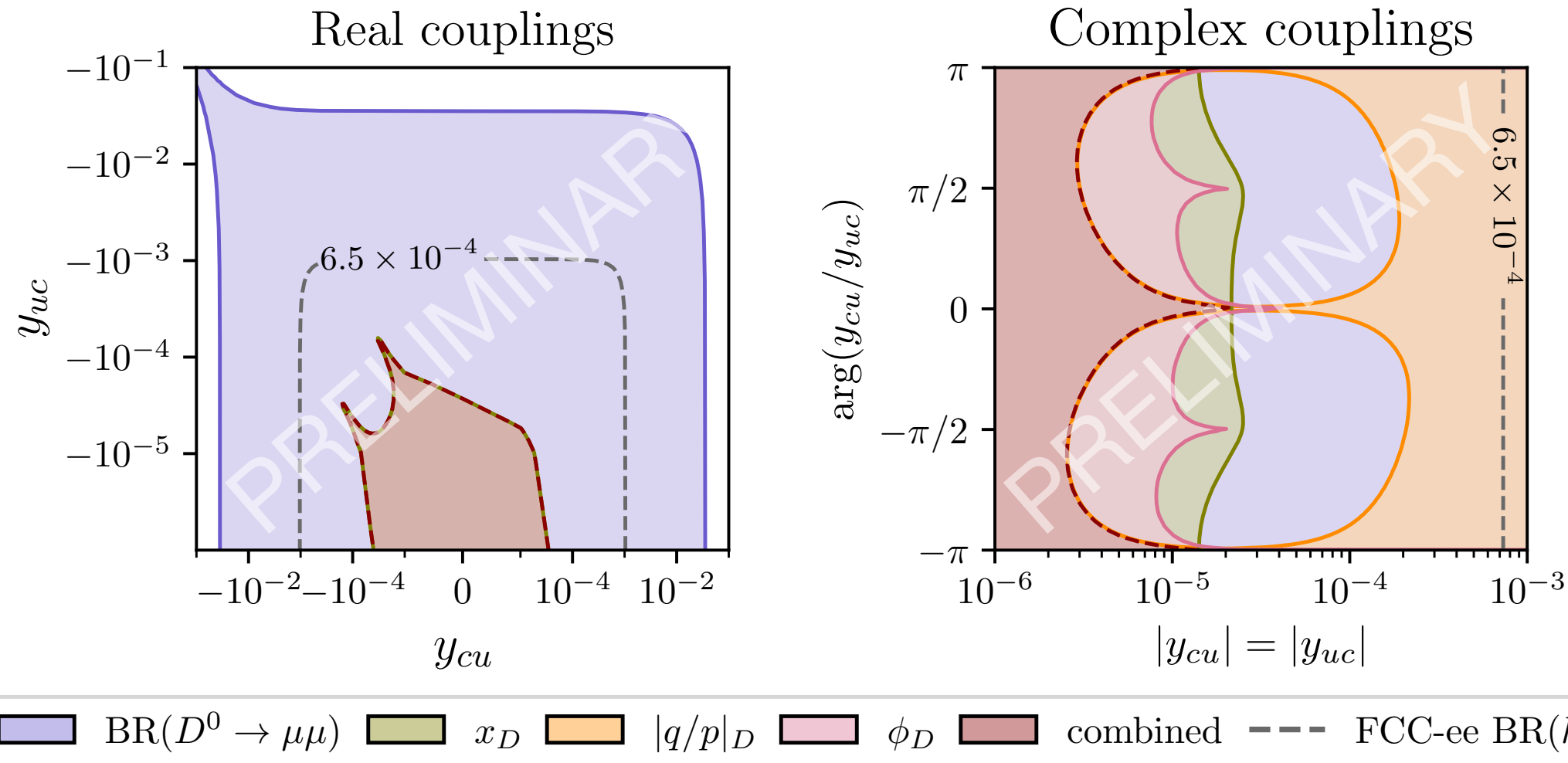
Upper limit **BR(h → bd)**

5.0 x 10 ⁻⁶	4.0 x 10 ⁻⁴

	5.5 x 10 ⁻⁶
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Bounds on y_{uc} and y_{cu}

- Challenge:** no reliable SM prediction $\rightarrow$ set **SM to zero** and inflate uncertainties to allow NP to saturate present experimental value within uncertainties

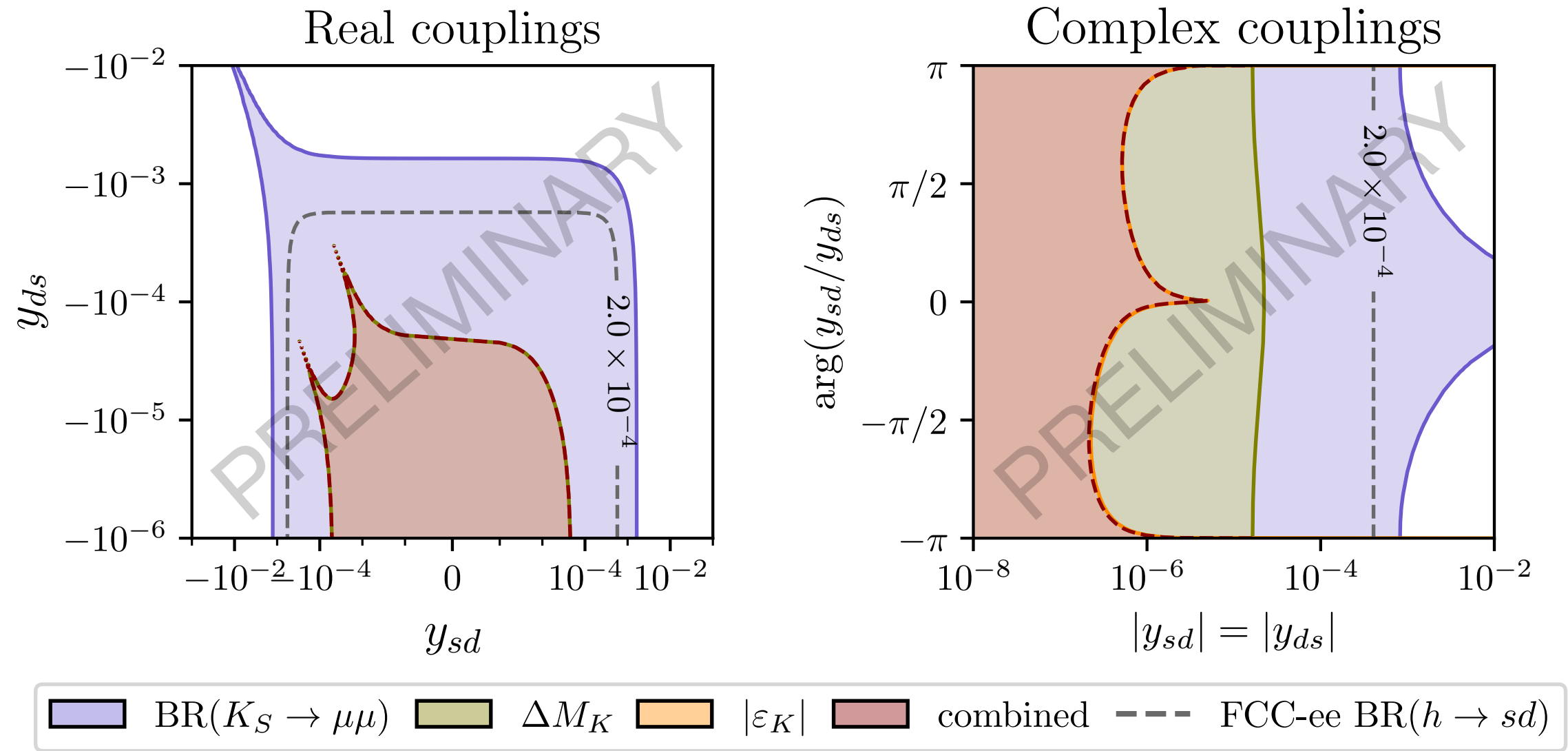


Upper limit $\text{BR}(h \rightarrow uc)$

1.0×10^{-6}	1.6×10^{-5}	8.0×10^{-1}

ξ	5.2×10^{-7}
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Bounds on y_{sd} and y_{ds}



Upper limit BR(h → sd)

1.1×10^{-6}	5.54×10^{-5}	1.5×10^{-3}

3.6×10^{-7}	3.1×10^{-8}

Projection of Bounds

Projections

- The mean of the projected (and current) observable measurement **set to SM**
- Projected uncertainties of CKM matrix parameters:

CKM parameter	$ V_{us} $	$ V_{ub} $	$ V_{cb} $	γ
Projected uncertainty	2.3×10^{-4}	3.2×10^{-5}	1.7×10^{-4}	3.5×10^{-3}
Current uncertainty	6.1×10^{-4}	1.4×10^{-4}	5.1×10^{-4}	6.3×10^{-2}

arXiv: [2006.04824](#)

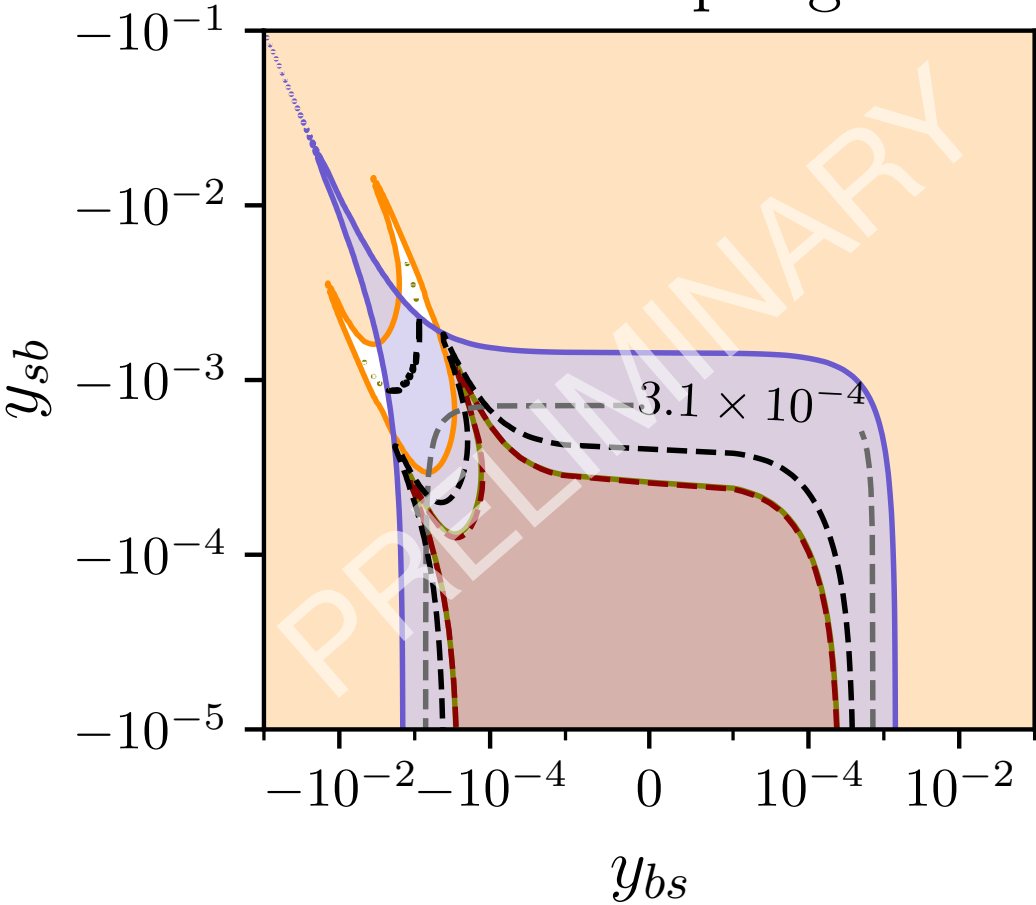
- Bag parameters $B_n^{(s(d))} \rightarrow$ **1% uncertainty**
- Projected uncertainties of (some) observable measurements:

Observable	$S_{\psi\phi}$	$S_{\psi K}$	a_{fs}^s	a_{fs}^d	$BR(B_s \rightarrow \mu\mu)$	$BR(B_d \rightarrow \mu\mu)$
Projected uncertainty	2.3×10^{-4}	8.0×10^{-4}	3.2×10^{-5}	3.2×10^{-5}	1.6×10^{-10}	1.2×10^{-11}
Current uncertainty	1.9×10^{-2}	1.7×10^{-2}	4.1×10^{-3}	1.7×10^{-3}	2.8×10^{-10}	5.4×10^{-11}

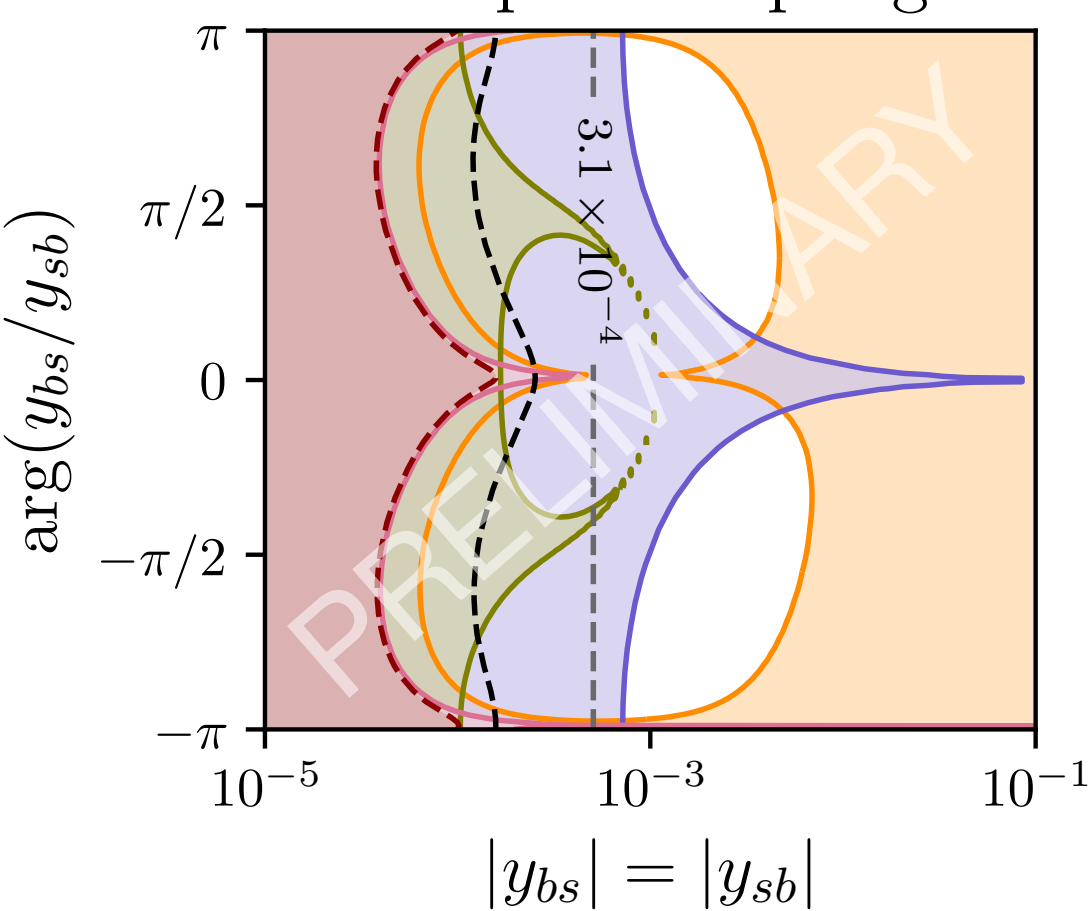
arXiv: [2006.04824](#), [ESPPU 2025 Calvi et. al.](#)

Bounds on y_{bs} and y_{sb}

Real couplings

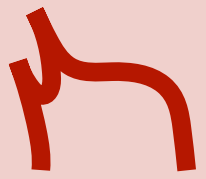


Complex couplings



BR($B_s \rightarrow \mu\mu$)
 ΔM_s
 a_{fs}^s
 $S_{\psi\phi}$
combined
FCC-ee $\text{BR}(h \rightarrow bs)$
current

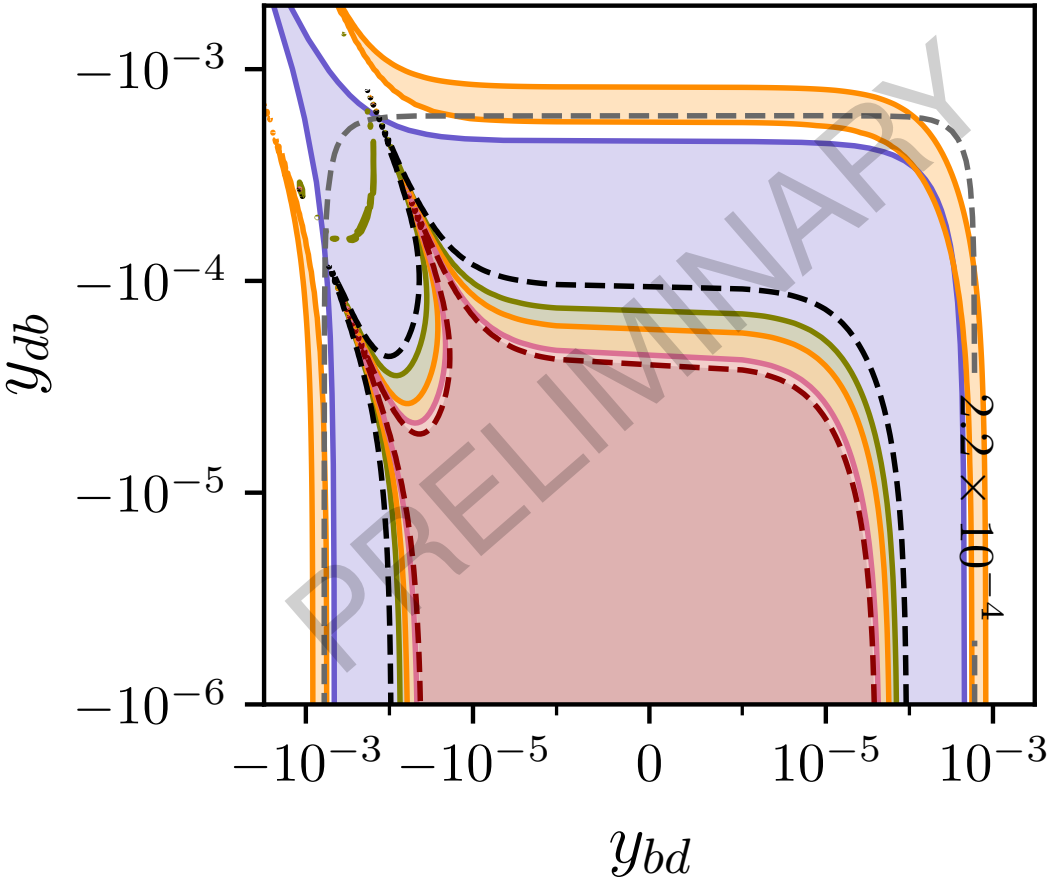
Upper limit $\text{BR}(h \rightarrow bs)$

	
4.0×10^{-5}	8.5×10^{-4}

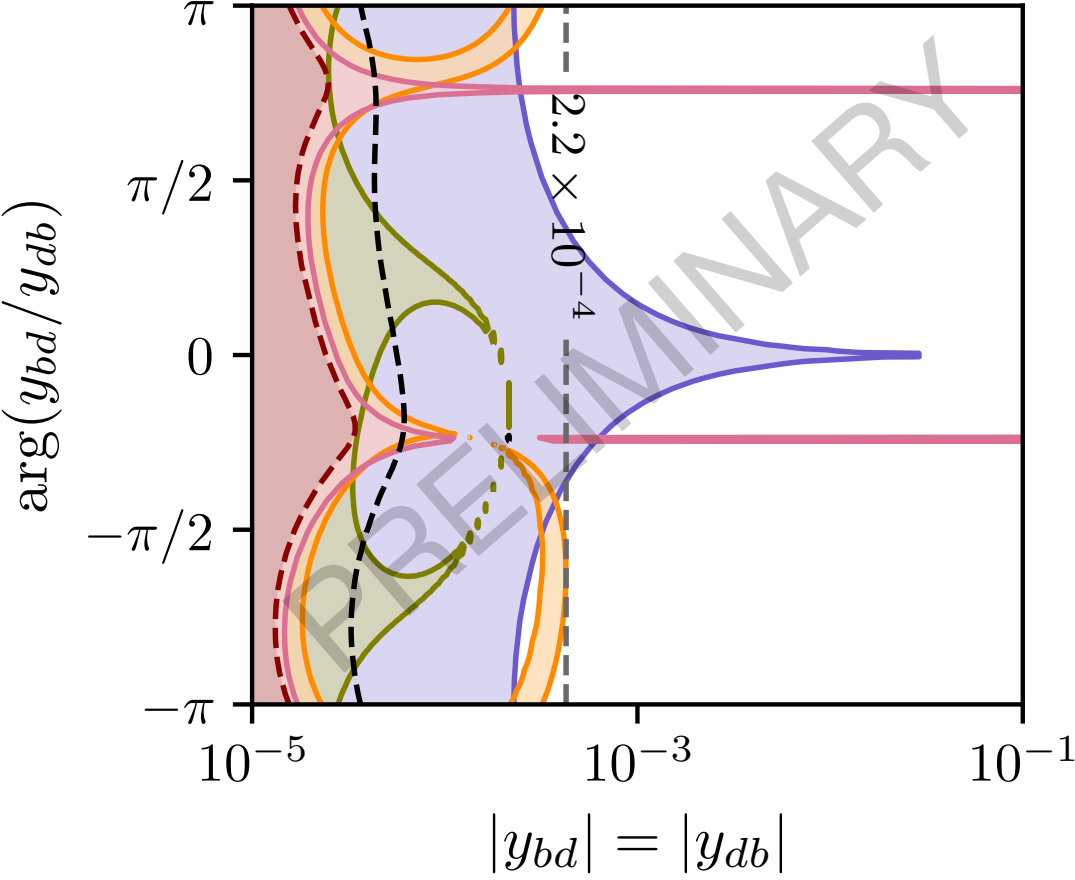
	3.0×10^{-5}
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Bounds on y_{bd} and y_{db}

Real couplings

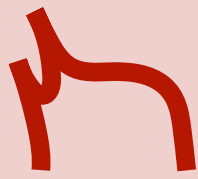


Complex couplings



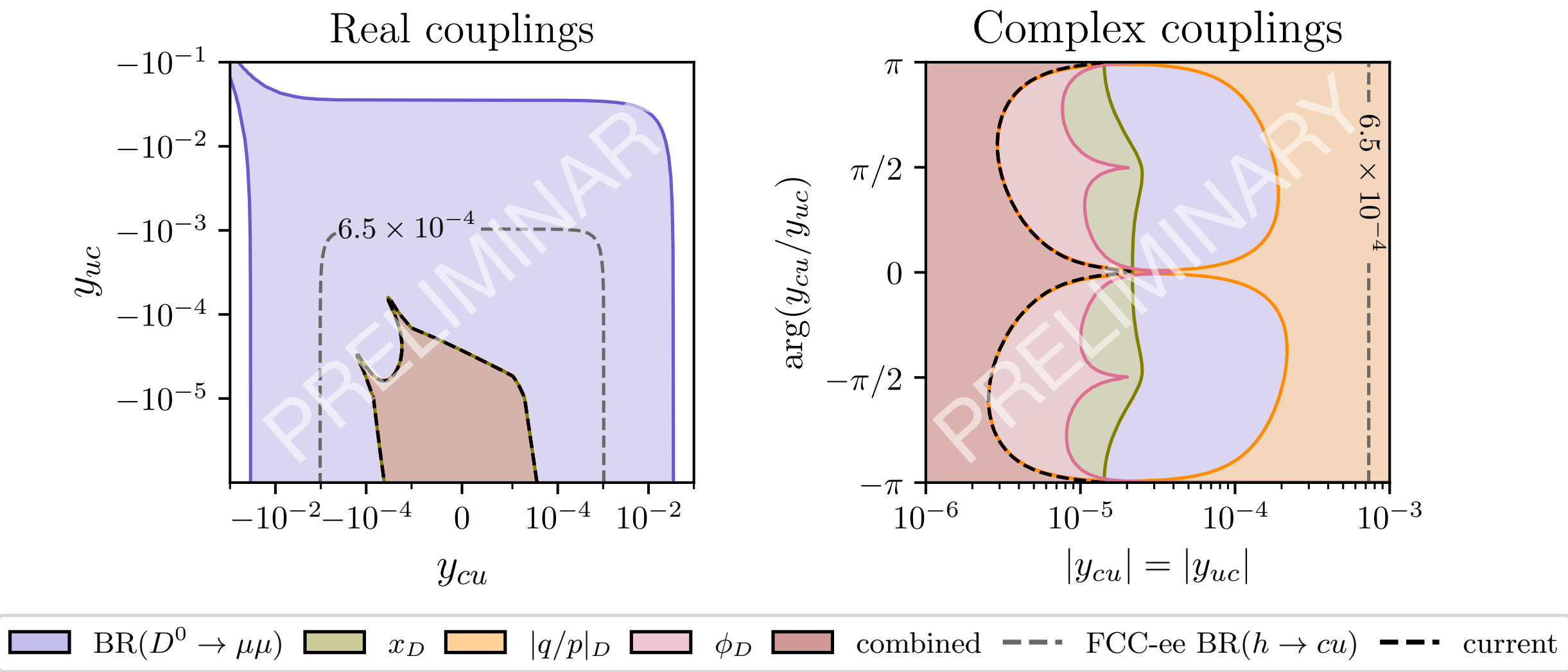
 $\text{BR}(B_d \rightarrow \mu\mu)$
 ΔM_d
 a_{fs}^d
 $S_{\psi K}$
 combined
 FCC-ee $\text{BR}(h \rightarrow bd)$
 current

Upper limit $\text{BR}(h \rightarrow bd)$

	
1.1×10^{-6}	7.8×10^{-5}

	1.4×10^{-6}
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Bounds on y_{uc} and y_{cu}

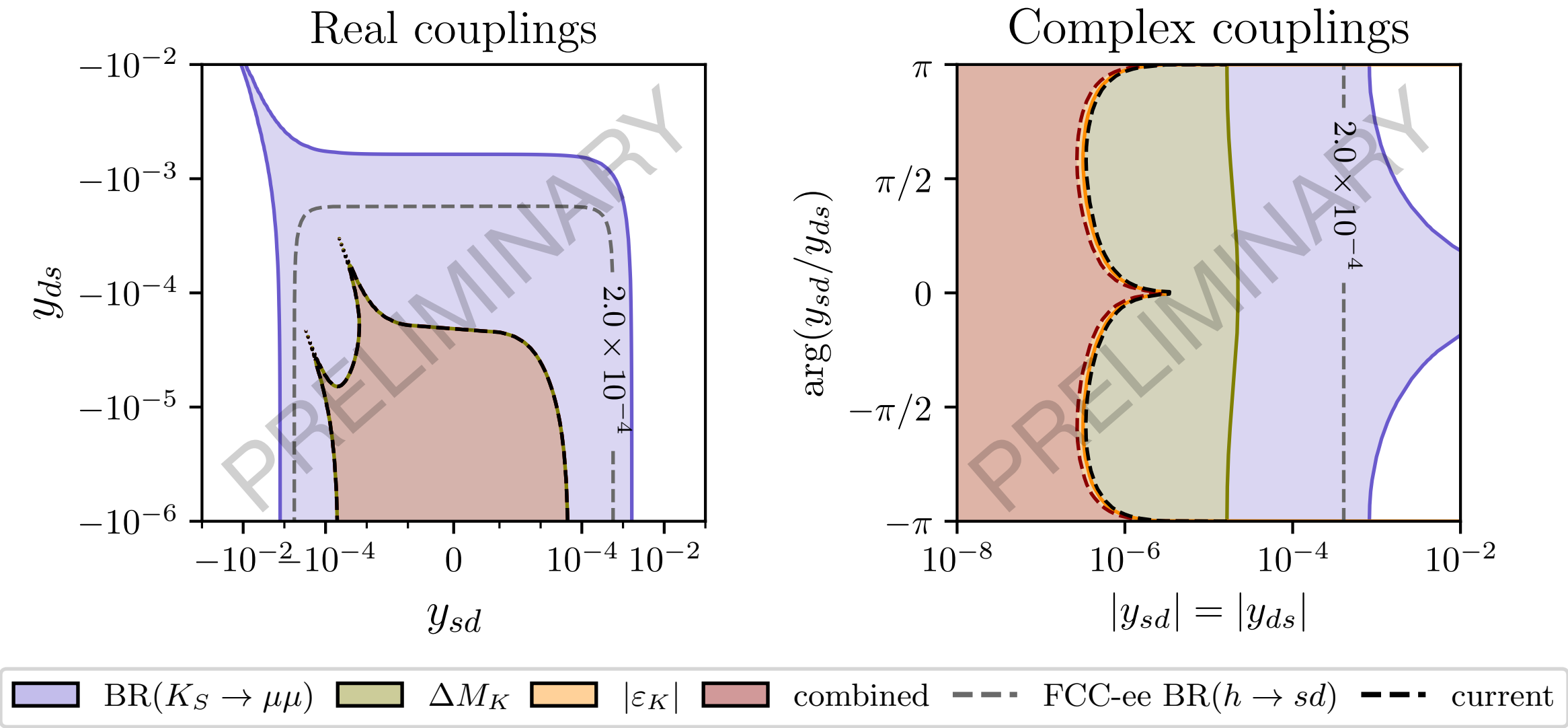


Upper limit **BR(h → uc)**

1.0×10^{-6}	1.6×10^{-5}	8.0×10^{-1}

Ξ	5.2×10^{-7}
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Bounds on y_{sd} and y_{ds}



Upper limit BR(h → sd)

1.1 x 10 ⁻⁶	5.54 x 10 ⁻⁵	1.5 x 10 ⁻³

3.3 x 10 ⁻⁷	8.5 x 10 ⁻⁹

Summary

Summary

- SM prediction of Higgs FCNCs is **tiny**
- Projected direct bounds on FCC-ee relatively **big**
- Indirect current bounds **mostly stronger** than projected direct bounds
- Future improvement of the indirect bounds
- All bounds **much higher** than SM prediction
- If any Higgs FCNCs measured → **New Physics**

Additional Slides

Matching the Effective Lagrangian to SMEFT

- SMEFT Lagrangian: $\mathcal{L}_{\text{SMEFT}} \supset -\frac{C_{dH}^{ij}}{\Lambda^2} \bar{Q}_L^i H d_R^j + \text{h.c.}$
- SM Lagrangian: $\mathcal{L}_{\text{SM}} \supset -y_{ij}^{SM} \bar{Q}_L^i H d_R^j + \text{h.c.}$

- Expand Higgs in unitary gauge: $H \rightarrow \begin{pmatrix} 0 \\ v + h \end{pmatrix}$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SMEFT}} \supset -\frac{1}{\sqrt{2}} \underbrace{\left(y_{ij}^{SM} v + \frac{v^3}{2\Lambda^2} C_{dH}^{ij} \right)}_{m_{ij}} \bar{d}_L^i d_R^j - \frac{1}{\sqrt{2}} \left(y_{ij}^{SM} + \frac{3v^2}{2\Lambda^2} C_{dH}^{ij} \right) \bar{d}_L^i d_R^j h$$

- The effective Lagrangian: $\mathcal{L}_{\text{eff}} \supset y_{ij} \bar{d}_L^i d_R^j h + \text{h.c.}$

$$\longrightarrow y_{ij} = -\frac{m_{ij}}{v} + \frac{1}{\sqrt{2}} \frac{v^2}{\Lambda^2} C_{dH}^{ij}$$

Definitions of the Mixing Observables

$$\frac{d}{dt} \begin{pmatrix} |P^0(t)\rangle \\ |\bar{P}^0(t)\rangle \end{pmatrix} = \left(-iM - \frac{1}{2}\Gamma \right) \begin{pmatrix} |P^0(t)\rangle \\ |\bar{P}^0(t)\rangle \end{pmatrix}$$

- $B_{s(d)} - \bar{B}_{s(d)}$ mixing observables:

$$\Delta M_q = M_H - M_L = 2\sqrt{\frac{|M_{12}|^2 - 1/4|\Gamma_{12}|^2 + \sqrt{(|M_{12}|^2 - 1/4|\Gamma_{12}|^2)^2 + 1/4(M_{12}\Gamma_{12}^* + M_{12}^*\Gamma_{12})^2}}{2}}$$

$$a_{fs}^q = \left| \frac{M_{12}}{\Gamma_{12}} \right| \sin(\arg(-M_{12}/\Gamma_{12})) \quad S_{\psi K(\phi)} = \sin(2\beta_{(s)})$$

- $D^0 - \bar{D}^0$ mixing observables:

$$x_D = \frac{\Delta M_D}{\Gamma_D} \quad y_D = \frac{\Delta \Gamma_D}{2\Gamma_D} \quad \frac{q}{p} = \sqrt{\left(\frac{M_{12}^* - i/2\Gamma_{12}^*}{M_{12} - i/2\Gamma_{12}} \right)} \quad \phi_D = \arg\left(\frac{q}{p}\right)$$

- $K^0 - \bar{K}^0$ mixing observables:

$$\varepsilon_K = \frac{A(K_L \rightarrow \pi\pi)}{A(K_S \rightarrow \pi\pi)} \quad \Delta M_K = M_{K_L} - M_{K_S}$$